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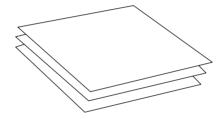
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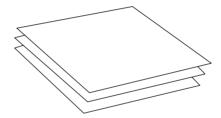


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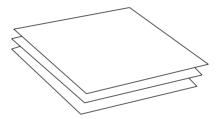
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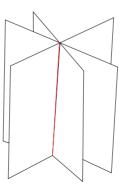
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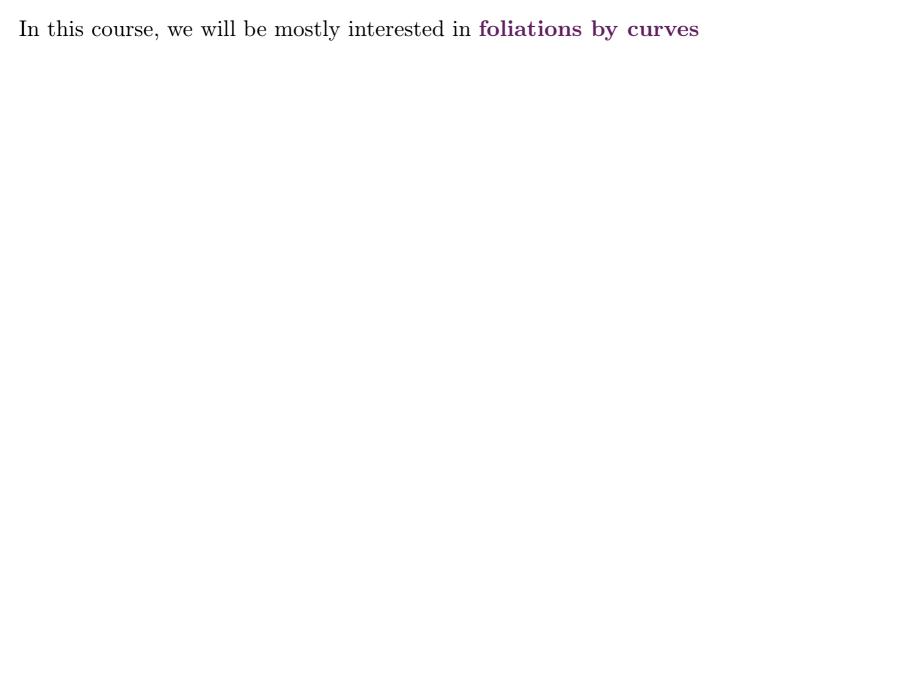
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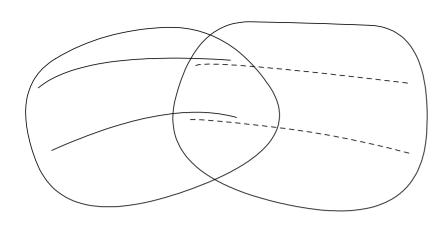
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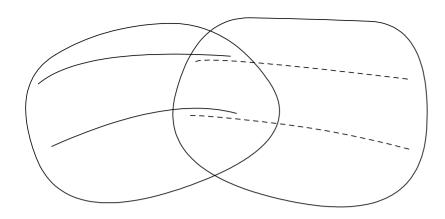
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Remark: In general, we cannot expect to have a single global generator for a foliation.

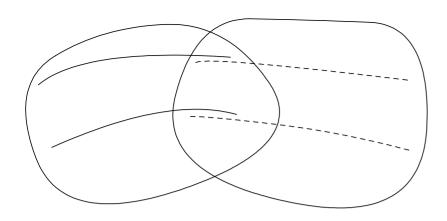


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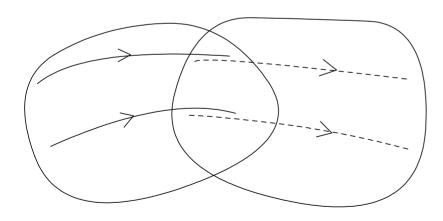
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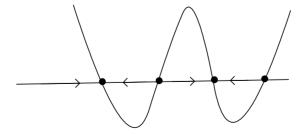
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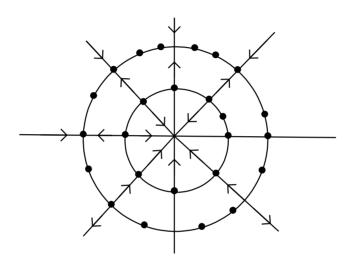
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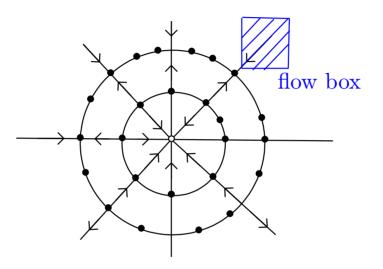


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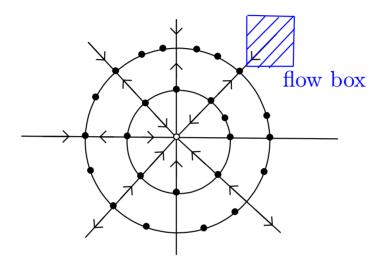
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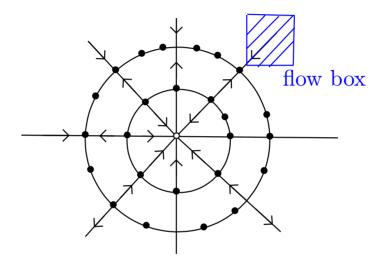


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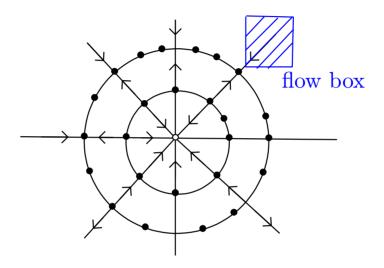
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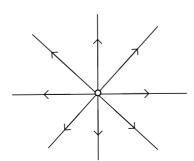
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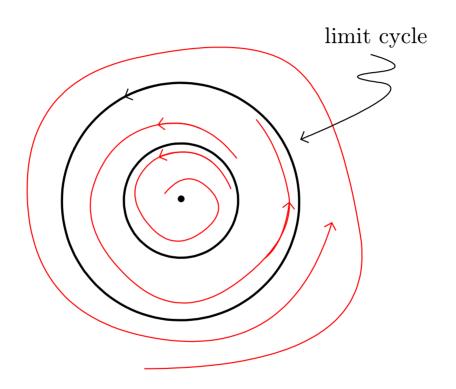


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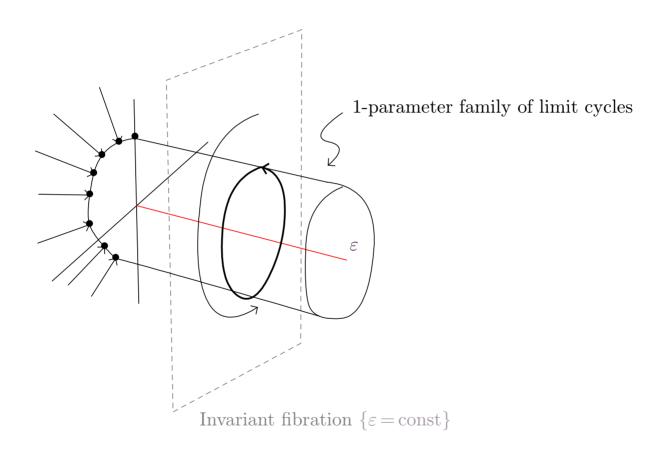
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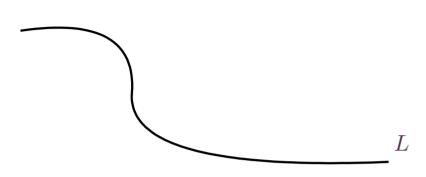
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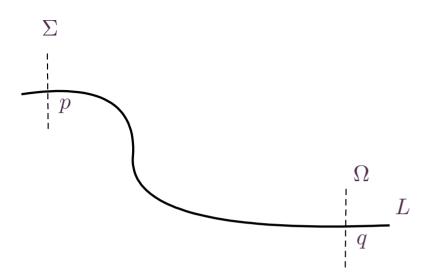
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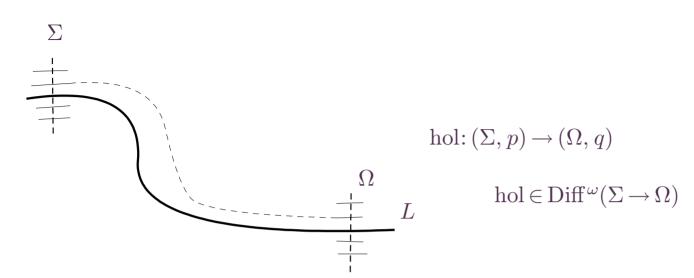
### **Holonomy Groupoid**



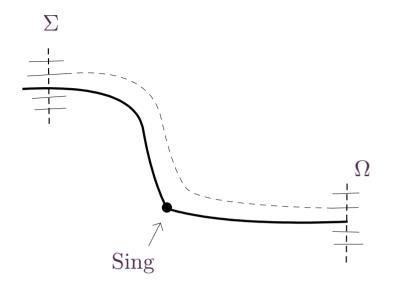


 $\begin{array}{c|c} \Sigma \\ \hline \downarrow \\ \hline \downarrow \\ \hline \end{array}$ 

any path  $p \rightarrow q$  on L can be lifted to nearby leafs

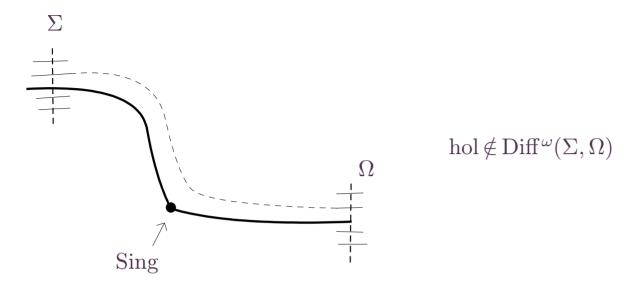


Adding a singularity on the path...



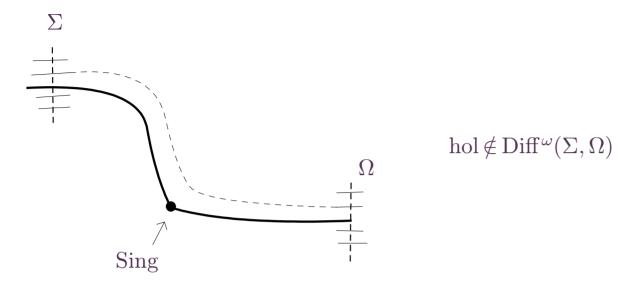
$$\operatorname{hol} \notin \operatorname{Diff}^{\omega}(\Sigma, \Omega)$$

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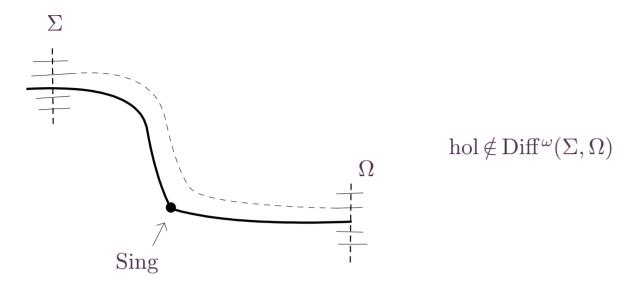
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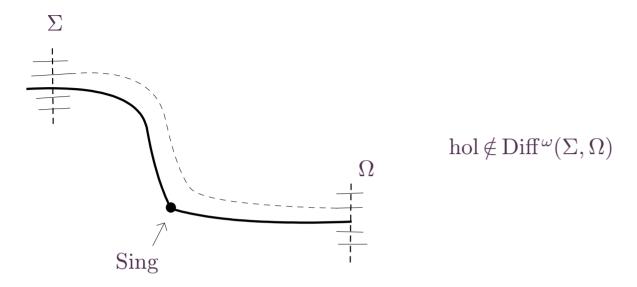


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(see the course of Patrick...)

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(i.e. that  $\partial \in \operatorname{End}_{\mathbb{C}}(\mathcal{O})$  stabilizes the maximal ideal)

**Flow-box Theorem** Then, there exists local analytic coordinates  $(f, g_1, ..., g_{n-1})$  such that

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Therefore  $f, \Phi(g_1), \ldots, \Phi(g_{n-1})$  is the required new coordinate system.

Then, (by Leibniz' rule)  $\partial(\mathbf{m}^{k+1}) \subset \mathbf{m}^{k+1}$  for each  $k \in \mathbb{N}$ , and  $\partial$  induces an sequence of endomorphism  $\{\partial_k\}_k$  on the jet spaces

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Moreover,  $\partial_s$  and  $\partial_n$  are derivations of  $\hat{\mathcal{O}} = \lim_{\longleftarrow} J^k$  (see Jean Martinet - Exposé Bourbaki'81).

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$$J^{k} \xrightarrow{O_{k}} J^{k}$$

$$\pi_{k,k-1} \downarrow \qquad \qquad \downarrow^{\pi_{k,k-1}}$$

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**Definition.** A germ of vector field  $\partial$  is *elementary* if:

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is generated (over  $\mathbb{C}$ ) by the monomials  $x^k = x_1^{k_1} \dots x_n^{k_n}$  such that  $\langle k, \lambda \rangle = \alpha$ .

The set of diagonal vector fields

$$L(\mu) = \sum_{i=1}^{n} \mu_i x_i \frac{\partial}{\partial x_i}, \qquad \mu \in \mathbb{C}^n$$

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where k ranges over the subset  $\mathbb{Z}^n \setminus \{0\}$  such that  $\langle \lambda, k \rangle = 0$ . These are the **resonant** monomials.

