# Blobbed Topological Recursion for Dirac Ensembles

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• This talk is based on joint work with Masoud Khalkhali, Random Finite Noncommutative Geometries and Topological Recursion, arXiv:1906.09362

• The major ideas discussed through the talk are certainly due to the following mathematicians (just to name a few):

- (Blobbed) Topological Recursion: Eynard, Chekhov, Borot, Orantin, ...
- Noncommutative Geometry: Connes, Marcolli, Barrett, ...

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#### 2 Dirac Ensembles

3 Schwinger-Dyson Eq. and (Blobbed) Topological Recursion

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### Outline

#### 1 Classical 1-Hermitian Matrix Models

#### 2 Dirac Ensembles

#### 3 Schwinger-Dyson Eq. and (Blobbed) Topological Recursion

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## **Classical 1-Hermitian Matrix Models**

• One starts from a  $U_N$ -invariant measure  $d\mu(H)$  on  $\mathcal{H}_N$ 

$$d\mu(H) = \exp\left[-(N/t)\operatorname{Tr}(\mathcal{V}(H))\right] dH$$

• Weyl integration formula:  $d\mu(H)$  induces a measure  $d\tilde{\mu}(\boldsymbol{\lambda})$  on  $\mathbb{R}^N$ 

$$d\tilde{\mu}(\boldsymbol{\lambda}) = \prod_{1 \le i < j \le N} |\lambda_j - \lambda_i|^2 \prod_{i=1}^N e^{-(N/t)\mathcal{V}(\lambda_i)} d\lambda_i$$

• Formal matrix models: all integrations are w.r.t. the normalized Gaussian measure

$$d\mu^0(H) = c \exp\left(-\frac{N\operatorname{Tr}(H^2)}{2t}\right) dH$$

## **Correlation Functions**

- Partition function  $Z_N = \int_{\mathcal{H}_N} d\mu(H)$ , and Free energy  $F = \log Z_N$
- Disconnected correlators are the moments of the following form

$$\hat{W}_n(x_1, \cdots, x_n) = \mathbb{E}\left[\prod_{j=1}^n \left(\sum_{i=1}^N \frac{1}{x_j - \lambda_i}\right)\right], \quad x_j \in \mathbb{C} \setminus \mathbb{R},$$

where  $\sum_{i=1}^{N} \frac{1}{x-\lambda_i}$  is the trace of the resolvent.

• Connected correlators are the joint cumulants of the following form

$$W_n(x_1, \cdots, x_n) = \sum_{K \vdash [1,n]} (-1)^{[K]-1} ([K]-1)! \prod_{i=1}^{[K]} \hat{W}_{|K_i|} (x_{K_i})$$

# Counting Discretized Surfaces

- Wick's Theorem: computation of the moments of the Gaussian measure can be reduced to enumeration of the maps (ribbon graphs)
- A term  $\tau_{\ell_i} = t_{\ell_i} \frac{\operatorname{Tr}(H^{\ell_i})}{\ell_i}$  in  $\operatorname{Tr}(\mathcal{V}(H)) \rightsquigarrow$  an  $\ell_i$ -gon of Boltzmann weight  $t_{\ell_i}$



Figure: A planar map with one marked rooted face (the colored 9-gon containing the point  $\infty$ )

# Topological Expansion

 $\bullet\,$  Using Wick's theorem, one obtains a large-N topological expansion

$$F = \sum_{g \ge 0} \left( N/t \right)^{2-2g} F_g , \qquad F_g = \sum_{[M] \in \mathbb{M}_g} \mathfrak{Bw}([M]) ,$$

$$W_n(x_1, \cdots, x_n) = \sum_{g \ge 0} (N/t)^{2-2g-n} W_{g,n}(x_1, \cdots, x_n),$$

where  $W_{g,n}(x_1, \dots, x_n)$  enumerates the connected closed maps of genus g with n marked rooted polygonal faces

Topological Recursion provides a machinery for computing the W<sub>g,n</sub>'s recursively, given certain initial data



