Tutorial for "How topological recursion organises quantum fields on noncommutative geometries"

This tutorial aims to understand in more details the derivation of the (closed) Dyson-Schwinger equations related to Raimar Wulkenhaar's Mini-course "How topological recursion organises quantum fields on noncommutative geometries".

Let H_N be the space of hermitian $N \times N$ matrices, then we are considering the following partition function

$$\mathcal{Z}^{k}[J] = \int_{H_{N}} dM \exp\left(-N\mathrm{Tr}\left(EM^{2} + \frac{\lambda}{k}M^{k} - JM\right)\right),$$

where $E, J \in H_N$ and E has distinct positive eigenvalues $(E_i)_i$. Denote by $S_{int}^k(M) = N \operatorname{Tr}\left(\frac{\lambda}{k}M^k\right)$ the interaction. We are in particular interested in \mathcal{Z}^3 (Kontsevich model) and \mathcal{Z}^4 (Grosse-Wulkenhaar model).

The correlation functions (the so-called $(N_1 + ... + N_b)$ -point function of genus g with b boundaries):

$$G^{(g)}_{|a_1^1 a_2^1 \dots a_{N_1}^1| \dots |a_1^b \dots a_{N_b}^b|} := N^{2g+b-2} \frac{\partial^{N_1 + \dots + N_b}}{\partial J_{a_1^1 a_2^1} \dots \partial J_{a_{N_1}^1 a_1^1} \dots \partial J_{a_{N_b}^b a_1^b}} \log \mathcal{Z}^k[J] \big|_{J=0}$$

where we assume for the definiton that all a_i^j are pairwise distinct.

Supporting Exercises

Exercise/Remark 1: Show that $\mathcal{Z}^3[0]$ is equivalent to

$$\mathcal{Z}^{3}[0] = C \int_{H_{N}} d\tilde{M} \exp\left(-N \operatorname{Tr}\left(\tilde{E}\tilde{M} + \tilde{\lambda}\tilde{M}^{3}\right)\right)$$

by determining $C, \tilde{E}, \tilde{\lambda}, \tilde{M}$. There is the notation of generalised Kontsevich model with the partition function

$$\mathcal{Z}^{gKont} = \int_{H_N} d\tilde{M} \exp\bigg(-N\mathrm{Tr}\big(\tilde{E}\tilde{M} + V(\tilde{M})\big)\bigg),$$

where V(x) is a polynomial with real coefficients. Show that \mathcal{Z}^4 can not be transformed to \mathcal{Z}^{gKont} . From this point of view \mathcal{Z}^k with k > 3 is a different type of generalisation of the classical Kontsevich model.

Exercise 2: Show that the partition function can be represented as

$$\mathcal{Z}^{k}[J] = K \exp\left(-S_{int}\left(\frac{1}{N}\frac{\partial}{\partial J}\right)\right) \exp\left(\frac{N}{2}\sum_{n,m}\frac{J_{nm}J_{mn}}{E_{n}+E_{m}}\right),$$

where K is some constant depending on E.

Exercise 3: Show that the correlation function is the connected expectation values

$$\frac{1}{N^K} \frac{\partial^K}{\partial J_{p_1q_1} \partial J_{p_2q_2} \dots \partial J_{p_Kq_K}} \log \mathcal{Z}[J] \Big|_{J=0} = \langle M_{q_1p_1} M_{q_2p_2} \dots M_{q_Kp_K} \rangle_c.$$

Exercise 4: Find an argument why any $(N_1 + ... + N_b)$ -point function with $\sum_{i=1}^b N_b$ odd vanishes for \mathcal{Z}^k , whenever k is even.

Exercise 5: Prove the Leibniz rule

$$e^{f(\partial_x)}(x \cdot g(x)) = f'(\partial_x)e^{f(\partial_x)}g(x) + xe^{f(\partial_x)}g(x)$$

Exercise 6: Compute the Ward identity

$$\frac{E_a - E_b}{N} \sum_m \frac{\partial}{\partial J_{am}} \frac{\partial}{\partial J_{mb}} \mathcal{Z}^k[J] = \sum_m \left(J_{ma} \frac{\partial}{\partial J_{mb}} - J_{bm} \frac{\partial}{\partial J_{am}} \right) \mathcal{Z}^k[J].$$

by considering invariance of the partition function under unitary transformation. Le $U = e^{iA} \in U(N)$, choose an infinitesimal transformation of the form $M \mapsto M' = UMU^{\dagger} = M + iAM - iMA + \mathcal{O}(A^2)$.

Exercise 7: Prove the recursive algebraic relation between correlation functions for $N_i > k - 2$ boundary length.

For k = 3

$$G_{|a_1^1..a_{N_1}^1|\mathcal{J}|} = -\lambda \frac{G_{|a_2^1a_3^1..a_{N_1}^1|\mathcal{J}|} - G_{|a_1^1a_3^1..a_{N_1}^1|\mathcal{J}|}}{(E_{a_1^1} + E_{a_2^1})(E_{a_1^1} - E_{a_2^1})}.$$
(1)

For k = 4

$$\begin{split} G_{|a_{1}^{1}a_{2}^{1}..a_{N_{1}}^{1}|\mathcal{J}|} &= -\frac{\lambda}{E_{a_{2}^{1}} - E_{a_{N_{1}}^{1}}} \bigg\{ \frac{1}{N^{2}} \sum_{k=2}^{N_{1}} \frac{G_{|a_{2}^{1}a_{3}^{1}..a_{k}^{1}|a_{k+1}^{1}..a_{N}^{1}a_{1}^{1}|\mathcal{J}|} - G_{|a_{1}^{1}a_{2}^{1}a_{3}^{1}..a_{k-1}^{1}|a_{k}^{1}..a_{N_{1}}^{1}|\mathcal{J}|} \\ &+ \sum_{\beta=2}^{b} \sum_{k=1}^{N_{\beta}} \frac{G_{|a_{1}^{\beta}a_{2}^{\beta}..a_{k}^{\beta}a_{2}^{1}a_{3}^{1}..a_{N_{1}}^{1}a_{1}^{1}a_{k+1}^{\beta}..a_{N_{\beta}}^{\beta}|\mathcal{J}\setminus\{J^{\beta}\}|}{E_{a_{k}^{\beta}} - E_{a_{1}^{1}}} \\ &+ \sum_{k=2}^{N} \sum_{\mathcal{I} \uplus \mathcal{I}'=\mathcal{J}} \frac{G_{|a_{k+1}^{1}..a_{N_{1}}^{1}a_{1}^{1}|\mathcal{I}|}G_{|a_{2}^{1}..a_{k}^{1}|\mathcal{I}'|} - G_{|a_{k}^{1}..a_{N_{1}}^{1}|\mathcal{I}|}G_{|a_{1}^{1}a_{2}^{1}..a_{k-1}^{1}|\mathcal{I}'|}}{E_{a_{k}^{1}} - E_{a_{1}^{1}}} \bigg\}. \end{split}$$