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The Poincaré-Dulac Theorem says that, up to a formal change of coordinates, we can write

$$\partial = \underbrace{\left(x\frac{\partial}{\partial x} - y\frac{\partial}{\partial y}\right)}_{\partial_s} + \underbrace{\sum_{k \geqslant 1} (xy)^k \left(a_k x\frac{\partial}{\partial x} + b_k y\frac{\partial}{\partial y}\right)}_{\partial_n}$$

where u = xy is the generator of the subring  $\ker(\partial_s)$ . By further reductions, we can write

$$(1+F)\left(\left(x\frac{\partial}{\partial x} - y\frac{\partial}{\partial y}\right) + \frac{u^n}{1+\rho u^n}\left(x\frac{\partial}{\partial x} + y\frac{\partial}{\partial y}\right)\right) \quad \text{or} \quad (1+F)\left(x\frac{\partial}{\partial x} - y\frac{\partial}{\partial y}\right)$$

for some  $F \in \mathbb{C}[[u]]$  of order  $\geqslant 1$ ,  $n \geqslant 1$  and  $\rho \in \mathbb{C}$ .

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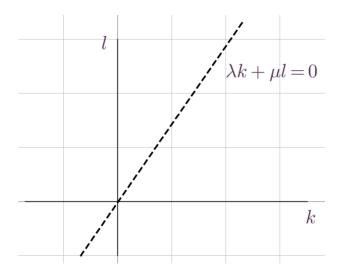
is a first integral of the vector field (namely,  $\partial I = 0$ ). It is an element of  $\mathbb{R}_{\mathrm{an,exp}}$ .

$$\partial = (\lambda x + \dots) \frac{\partial}{\partial x} - (\mu y + \dots) \frac{\partial}{\partial y}$$

Then,  $\operatorname{Spec}(\partial|_{J^1}) = \{\lambda, -\mu\}$ 

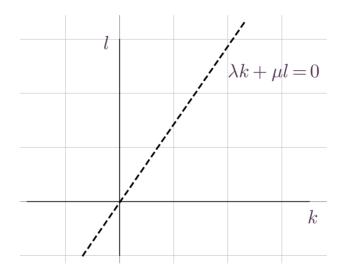
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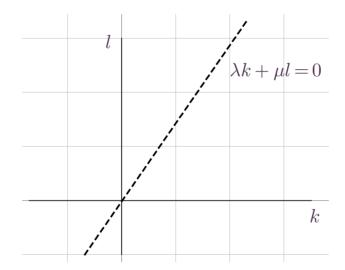


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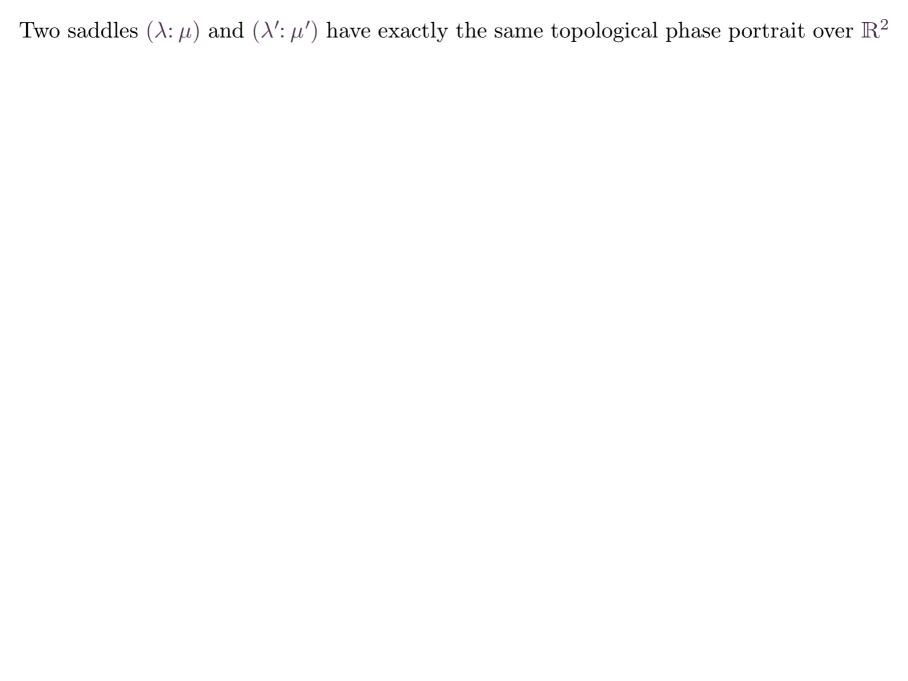
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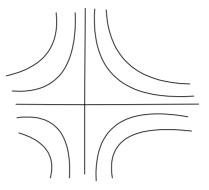
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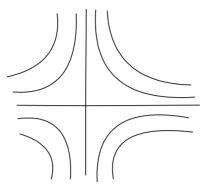
and the first integral is simply  $I = x^{\mu}y^{\lambda}$ .



Two saddles  $(\lambda:\mu)$  and  $(\lambda':\mu')$  have exactly the same topological phase portrait over  $\mathbb{R}^2$ 

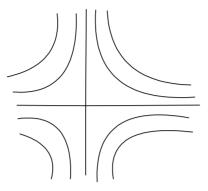


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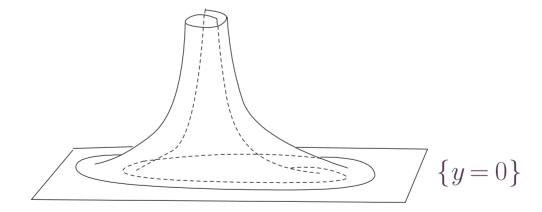


but they are completely different over  $\mathbb{C}^2$  for  $\lambda/\mu \neq \lambda'/\mu'$ .

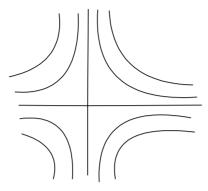
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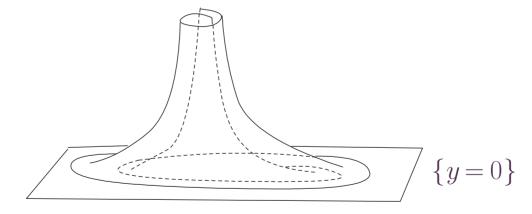
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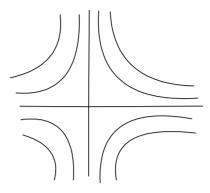


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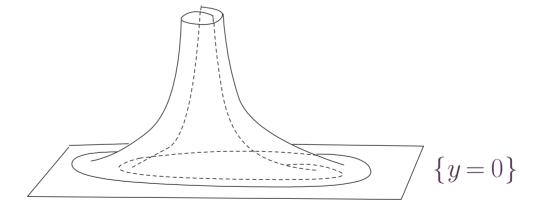


Over  $\mathbb{C}^2$ : There are several **rigidity phenomena** 

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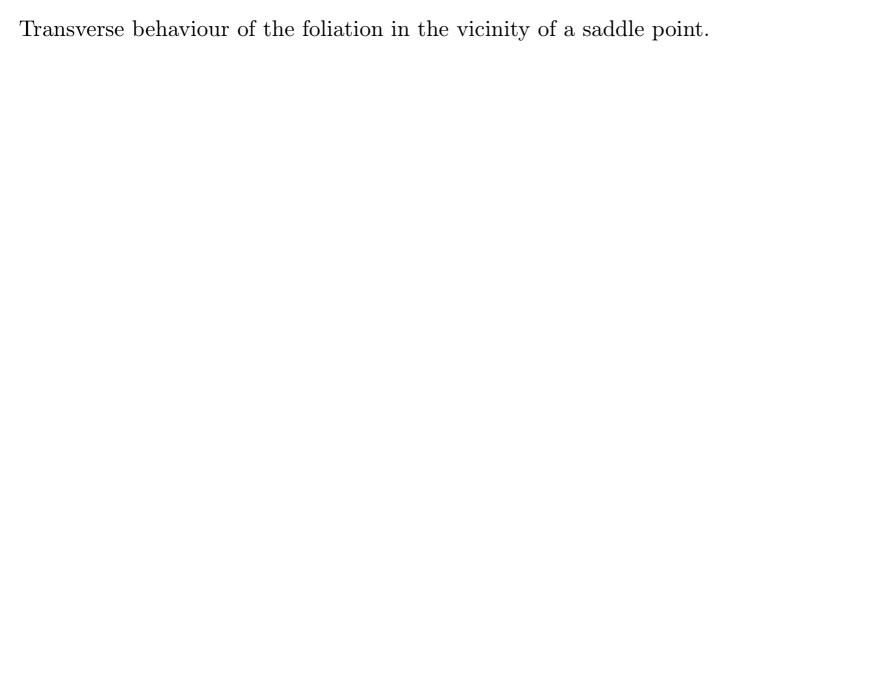


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Over  $\mathbb{C}^2$ : There are several **rigidity phenomena** 

E.g. Some analytic invariants are topologically determined (for instance, linearizability).

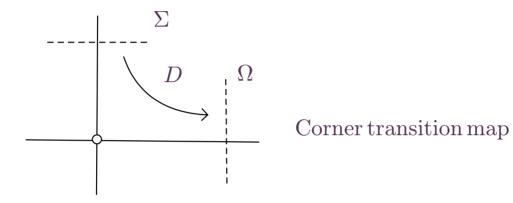


There are two holonomy maps of interest:

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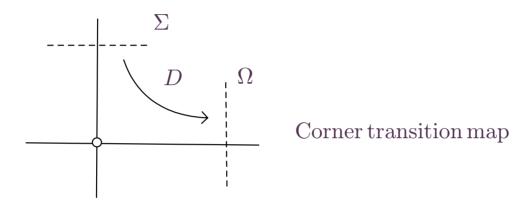
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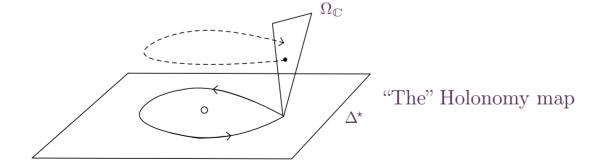
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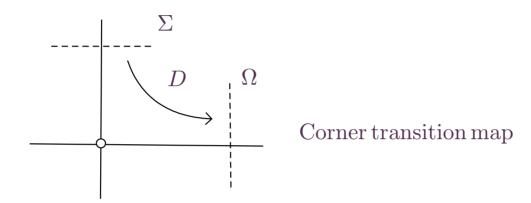


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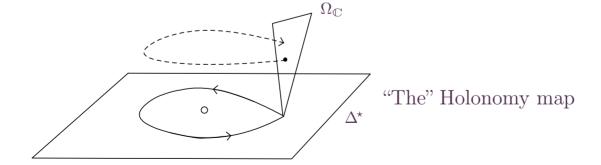


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2) In the complex setting...



We can recover the (orbital) analytic class of the saddle from the analytic class of one of these maps (once we fix the ratio  $\mu/\lambda$ )

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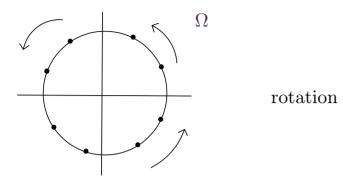
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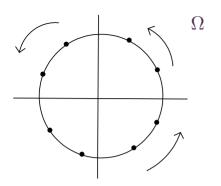
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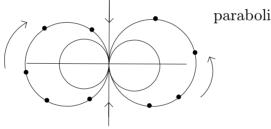
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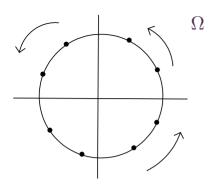


rotation

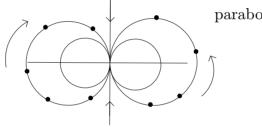


parabolic fixed point

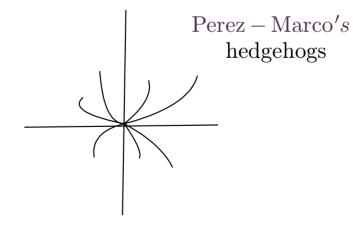
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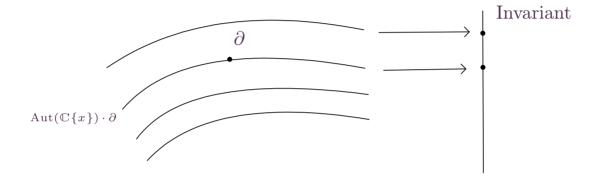
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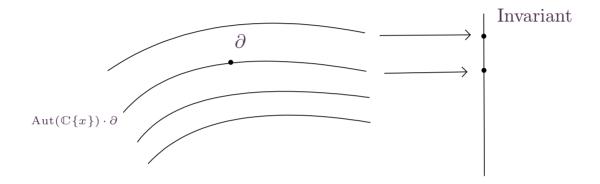
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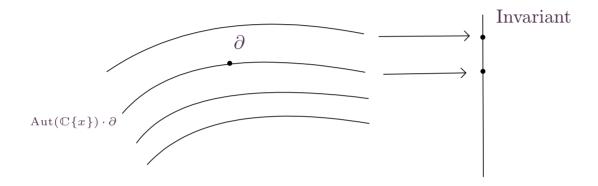
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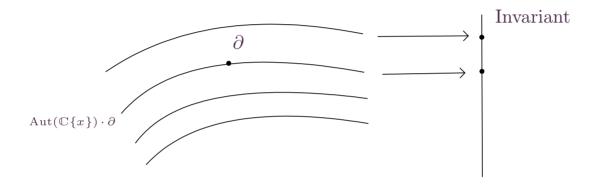


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The problem is reasonably well-understood for **elementary singularities in dimension two** (modulo some very hard *small divisor problems*) see e.g. Dulac, Ecalle, Ilyashenko, Martinet, Ramis, Yoccoz and Perez Marco, . . . works.