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## **Real Spectral Triples**

- A spectral triple  $(\mathcal{A}, \mathcal{H}, D)$  equipped with a real structure  $J : \mathcal{H} \to \mathcal{H}$  and a chirality operator  $\gamma : \mathcal{H} \to \mathcal{H}$
- The real spectral triple  $(C^{\infty}(\mathcal{M}), L^2(\mathcal{M}, \$), \not D, \gamma, J)$  associated to a spin Riemannian manifold  $\mathcal{M}$
- Connes' distance formula: The Dirac operator D encapsulates all the information about the Riemannian metric over M

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### **Dirac Ensembles**

- Moduli space of Dirac operators D encodes all possible geometries ("metrics") over a fixed fermion space (A, H, γ, J)
- As a model for *Quantum Gravity* on a finite noncommutative space, one considers a distribution of the following form over the moduli space of Dirac operators

 $e^{-\mathcal{S}(D)} \,\mathrm{d}D$ 

• The action functional is *spectral* 

$$\mathcal{S}(D) = \operatorname{Tr}(f(D)) = \sum_{\lambda \in \operatorname{Spec}(D)} f(\lambda)$$

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# Matrix Geometries of type (p, q) [Barrett (2015)]

- A particular class of finite real spectral triples
  - $\mathcal{A} = M_N(\mathbb{C})$
  - $\mathcal{H} = V_{p,q} \otimes M_N(\mathbb{C})$
  - $\langle v \otimes A, u \otimes B \rangle = \langle v, u \rangle \operatorname{Tr} (AB^*) , \quad v, u \in V_{p,q} , A, B \in M_N(\mathbb{C})$
  - $\pi(A)(v \otimes B) = v \otimes (AB)$  ,

where  $V_{p,q}$  is an irreducible complex module over the Clifford algebra  $\mathbf{C}\ell_{p,q}$ 

 The Dirac operators are expressed in term of commutators or anticommutators with given Hermitian matrices H and anti-Hermitian matrices L

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## Random matrix geometries of type (1,0)

• The Dirac operator

$$D = \{H, \cdot\}, \quad H \in \mathcal{H}_N$$

• We consider a model with the following action functional

$$\mathcal{S}(D) = \mathcal{S}_{\text{unstable}}(D) + \mathcal{S}_{\text{stable}}(D),$$

where

$$\mathcal{S}_{\text{unstable}}(D) = \text{Tr}\left(\mathcal{V}(D)\right), \quad \mathcal{V}(x) = \frac{1}{2t} \left(\frac{x^2}{2} - \sum_{n=3}^d \alpha_n \frac{x^n}{n}\right),$$

and

$$\mathcal{S}_{\text{stable}}(D) = -\sum_{s=1}^{\mathfrak{g}} \left(N/t\right)^{-4s} \sum_{n_I \in \mathbb{N}^s_{\uparrow}} \hat{\alpha}_{n_I} \prod_{i=1}^s \operatorname{Tr}\left(D^{n_i}\right) \,.$$

## Topological expansion of the action functional

An elementary 2-cell of topology (g, n) with polygonal boundaries of perimeters {ℓ<sub>i</sub>}<sup>n</sup><sub>i=1</sub>, ℓ<sub>i</sub> ≥ 1, is a (equivalence class of) surface of genus g whose boundary has n connected components, and consists of the 1-skeleton of ℓ<sub>i</sub>-gons



Figure: An elementary 2-cell of topology (g, n) = (3, 2) with polygonal boundaries of perimeters  $(\ell_1, \ell_2) = (5, 6)$ 