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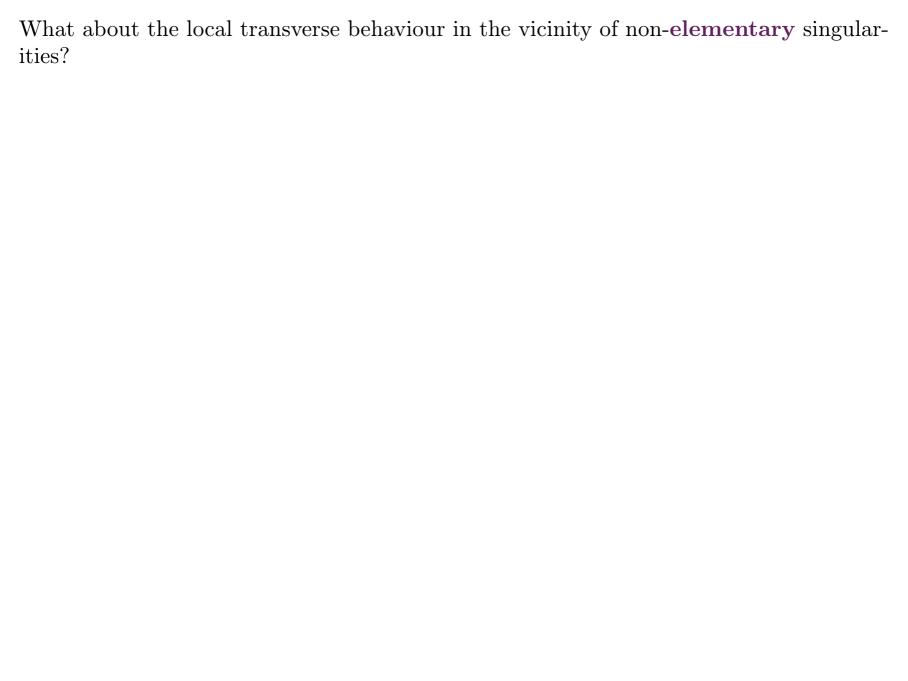
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This problem is much less understood for vector fields higher dimensions.



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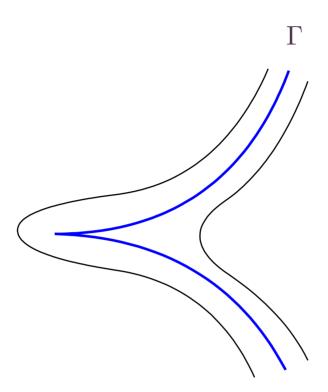
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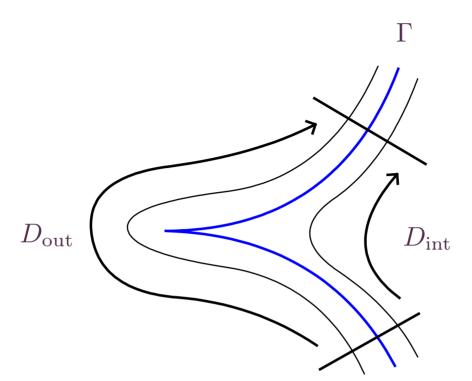
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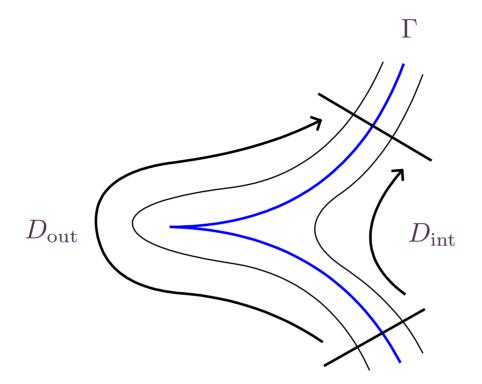
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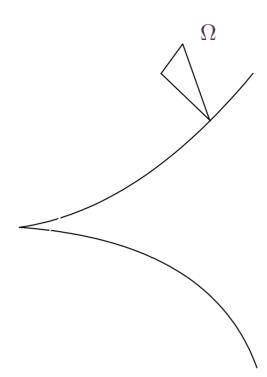
The cusp  $\Gamma = \{f = 0\}$  is an invariant curve.

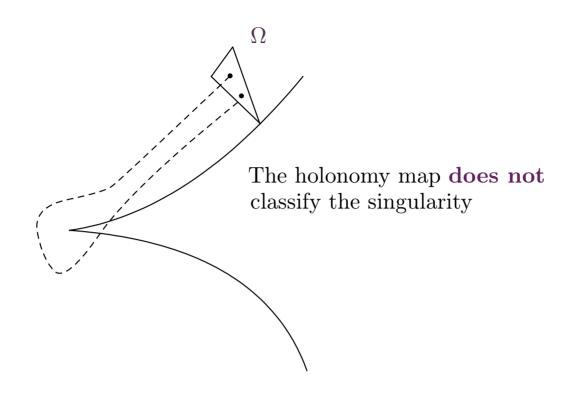


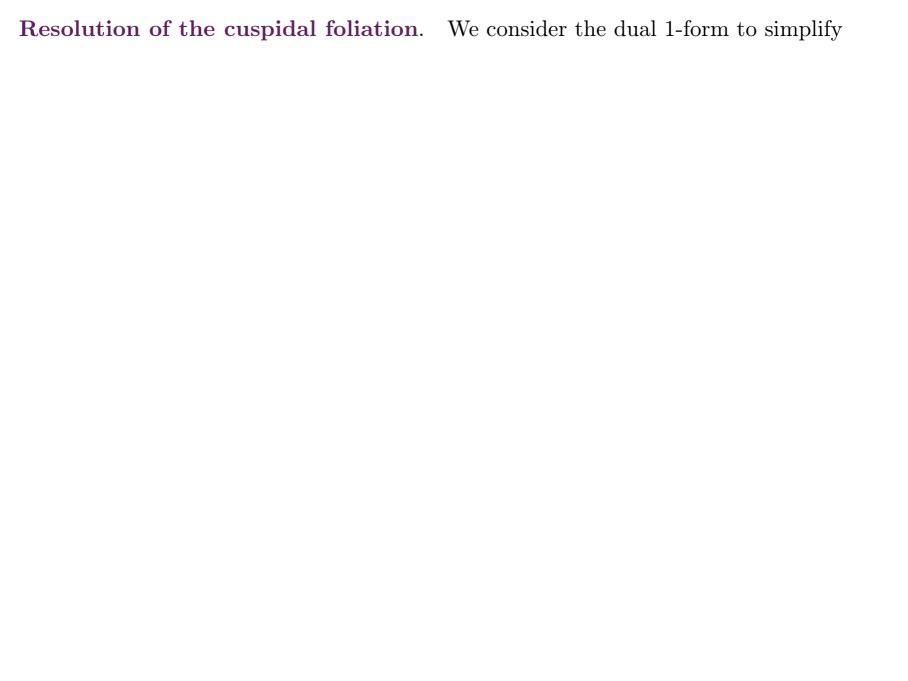


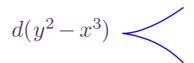


There are two **distinct** corner transition maps.









$$d(y^2 - x^3) <$$

Blow-up 1:  $x \rightarrow x$ ,  $y \rightarrow xy$ 

$$d(x^2(y^2-x))$$

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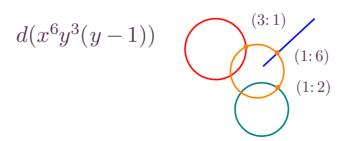
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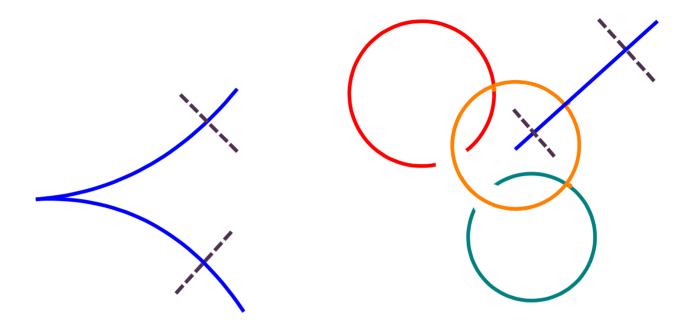
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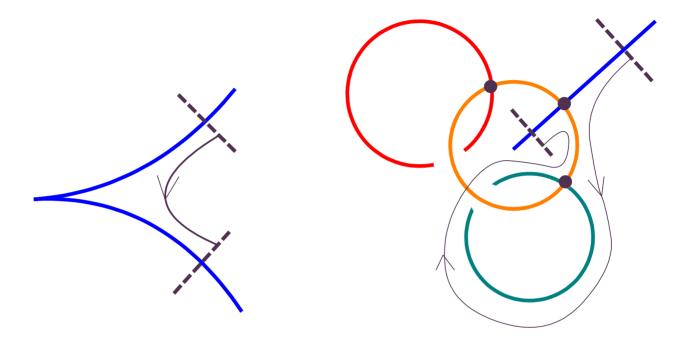
Blow-up 3:  $x \rightarrow x$ ,  $y \rightarrow xy$ 



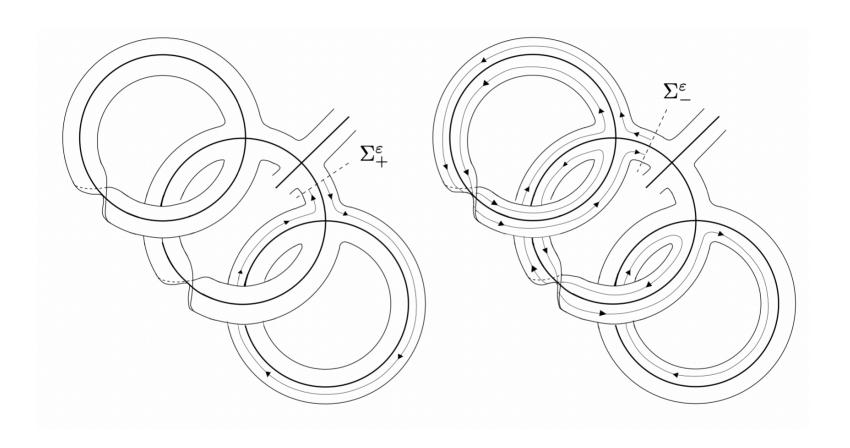
All singularities are now elementary saddles.



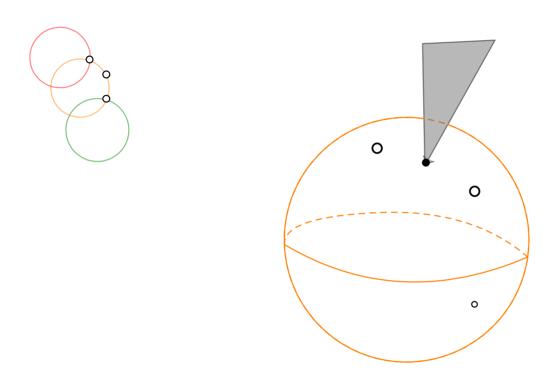
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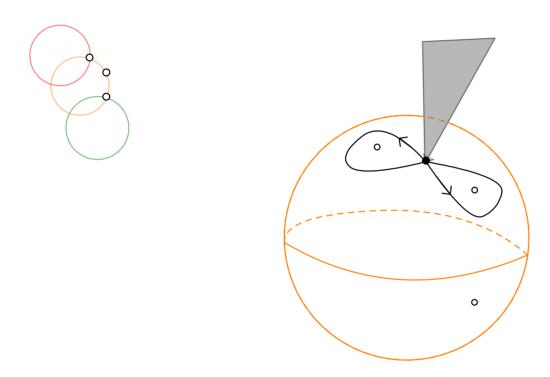
The foliation is now organized in a neighborhood of the exceptional divisor..



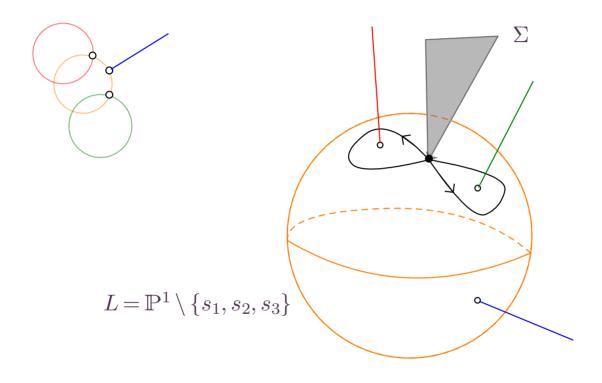
Can we recover the analytic moduli from the transverse behaviour?



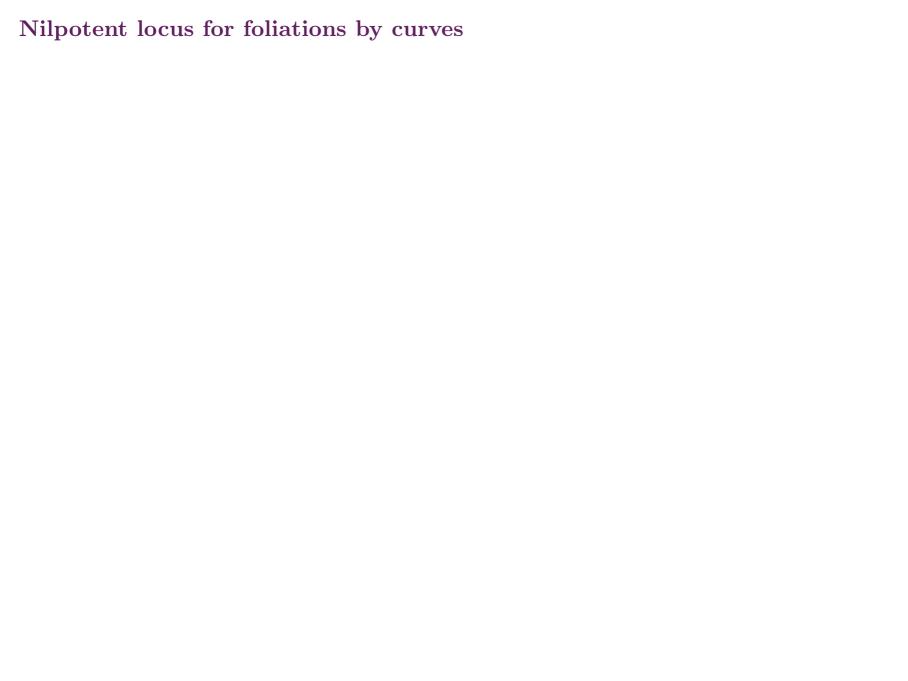
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(Moussu) The vanishing holonomy  $\operatorname{Hol}(\mathcal{F}, L) = \langle f, g \in \operatorname{Diff}(\mathbb{C}, 0) \mid f^2 = g^3 = \operatorname{id} \rangle$  characterizes the analytic class of the germ of foliation.



## Nilpotent locus for foliations by curves

The *nilpotent locus* of a foliated manifold is the subset  $\mathrm{Nilp}(M,\mathcal{F})$  of points where  $\mathcal{F}$  is not elementary.

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Claim:  $Nilp(M, \mathcal{F})$  is an analytic (or algebraic) subset of M.

(in fact,  $p \in \text{Nilp}(M, \mathcal{F}) \iff \partial(\mathbf{m}_p) \subset \mathbf{m}_p \text{ and } \partial_1 \in \text{End}_{\mathbb{C}}(\mathbf{m}_p/\mathbf{m}_p^2) \text{ is a nilpotent endomorphism, for } \partial \text{ some arbitrarily chosen local generator}).$ 

## Nilpotent locus for foliations by curves

The *nilpotent locus* of a foliated manifold is the subset  $\mathrm{Nilp}(M,\mathcal{F})$  of points where  $\mathcal{F}$  is not elementary.

Claim:  $Nilp(M, \mathcal{F})$  is an analytic (or algebraic) subset of M.

(in fact,  $p \in \text{Nilp}(M, \mathcal{F}) \iff \partial(\mathbf{m}_p) \subset \mathbf{m}_p \text{ and } \partial_1 \in \text{End}_{\mathbb{C}}(\mathbf{m}_p/\mathbf{m}_p^2) \text{ is a nilpotent endomorphism, for } \partial \text{ some arbitrarily chosen local generator}).$ 

Alternatively,

$$p \in \text{Nilp}(M, \mathcal{F}) \iff \forall k \in \mathbb{N} \exists n \in \mathbb{N} : (\partial_k)^n = 0$$

where  $\partial_k: J^k \to J^k$  is the induced derivation on the  $k^{\text{th}}$  jet.

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$$\forall i \in \mathbb{N} : \partial(\langle f^i \rangle) \subset \langle f^i \rangle$$

We further say that  $\mathcal{F}$  is **tightly adapted** to D if there exists an index i such that

$$\partial(\langle f^i \rangle) \not\subset \langle f^{i+1} \rangle$$

In other words, for  $E = (x_1 \dots x_k = 0)$ ,

$$\partial = \sum_{i=1}^{k} a_i \left( x_i \frac{\partial}{\partial x_i} \right) + \sum_{i=k+1}^{n} a_i \frac{\partial}{\partial x_i}$$

with  $a_1, \ldots, a_n \in \mathbb{C}\{x\}$  such that  $\langle a_1, \ldots, a_n \rangle \not\subset \langle x_i \rangle$ , for each  $i = 1, \ldots, k$ .

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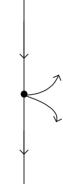
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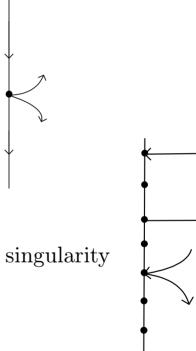
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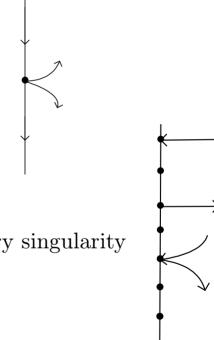


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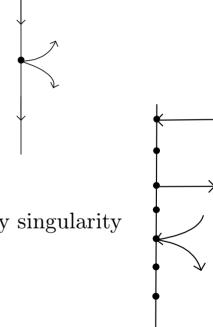


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 $\mathcal{F}$  is tightly adapted to  $E \iff$  no irreducible component of E lies on  $\mathrm{Nilp}(M,\mathcal{F})$ 

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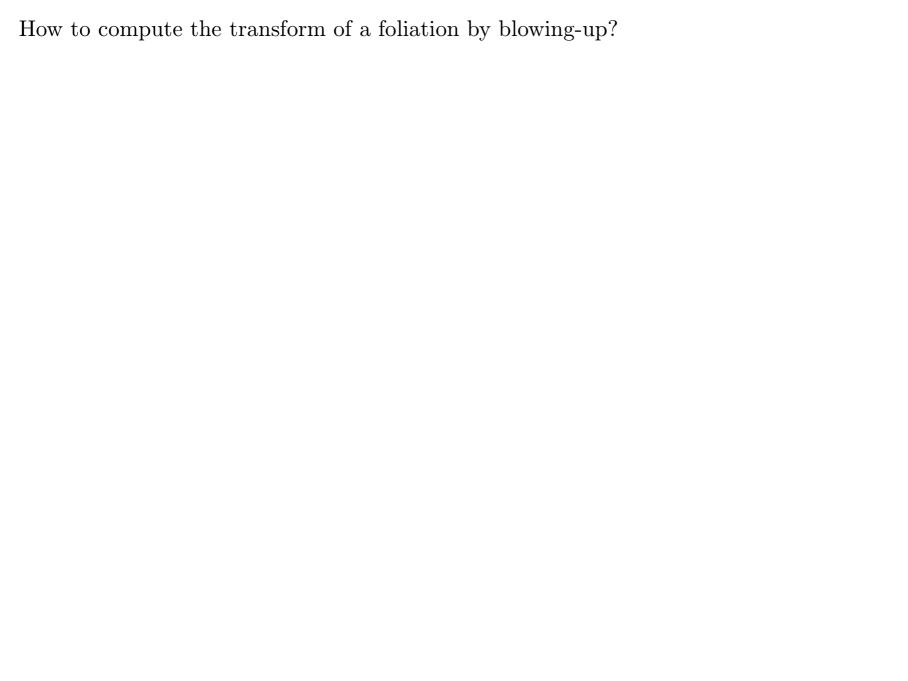
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- 2) Nilp $(M_n, \mathcal{F}_n) = \emptyset$ .



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It is easier to compute the strict transform of the **logarithmic basis**  $\left\{x_1 \frac{\partial}{\partial x_1}, \dots, x_n \frac{\partial}{\partial x_n}\right\}$ .

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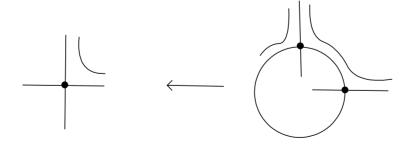
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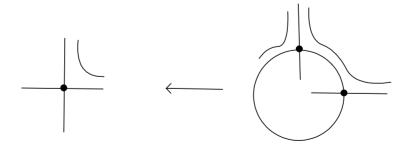
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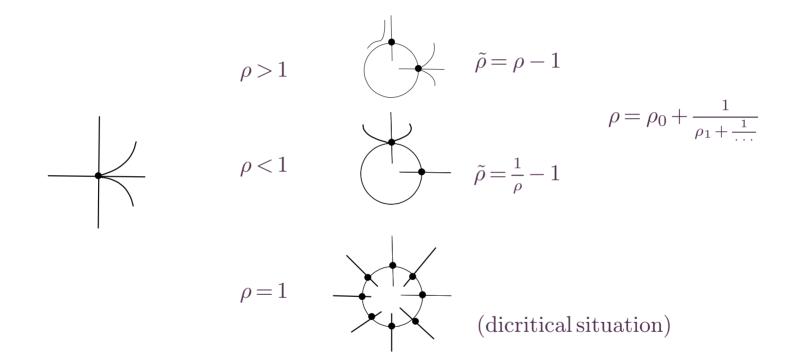


We can never get rid of saddle points...

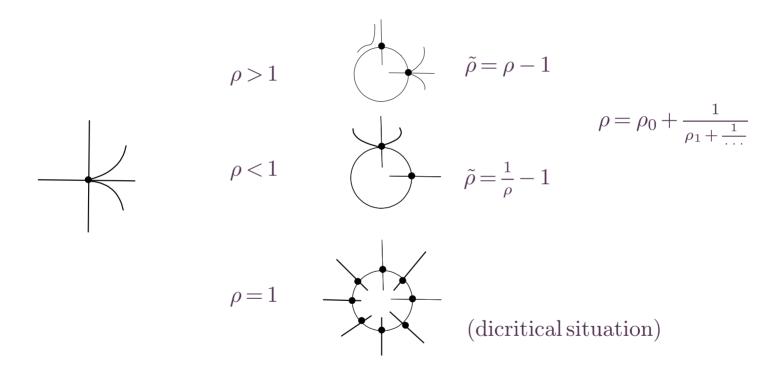
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We can never get rid of a node if  $\rho \notin \mathbb{Q}$ .

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This model is completely stable. It is a final model.

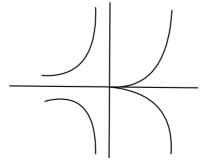
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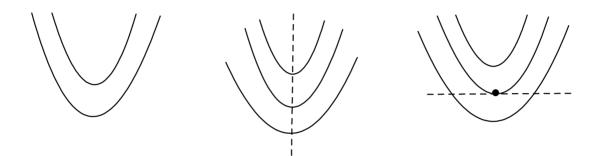
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First integral 
$$h = (x^m y) \exp\left(\frac{1}{kx^k}\right)$$

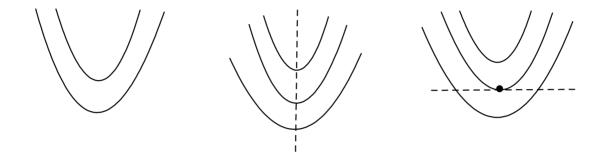


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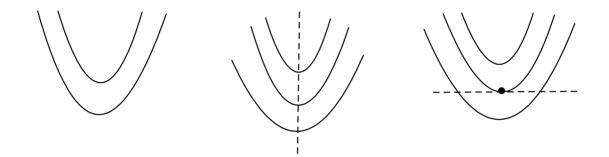


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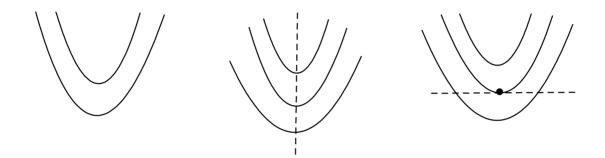
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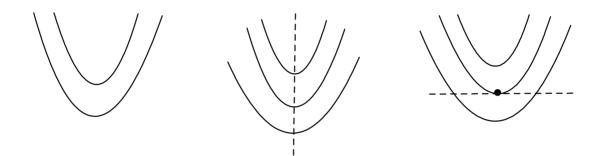


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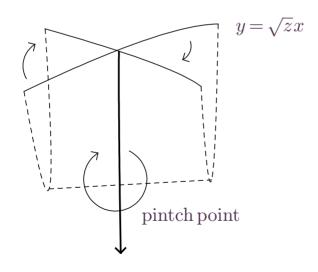
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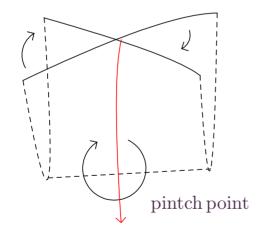


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with  $\beta \notin \frac{1}{2}\mathbb{Z}_{>0}$ ,  $\lambda \in \mathbb{C}^*$ .



Formal expansion of the "handle"

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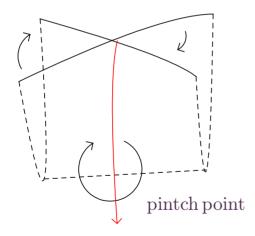
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We cannot take the handle as a blowing-up center because it is non-analytic.

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We consider the graph of the quotient mapping as a subset of  $\mathbb{C}^n \times \mathbb{P}^{n-1}_{\omega}$ 

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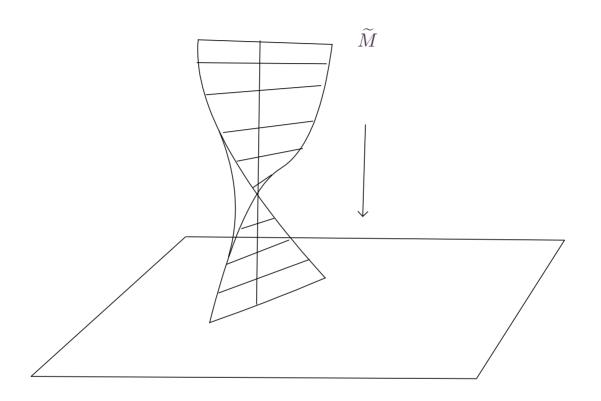
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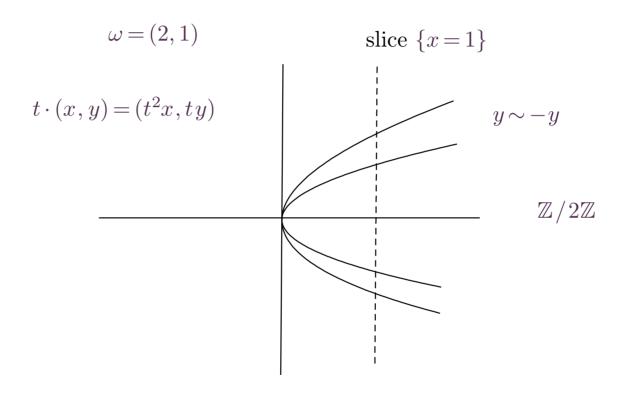
and the projection  $\pi: \widetilde{M} \to \mathbb{C}^n$  is the weighted blowing-up of the origin in  $\mathbb{C}^n$ .