• Decomposition of  $\mathcal{S}(D)$  in terms of underlying elementary 2-cells

$$\begin{split} \mathcal{S}(D) &= \mathcal{S}_0(H) + \mathcal{S}_{\text{int}}(H) \\ &= \frac{N}{2t} \operatorname{Tr} \left( H^2 \right) - \sum_{[C] \in \mathcal{C}} \frac{(N/t)^{\chi(C)}}{(\beta_0 \ (\partial C))!} \ T_{[C]}(H) \,, \end{split}$$

where

- $\chi(C)$  = the Euler characteristic of an elementary 2-cell C
- $\beta_0(\partial C) =$  the number of connected components of the boundary of C
- Classifying the elementary 2-cells based on whether  $\chi(C) \ge 0$  or  $\chi(C) < 0$

$$\mathcal{C} = \mathcal{C}_{\mathrm{unstable}} \cup \mathcal{C}_{\mathrm{stable}}$$

• For an elementary 2-cell C of topology (g, n) with polygonal boundaries of perimeters  $\{\ell_i\}_{i=1}^n$ 

$$T_{[C]}(H) \coloneqq \mathbf{t}_{\vec{\ell}}^{(g)} \prod_{i=1}^{n} \frac{\operatorname{Tr}\left(H_{i}^{\ell}\right)}{\ell_{i}}$$
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# Corresponding Hermitian matrix model

- A formal multi-trace 1-Hermitian matrix model
- The term  $\exp(-S_{int}(H))$  is considered as a formal power series in the Boltzmann weights  $t^{(g)}_{\vec{r'}}$
- $\bullet\,$  Using Wick's theorem, we get the following large- N topological expansion

$$W_n(x_1, \cdots, x_n) = \sum_{g \ge 0} (N/t)^{2-2g-n} W_{g,n}(x_1, \cdots, x_n),$$

where  $W_{g,n}(x_1, \dots, x_n) \in \mathbb{Q}[\mathbf{t}][[t]][[(x_j^{-1})_j]]$  enumerates the connected closed stuffed maps of genus g with n marked rooted polygonal faces

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# Stuffed maps

• Each term of the form  $\frac{(N/t)^{\chi(C)}}{(\beta_0(\partial C))!} T_{[C]}(H)$  is represented by the corresponding elementary 2-cell C



Figure: A closed stuffed map of genus two with two marked rooted polygonal faces (brown disk) obtained by gluing the unstable elementary 2-cells of topology (g, n) = (0, 1) (green disk) and (g, n) = (0, 2) (red cylinder)

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# Large-N expected spectral distribution ("equilibrium measure")

- Tame Boltzmann weights are those numerical values of  $t_{\ell}^{(g)}$  for which the generating function of the rooted planar stuffed maps with topology of a disk and perimeter  $\ell$  is finite for all  $\ell \in \mathbb{N}$ .
- The formal series  $W_{0,1}(x)$  is a holomorphic function with discontinuity locus  $\Gamma = [\mathfrak{a}, \mathfrak{b}] \subset \mathbb{R}$ .
- The large-N expected spectral density

$$\varphi(s) = \frac{1}{2\pi i} \lim_{\epsilon \to 0^+} \left( W_{0,1}(s - i\epsilon) - W_{0,1}(s + i\epsilon) \right), \quad \forall s \in \Gamma_{\text{interior}}$$

is supported on  $\Gamma$ , and assumes positive values on  $\Gamma_{\text{interior}}$ .

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## Schwinger-Dyson Equations (SDEs)

- The Schwinger-Dyson equations for a matrix model are a tower of equations satisfied by the *n*-point correlation functions of the model.
- The root of SDEs is the invariance of the integral of a top-degree differential form under a 1-parameter family of orientation-preserving diffeomorphisms on a manifold, i.e.

$$\int_{\phi_t(\Omega)} \Psi \omega = \int_{\Omega} \phi_t^{*} \left( \Psi \omega \right) \,, \quad \forall t \in \left( -\epsilon, \epsilon \right),$$

where  $\Psi : \mathcal{M} \to \mathbb{R}$  is a smooth function over a Riemannian *n*-manifold  $\mathcal{M}$ with the canonical volume form  $\omega$ , and  $\phi_t : \Omega \to \Omega$  is a local flow over a compact *n*-dimensional submanifold  $\Omega \subset \mathcal{M}$ .