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### Example



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We interpret  $(y_{1,..}, y_n)$  as an **orbifold chart** on  $\widetilde{M}$ . Namely the affine space  $\mathbb{C}^n$  equipped with an action of the cyclic group  $\mathbb{Z}/\omega_1\mathbb{Z}$ , defined by

$$y_1 \to \xi y_1$$
, For  $2 \leqslant k \leqslant n$ :  $y_k \longrightarrow \xi^{-\omega_k} y_k$ 

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The glueing of these charts equipps  $\widetilde{M}$  with the structure of an **orbifold**.

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An orbifold atlas on M is a collection  $\mathcal{U} = \{(U_i, G_i, \phi_i)\}_{i \in I}$  of pairwise compatible orbifold charts such that  $\{\phi(U_i)\}_{i \in I}$  forms an open cover of M.

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An **orbifold** is a pair  $(M, \mathcal{U})$  where M is paracompact Hausdorff topological space and  $\mathcal{U}$  is a maximal orbifold atlas on M.

A sub-variety  $Y \subset M$  is a **sub-orbifold** if for each point  $p \in Y$  there exists a local chart  $(U, G, \phi)$  such that  $\phi^{-1}(Y \cap U)$  is a G-invariant submanifold of U.

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 $X = \operatorname{Spec} \mathbb{C}[x, y]^G$  (ring of invariants)

$$\mathbb{C}[x,y]^G = \mathbb{C}[x^2, xy, y^2]$$

$$X = \operatorname{spec} \mathbb{C}[u, v, w] / (v^2 - uw)$$

X is the quadratic cone.

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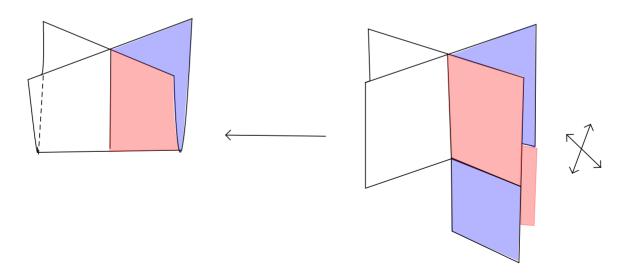
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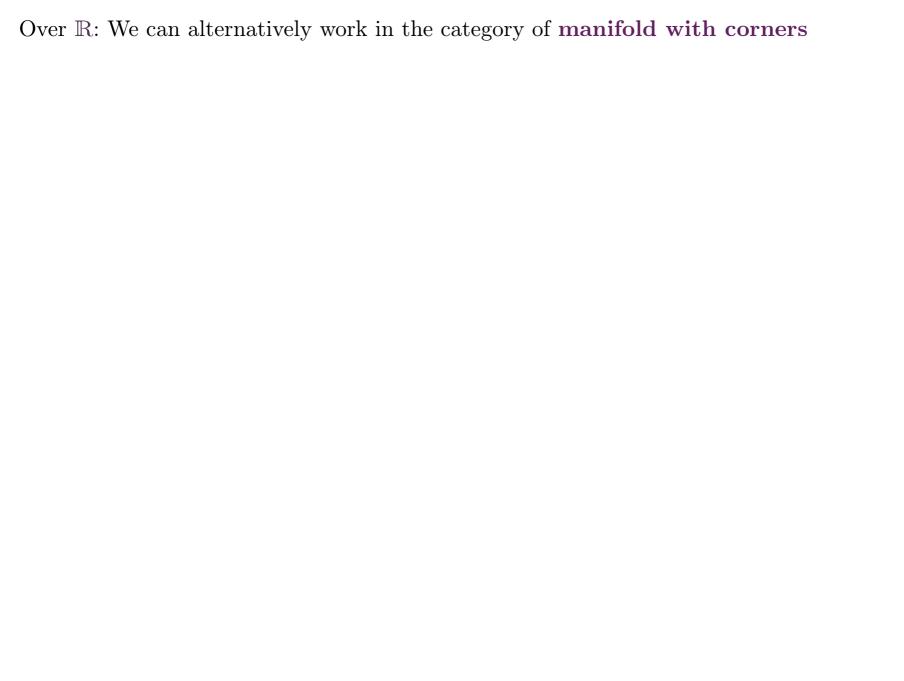
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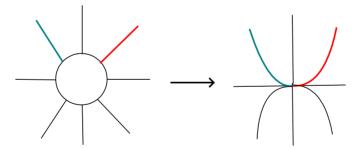
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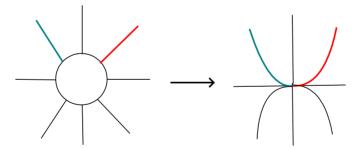


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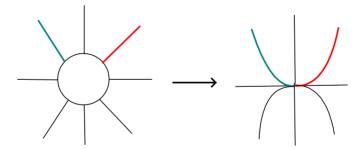
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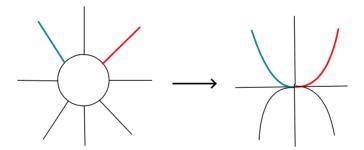
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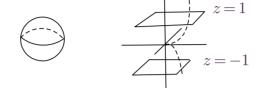
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(c.f. Melrose's "Analysis on manifolds with corners" - online)

Example: Spherical blowing-up of the (real) Whitney umbrella

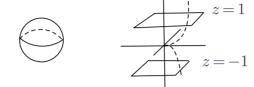
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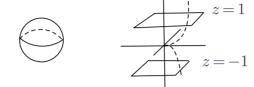
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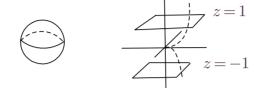
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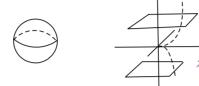
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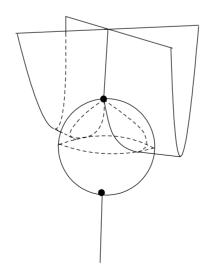
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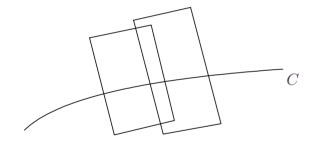
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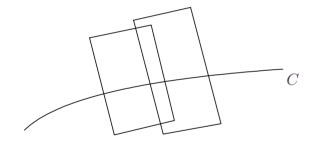


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$$\mathcal{O}_p = \mathcal{O}_0 \supset \mathcal{O}_1 \supset \mathcal{O}_2 \supset \cdots \qquad \mathcal{O}_k \cdot \mathcal{O}_l \subset \mathcal{O}_{k+l},$$

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such that  $F_i F_j \subset F_{i+j}$  and such that, for each point p on the support, the stalk of this filtration coincides with a quasi-homogeneous filtration as defined above.

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More generally, all automorphisms obtained by integrating the Lie algebra (over  $\mathbb{C}$ ) generated by

$$\left\{ x \frac{\partial}{\partial x}, y \frac{\partial}{\partial y}, x^{l} \frac{\partial}{\partial y}, y^{m} \frac{\partial}{\partial x} \mid m \geqslant 1, l \geqslant \beta \right\}$$

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The solution curves of  $\partial$  are precisely the orbits of the torus action  $t \cdot (x, y) = (tx, t^n y)$ .

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In the x-chart:  $x \to x^2$ ,  $y \to x^3y$ : (Using the assumption of the (2,3)-order of  $\Delta$ )

$$\partial = xy\left(x\frac{\partial}{\partial x} - 3y\frac{\partial}{\partial y}\right) + 3xy^{-1}\left(y\frac{\partial}{\partial y}\right) + x^2\Delta = x\left(xy\frac{\partial}{\partial x} + 3(1-y^2)\frac{\partial}{\partial y}\right) + x^2\Delta$$

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The divisor  $\{x=0\}$  is contained in the nilpotent locus. We factor out x and write

$$\partial_1 = xy \frac{\partial}{\partial x} + 3(1-y^2) \frac{\partial}{\partial y} + \Delta_1$$

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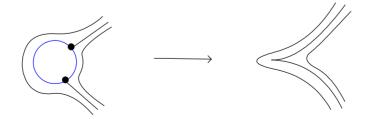
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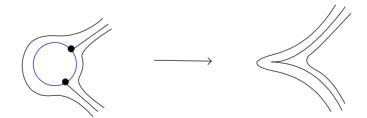
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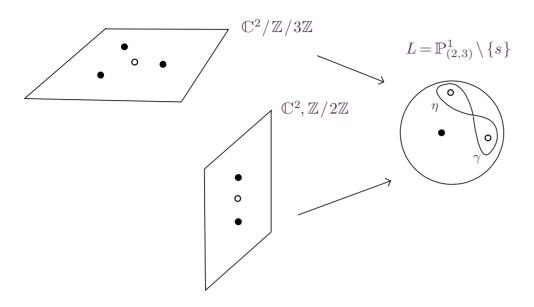
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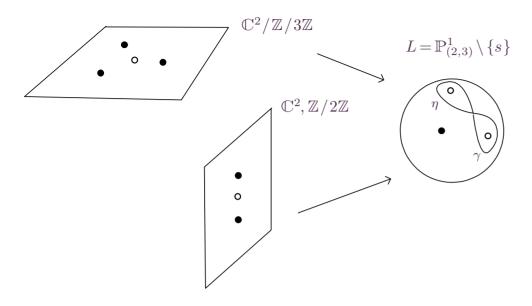
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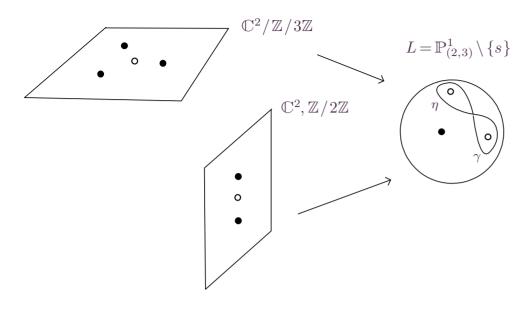
The resulting perturbation  $\Delta$  is of quadratic order along E (does not change the eingenvalues at the singular point)





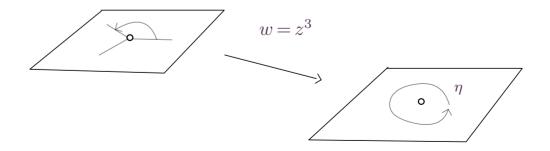
The fundamental group of the (orbi-)leaf L is

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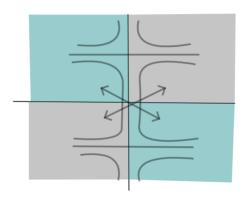


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$$g \cdot x = -x, \quad g \cdot y \to -y$$

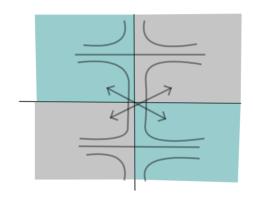
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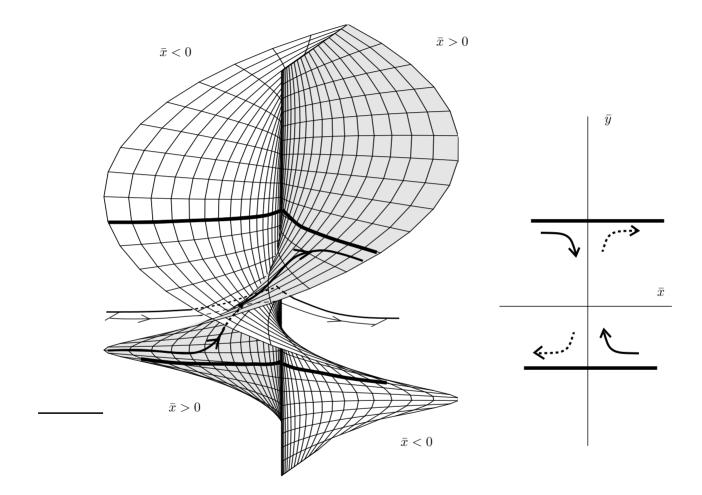


Other chart

$$\partial_2 = 2(1-x^3)\frac{\partial}{\partial x} - x^2y\frac{\partial}{\partial y}$$

$$g \cdot x = \xi^{-2} x$$
,  $g \cdot y = \xi y$ ,  $(\xi^3 = id)$ 

$$g \cdot \partial_2 = \xi^2 \partial_2$$



Van der Essen's proof (c.f. Ilyashenko-Yakovenko's book) We write  $\partial = a \frac{\partial}{\partial x} + b \frac{\partial}{\partial y}$ 

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Suppose that the germ is singular. We can assume that  $a,b \in \mathbb{C}\{x,y\}$  have no common factor and consider

$$m(0) = \dim_{\mathbb{C}} \frac{\mathbb{C}\{x, y\}}{(a, b)} \ge 1, \quad \mu(0) = \min_{k} \{(J^{k}a, J^{k}b) \ne (0, 0)\}$$

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- If  $l(0) \geqslant 2$  then  $m(\tilde{p}_i) < m(p)$
- If l(0) = 1 then this is a special case which has to be treated separately...

Example of "special case".

$$y\frac{\partial}{\partial x} + x^M \frac{\partial}{\partial y}$$

$$\mu = 1, m = M \geqslant 3$$

$$x \to x$$
,  $y \to xy$ 

$$xy\frac{\partial}{\partial x} + (x^{M-1} - y^2)\frac{\partial}{\partial y}$$

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The "invariant" increases and this case needs to be treates separately...

Example of "special case".

$$y\frac{\partial}{\partial x} + x^M \frac{\partial}{\partial y}$$

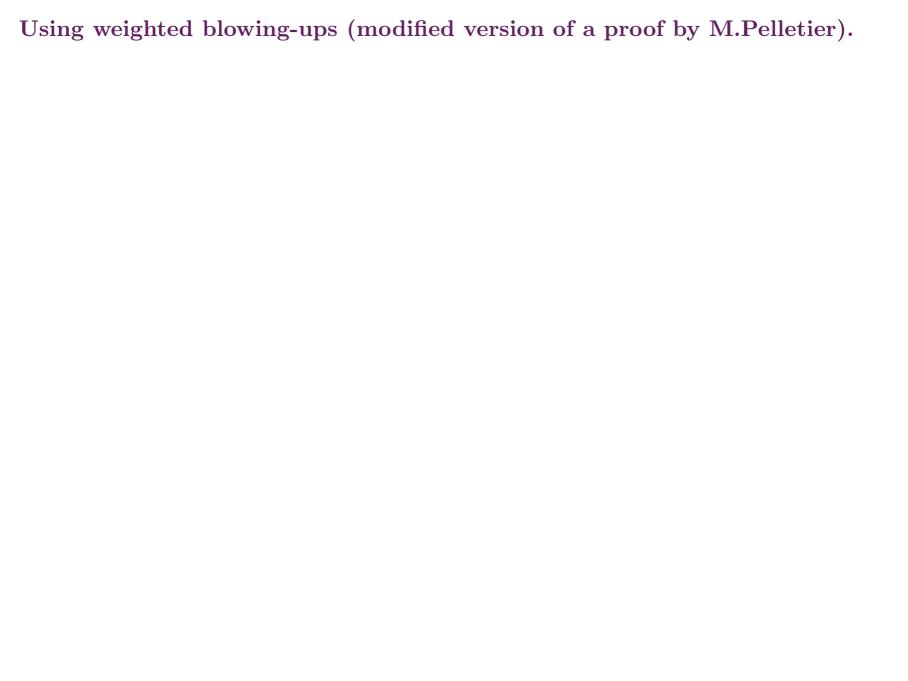
$$\mu = 1, m = M \geqslant 3$$

 $x \to x$ ,  $y \to xy$ 

$$xy\frac{\partial}{\partial x} + (x^{M-1} - y^2)\frac{\partial}{\partial y}$$

$$\mu = 2, m = M + 1$$

The "invariant" increases and this case needs to be treates separately...



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which will define the blowing-up...

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# Intermezzo: The Newton polyhedron of a germ of vector field

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We can reorder the expansion and write the monomial expansion

$$\partial = \sum_{k \in \mathbb{Z}^n} x^k L(\mu_k)$$

where, we recall, each  $L(\mu) = \sum \mu_i x_i \frac{\partial}{\partial x_i}$  is a diagonal vector field, i.e. an element of the  $\mathbb{C}$ -maximal toral subalgebra

$$\mathfrak{t} = \left\langle x_1 \frac{\partial}{\partial x_1}, \dots, x_n \frac{\partial}{\partial x_n} \right\rangle$$

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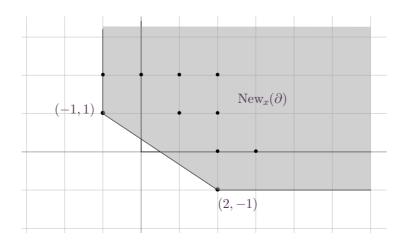
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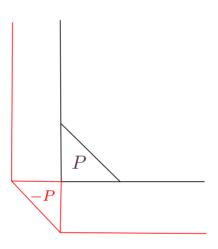
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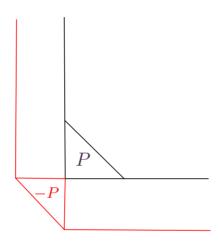
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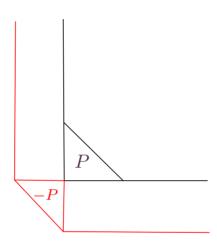


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$$\operatorname{supp}_{x}(\partial) \subset \{k: k_{i} \geqslant 0\} \quad \land \quad \operatorname{supp}_{x}(\partial) \cap \{k: k_{i} = 0\} \neq \emptyset$$

 $(x=0) \text{ invariant} \Longleftrightarrow \partial(\langle x \rangle) \subset \langle x \rangle \Longleftrightarrow a \in \mathbb{C}\{x,y\} \Longleftrightarrow [(k,l) \in \operatorname{supp}(\partial) \Longrightarrow k \geqslant 0]$ 

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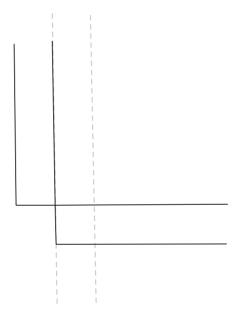
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The resolution of singularities should correspond to a combinatorial game based on the Newton polyhedron.

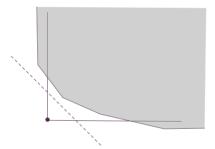
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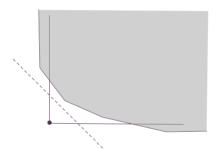
**Proposition:**  $\partial \in \text{Der}(\mathcal{O})$  is a nilpotent germ if and only if there exists a local system of coordinates  $x = (x_1, \dots, x_n)$  such that  $0 \notin \text{New}_x(\partial)$ .



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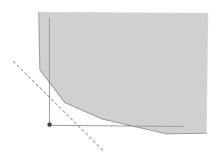


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$$\operatorname{New}_x(\partial) \subset H = \{ \langle \omega, \cdot \rangle \geqslant \alpha \}$$

(indeed, if some  $\omega_i < 0$  then for  $v \in \operatorname{supp}_x(\partial)$ ,  $\langle \omega, v + t e_i \rangle \to -\infty$  as  $t \to +\infty$ ).

$$\lambda(t) \cdot x_i = t^{\omega_i} x_i, \qquad i = 1, \dots, n$$

(or, equivalently, the graduation associted to the infinitesimal semisimple generator  $\delta = \sum \omega_i x_i \frac{\partial}{\partial x_i}$ ). This action is diagonalizable and we have a direct sum decomposition

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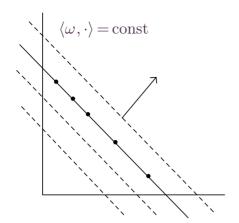
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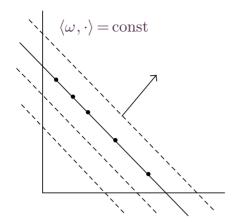
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And, naturally  $\partial \in \operatorname{Gr}_{\alpha}$ ,  $f \in \operatorname{Gr}_{\beta} \Longrightarrow \partial f \in \operatorname{Gr}_{\alpha+\beta}$ .

$$\partial \in \operatorname{Gr}_{\alpha}(\operatorname{Der}, \lambda) \Longleftrightarrow \operatorname{supp}_{x}(\partial) \subset \{k: \langle \omega, k \rangle = \alpha\}$$



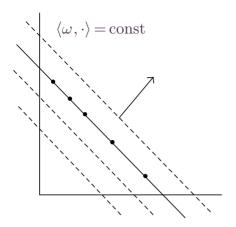
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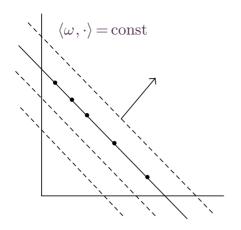


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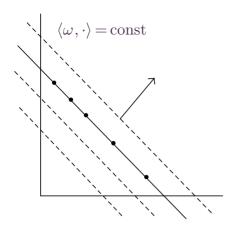


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As a consequence, for  $\mathbf{m} = \langle x_1, \dots, x_m \rangle$  the maximal ideal, for each s there exists a  $r \geqslant 1$  such that

$$\partial^r(\mathfrak{m}^s) \subset \mathfrak{m}^{s+1}$$

(because for  $k \in \mathbb{Z}_{\geq 0}^n$ ,  $|k| \geq \langle \omega, k \rangle / \max{\{\omega_i\}}$ ). Hence,  $\partial$  is nilpotent.

Reciprocally, assume that  $\partial$  is nilpotent. Then,  $\partial(\mathbf{m}) \subset \mathbf{m}$  and  $\partial_S = 0$ . There exists a local coordinate system such that  $\partial|_{J^1} = \begin{pmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$ , i.e. such that

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where  $\varepsilon_i \in \{0, 1\}$ . In other words, in the logarithmic basis, we obtain

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We now consider the weight-vector  $\rho = (-n/2, \dots, n/2)$ , or any other rational vector satisfying.

$$\langle \mathbf{1}, \rho \rangle = 0, \qquad \langle \rho, e_{i+1} - e_i \rangle > 0, \qquad e_i = (0, \dots, 1, \dots 0)$$

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We now consider the weight-vector  $\rho = (-n/2, \dots, n/2)$ , or any other rational vector satisfying.

$$\langle \mathbf{1}, \rho \rangle = 0, \qquad \langle \rho, e_{i+1} - e_i \rangle > 0, \qquad e_i = (0, \dots, 1, \dots 0)$$

Then, for all sufficiently small  $\varepsilon \in \mathbb{Q}_{>0}$ , the semi-simple derivation  $\omega = h + \varepsilon L(\rho)$  defines a half-space which separates  $\mathrm{New}_x(\partial)$  from 0.

Reciprocally, assume that  $\partial$  is nilpotent. Then,  $\partial(\mathbf{m}) \subset \mathbf{m}$  and  $\partial_S = 0$ . There exists a local coordinate system such that  $\partial|_{J^1} = \begin{pmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$ , i.e. such that

$$\partial(x_i) = \varepsilon_i x_{i+1} \qquad (\operatorname{mod} \mathfrak{m}^2)$$

where  $\varepsilon_i \in \{0, 1\}$ . In other words, in the logarithmic basis, we obtain

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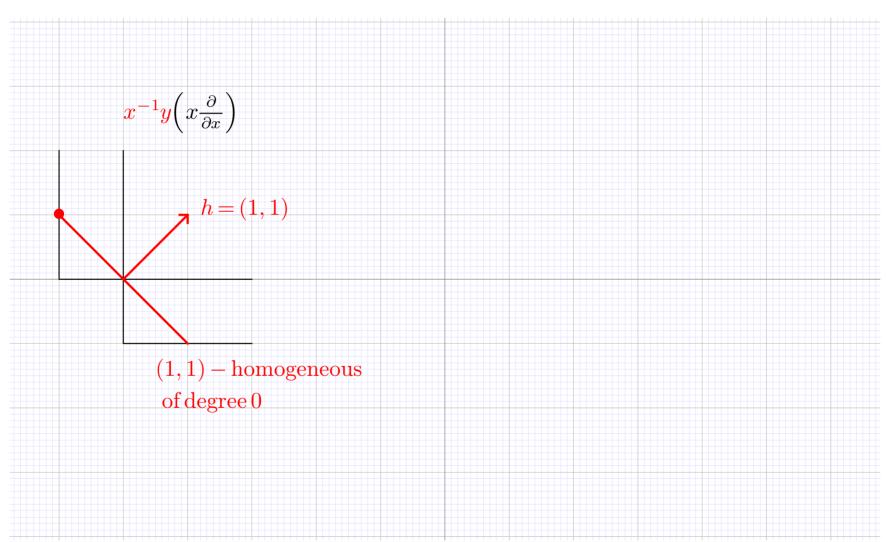
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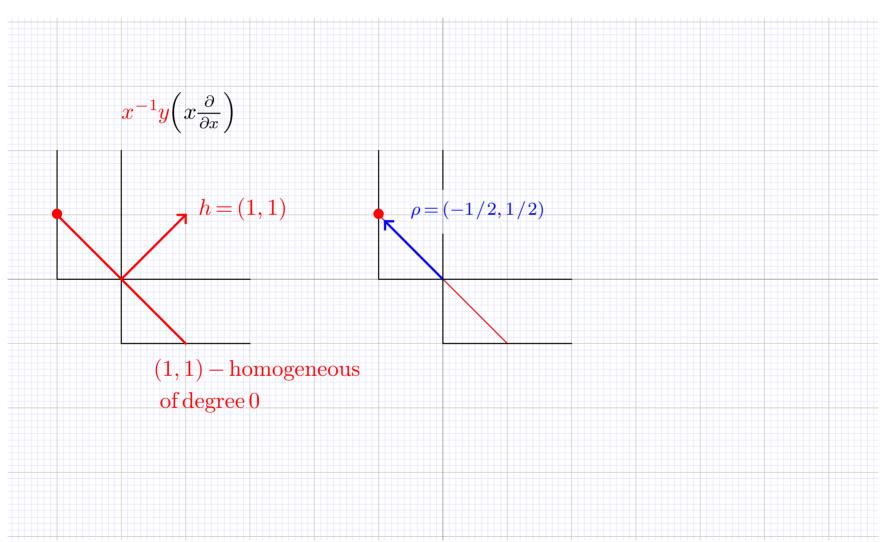
(because for  $|k| \ge 2$ ,  $\langle \omega, k \rangle \ge 2 - n\varepsilon |k|$ , and New<sub>x</sub>( $\partial$ ) has finitely many vertices)