## SDEs for Random Matrix Geometries of type (1,0)

The rank n SDE is given by:

$$\begin{split} W_{n+1}(x, x, x_I) &+ \sum_{J \subseteq I} W_{|J|+1}(x, x_J) W_{n-|J|}(x, x_{I\setminus J}) \\ &+ \sum_{i \in I} \oint_{\Gamma} \frac{\mathrm{d}\xi}{2\pi \mathrm{i}} \frac{W_{n-1}(\xi, x_{I\setminus \{i\}})}{(x-\xi)(x_i-\xi)^2} \\ &+ \sum_{k=1}^{2\mathfrak{g}} \sum_{\substack{K \vdash [\![1,k]]\\J_1 \sqcup \cdots \sqcup J_{[K]} = I}} \oint_{\Gamma} \left[ \prod_{r=1}^k \frac{\mathrm{d}\xi_r}{2\pi \mathrm{i}} \right] \frac{\partial_{\xi_1} T_k(\xi_1, \cdots, \xi_k)}{(k-1)! (x-\xi_1)} \prod_{i=1}^{[K]} W_{|K_i|+|J_i|} \left(\xi_{K_i}, x_{J_i}\right) \\ &= 0 \,, \end{split}$$

where the symmetric k-point interactions  $T_k$  are defined by

$$\sum_{\lambda} T_k(\lambda_{i_1}, \lambda_{i_2}, \cdots, \lambda_{i_k}) = \sum_{\substack{[C] \in \mathcal{C} \\ \beta_0(\partial C) = k}} (N/t)^{\chi(C)} T_{[C]}(H) .$$

By considering a large-N expansion of topological type for the correlation functions  $W_n$  and the k-point interactions  $T_k$ , the rank n SDE to order  $N^{3-2g-n}$  is given by:

$$\begin{split} W_{g-1,n+1}(x,x,x_{I}) &+ \sum_{J \subseteq I, \ 0 \leqslant f \leqslant g} W_{f,\,|J|+1}(x,x_{J}) \, W_{g-f,\,n-|J|}(x,x_{I\setminus J}) \\ &+ \sum_{i \in I} \oint_{\Gamma} \frac{\mathrm{d}\xi}{2\pi \mathrm{i}} \, \frac{W_{g,\,n-1}(\xi,x_{I\setminus\{i\}})}{(x-\xi)(x_{i}-\xi)^{2}} \\ &+ \sum_{\substack{1 \leqslant k \leqslant 2\mathfrak{g} \\ 0 \leqslant h}} \sum_{\substack{K \vdash [\![1,k]\!] \\ J_{1} \sqcup \cdots \sqcup J_{[K]} = I}} \sum_{\substack{0 \leqslant f_{1}, \cdots, f_{[K]} \\ h+k-[K] + \sum_{i} f_{i} = g}} \\ &\int_{\Gamma} \Big[ \prod_{r=1}^{k} \frac{\mathrm{d}\xi_{r}}{2\pi \mathrm{i}} \Big] \, \frac{\partial_{\xi_{1}} T_{h,k}(\xi_{1}, \cdots, \xi_{k})}{(k-1)! \, (x-\xi_{1})} \, \prod_{i=1}^{[K]} W_{f_{i},\,|K_{i}|+|J_{i}|} \, (\xi_{K_{i}}, x_{J_{i}}) \end{split}$$

= 0.

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# Spectral Curve

• The rank one SDE to leading order in N

$$(W_{0,1}(x))^2 + \sum_{k=1}^2 \oint_{\Gamma} \left[ \prod_{r=1}^k \frac{\mathrm{d}\xi_r}{2\pi \mathrm{i}} \right] \frac{\partial_{\xi_1} T_{0,k}(\xi_1, \cdots, \xi_k)}{x - \xi_1} \prod_{r=1}^k W_{0,1}(\xi_r) = 0.$$

• The Stieltjes transform  $W_{0,1}(x)$  of the large-N spectral distribution  $\mu = \varphi(s) \,\mathrm{d}s$  of the model satisfies a quadratic algebraic equation:

$$W_{0,1}(x) = Q(x) + M(x)\sqrt{(x-\mathfrak{a})(x-\mathfrak{b})},$$

where the coefficients of the polynomials Q(x) and M(x) depend on the Boltzmann weights and the moments of  $\mu$ .