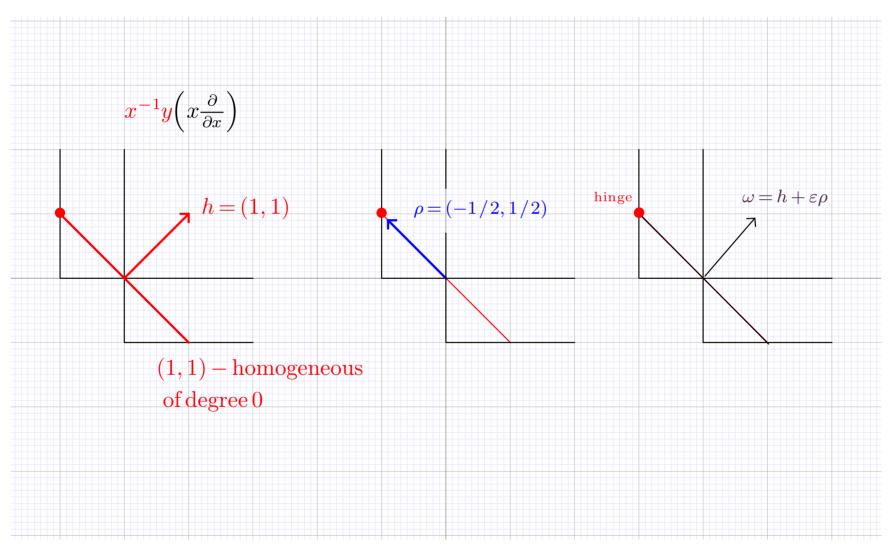




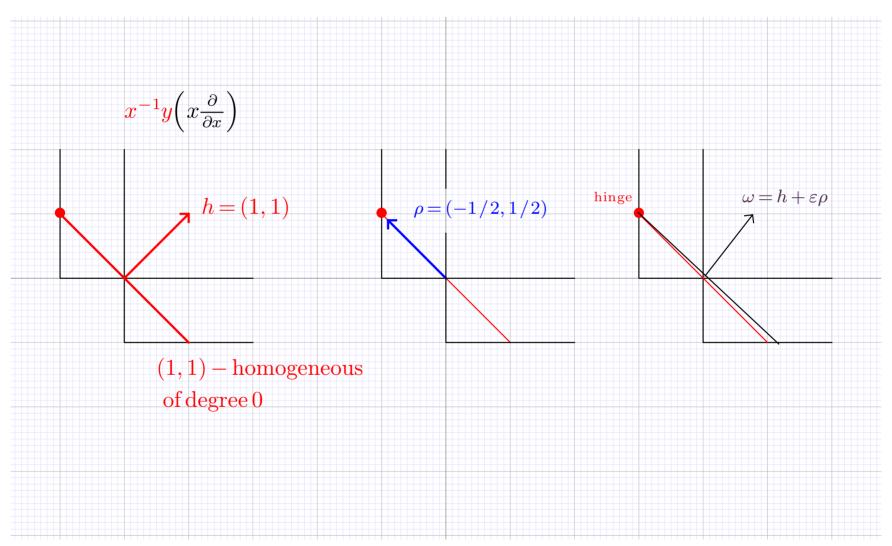




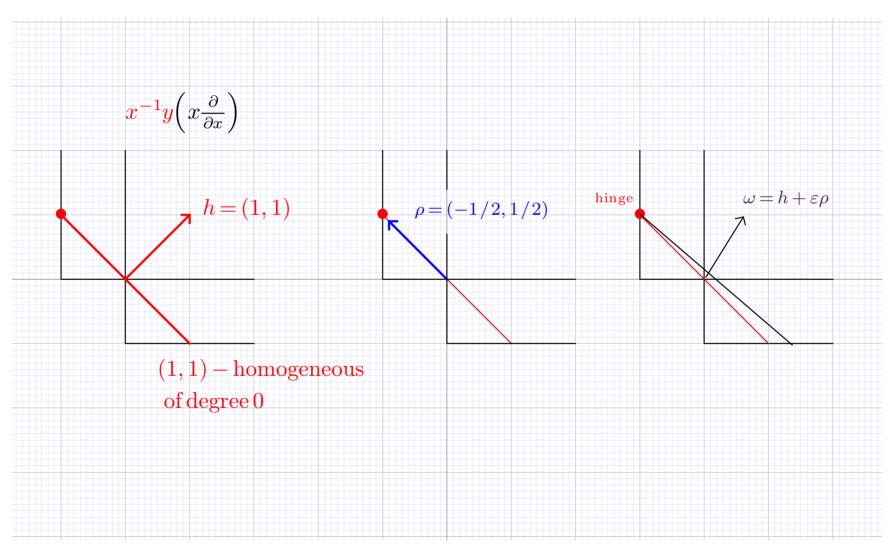




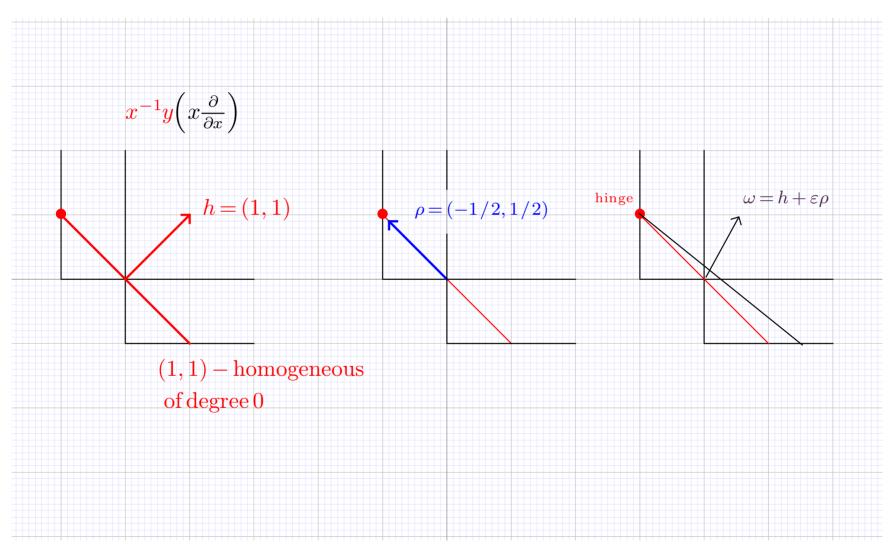




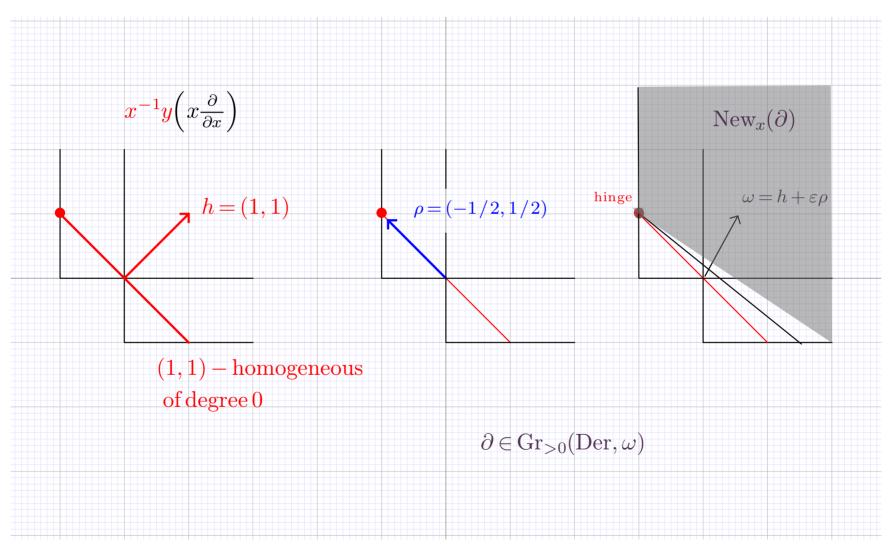












Claim: Suppose that  $\partial$  is elementary (i.e. not-nilpotent). Then, for all choices of coordinate systems  $(x_1, \ldots, x_n)$ ,

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The case  $\partial(\mathfrak{m}) \not\subset \mathfrak{m}$  is even easier. In fact,  $\partial(\mathfrak{m}) \not\subset \mathfrak{m}$  if and only if

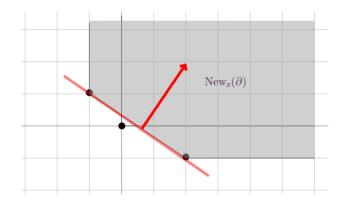
$$\exists i \in \{1, \dots, n\}: -e_i = (0, \dots, -1, \dots, 0) \in \text{New}_x(\partial)$$

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is such that  $\partial \in \operatorname{Gr}_{\geqslant 1}$ . (write  $x^{-1}y\left(x\frac{\partial}{\partial x}\right) + x^2y^{-1}\left(y\frac{\partial}{\partial y}\right)$  and  $x^{-1}y, x^2y^{-1} \in \operatorname{Gr}_1$ )

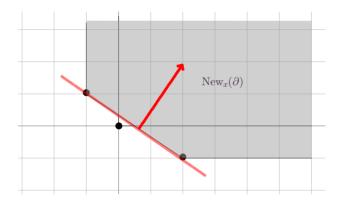
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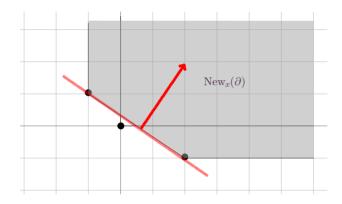
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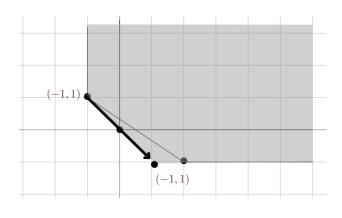
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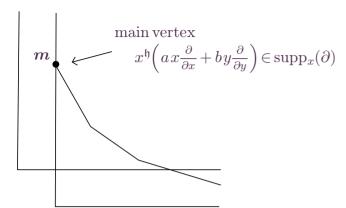
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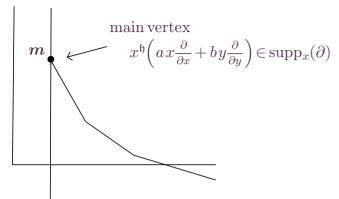
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To simplify, we will assume that e(p) = 1 for all points  $p \in \text{Nilp}(M, \mathcal{F})$ .

(otherwise it suffices to slightly modify the invariant by including e(p) lexicographically).

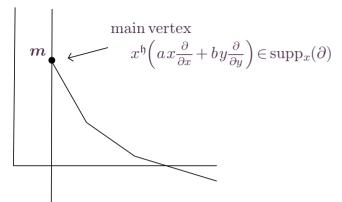


Invariance of (x = 0) implies that  $\partial \in \operatorname{Gr}_{\geqslant 0}(\cdot, x \frac{\partial}{\partial x})$  (i.e.  $\partial(\langle x \rangle) \subset \langle x \rangle$ )



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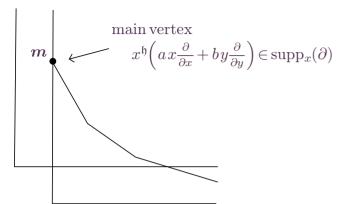
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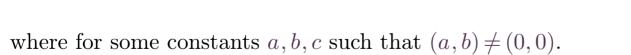


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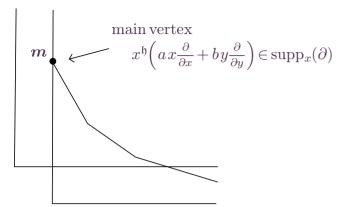
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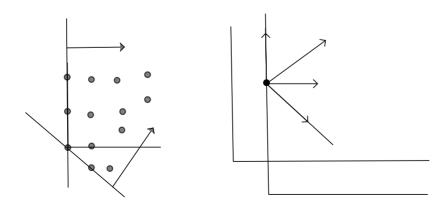
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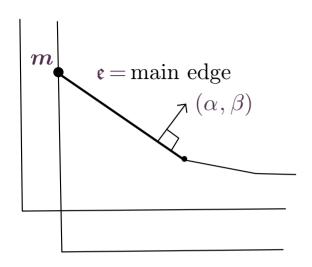
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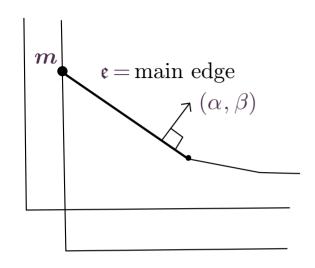
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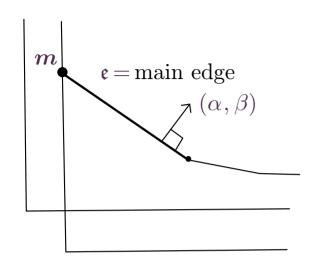






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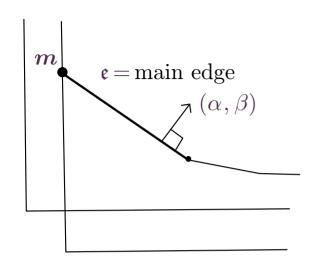
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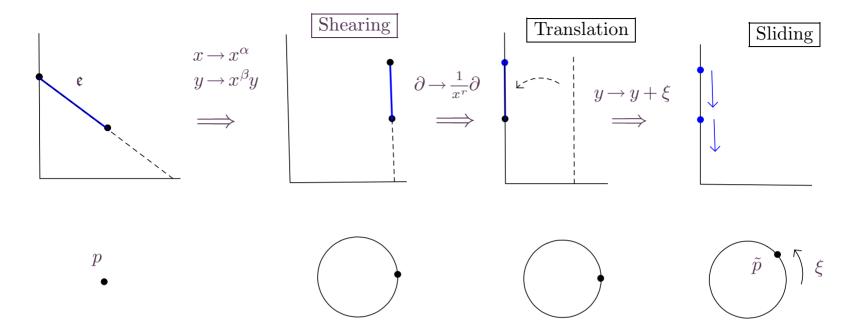