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# Spectral Curve

• Using the Joukowski map  $x : \mathbb{C} \setminus \{0\} \to \mathbb{C}$ , given by

$$x(z) = \frac{\mathfrak{a} + \mathfrak{b}}{2} + \frac{\mathfrak{b} - \mathfrak{a}}{4} \left( z + \frac{1}{z} \right)$$

the function  $W_{0,1}(x(z))$  gets an *analytic continuation* over the spectral curve  $\Sigma$  of the model.



Figure: Illustration of the Joukowski map an the spectral curve  $\Sigma$  of the model

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• From the coefficients  $W_{g,n}$  of the correlation functions to meromorphic symmetric differentials  $\omega_{g,n}$  of degree n, i.e. sections of the *n*-times external tensor product

 $K_{\Sigma}^{\boxtimes n} \to \Sigma^n$  of the canonical line bundle  $K_{\Sigma} \to \Sigma$ , given by

$$\omega_{g,n}(z_1, \cdots, z_n) = W_{g,n}(x(z_1), \cdots, x(z_n)) dx(z_1) dx(z_2) \cdots dx(z_n) + \delta_{n,2} \,\delta_{g,0} \, \frac{dx(z_1) \, dx(z_2)}{(x(z_1) - x(z_2))^2}$$

- Input for the (Blobbed) Topological Recursion Formula
  - The Riemann surface  $\Sigma$  equipped with a local biholomorphic involution
  - The 1-form  $\omega_1^0(z)$
  - The symmetric bidifferential  $\omega_2^0(z, z_1)$

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## Blobbed Topological Recursion [Borot (2014)]

#### Main Result

For random matrix geometries of type (1,0) with the distribution  $d\rho = e^{-S(D)} dD$ , all the stable  $\omega_{g,n}$ , 2g + n - 2 > 0, can be computed recursively, using the blobbed topological recursion formula given by

$$\omega_{g,n}(z,z_I) = \sum_{p \in \mathfrak{R}} \operatorname{Res}_{\zeta=p} K(z,\zeta) \,\mathcal{E}_{g,n}(\zeta,\iota(\zeta);z_I) - \frac{1}{2\pi \mathrm{i}} \oint_{\partial \Sigma} \omega_{0,2}(z,\zeta) \,\mathcal{V}_{g,n}(\zeta;z_I) \,,$$

where

$$K(z,\zeta) = \frac{1}{2} \frac{\int_{\iota(\zeta)}^{\zeta} \omega_{0,2}(z,\tau)}{\omega_{0,1}(\zeta) - \omega_{0,1}(\iota(\zeta))}$$

$$\mathcal{E}_{g,n}(z,\iota(z);z_{I}) = \omega_{g-1,\,n+1}(z,\iota(z),z_{I}) + \sum_{\substack{J \subseteq I , \ 0 \leqslant f \leqslant g \\ (J,f) \neq (\emptyset,0) , \ (I,g)}} \omega_{f,\,|J|+1}(z,z_{J}) \, \omega_{g-f,\,n-|J|}(\iota(z),z_{I \setminus J}) \, \omega_{g-f,\,n-|J|}(\iota(z),z_{I}) \, \omega_{g$$

## Schematic illustration of the Topological Recursion

• The operator  $\sum_{p \in \Re} \underset{\zeta = p}{\operatorname{Res}} K(z, \zeta)$  is represented by a pair of pants

 A differential ω<sub>g,n</sub> of degree n is represented by a surface of genus g with n boundary components



Image courtesy of Wikipedia

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