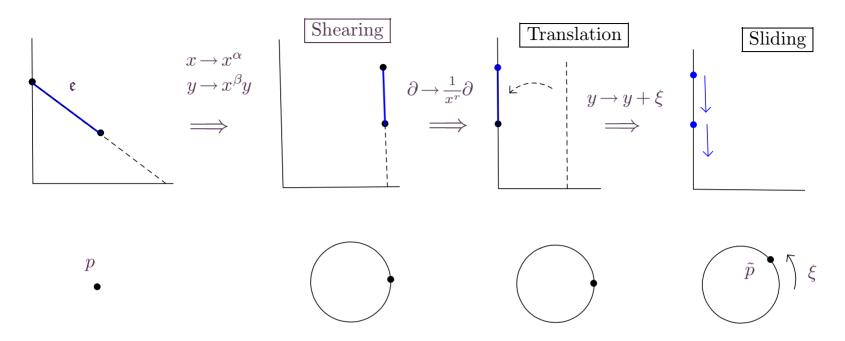
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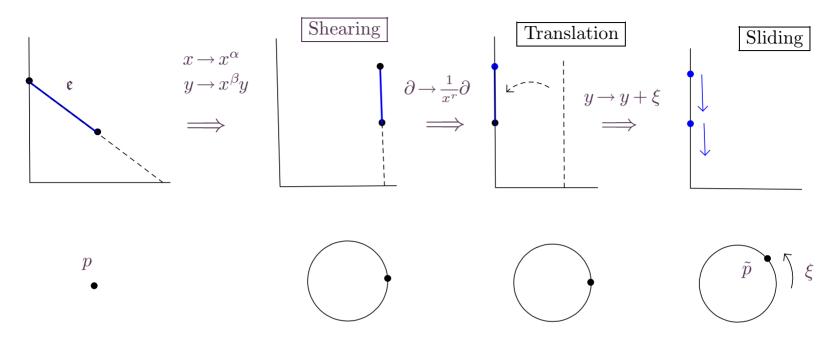
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Let  $\operatorname{wt}(\mathfrak{e}) = \alpha x \frac{\partial}{\partial x} + \beta y \frac{\partial}{\partial x}$  denote the irreducible weight-vector determined by  $\mathfrak{e}$ .

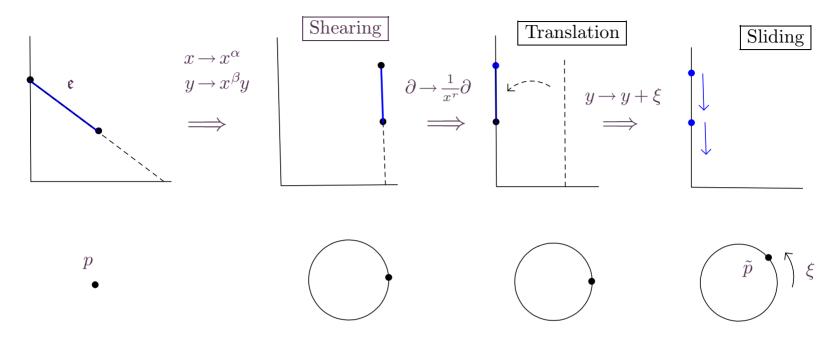




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We can have a full compensation phenomena in the "sliding phase".

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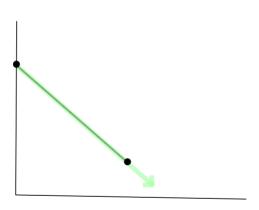
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The above map slides the monomials in the direction of the main edge.

**Theorem (Local resolution)** Suppose that  $\text{New}_{(x,y)}(\partial)$  is edge stable, and let

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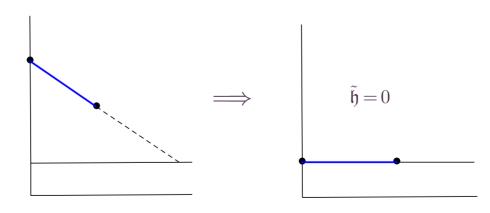
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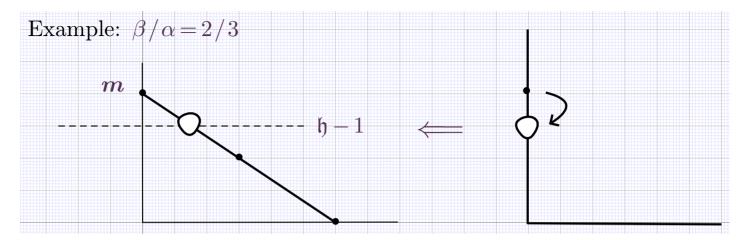
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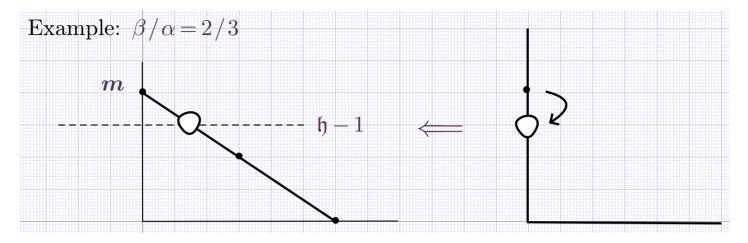
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The assumption  $\tilde{\mathfrak{h}} = \mathfrak{h}$  is equivalent to say that  $\operatorname{New}_{(x,y)}(\partial)$  is edge-unstable, which contradicts the hypothesis of the Theorem.

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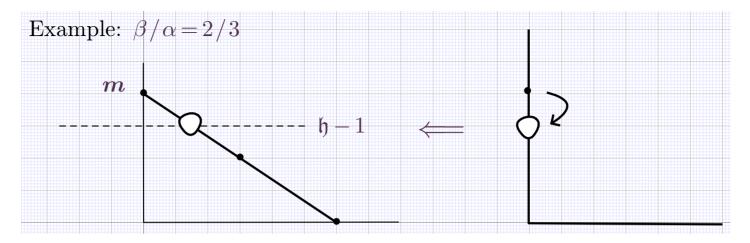


After blowing-up, followed by an arbitrary translation  $y \rightarrow y + \xi$ , we have

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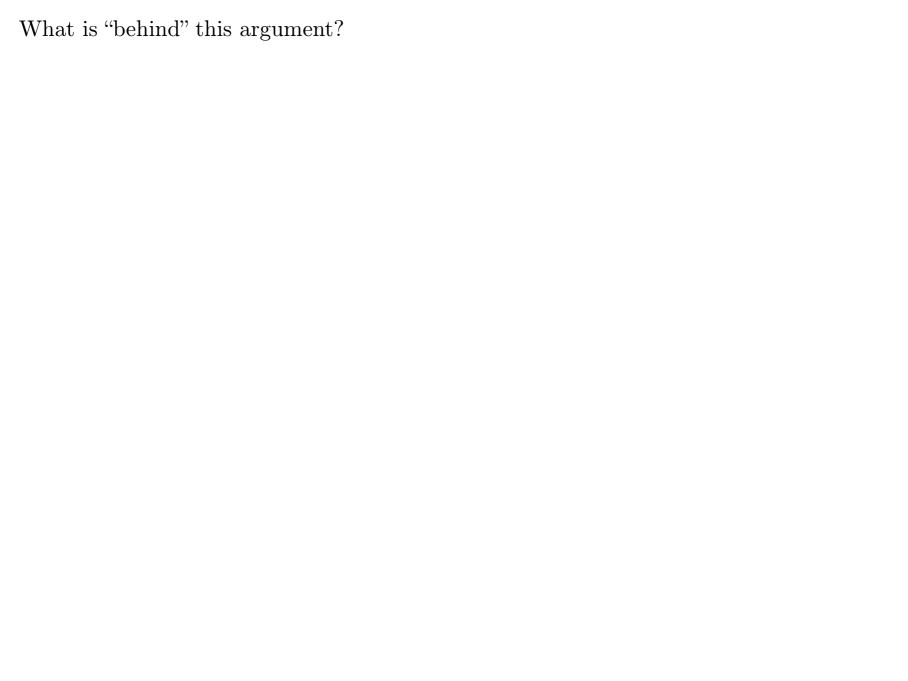


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(Abhyankar called this argument the "lazy Tschirnhaussen").



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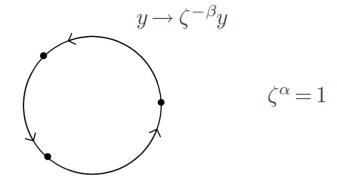
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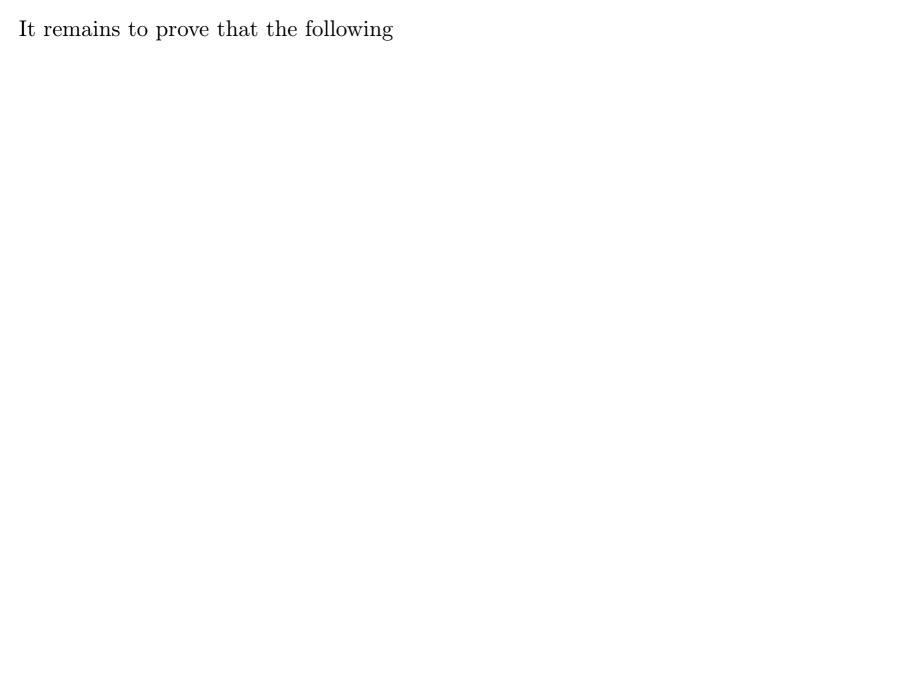
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**Proof:** We start with an arbitrary adapted coordinate system  $(x, y_0)$ .

- 1) If  $New_{(x,y_0)}(\partial)$  is edge-stable, we stop
- 2) If  $\text{New}_{(x,y_0)}(\partial)$  is edge-unstable, we choose a polynomial coordinate change  $(x,y_0) \rightarrow (x,y_1)$ , where

$$y_1 = y_0 + \xi_0 x^{k_0}, \qquad k_0 = \beta_0 / \alpha_0$$

eliminates the main edge  $\mathfrak{e}_0$ .

Indeed, assume the contrary. Then, we end-up with an infinite sequence of coordinate changes

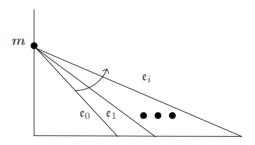
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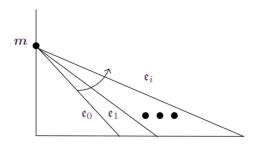
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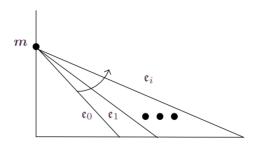


The composition of these maps converges to a formal coordinate change  $\hat{y_{\infty}} = y_0 + \sum \xi_i x^{ki}$ 

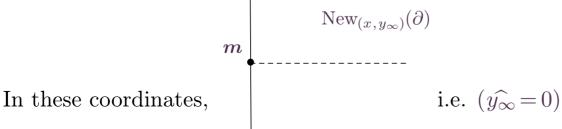
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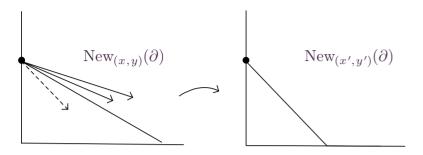
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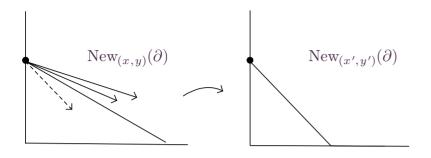


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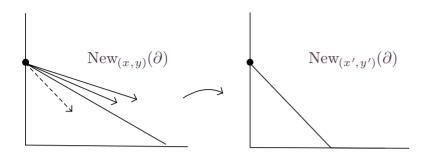


i.e.  $(\widehat{y_{\infty}} = 0) \subset \text{Nilp}(M, \mathcal{F})$ . Contradiction.



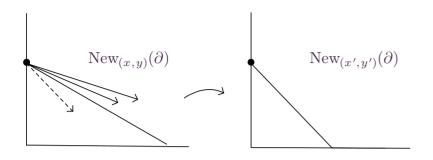


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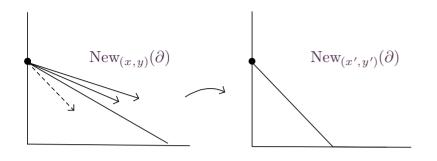
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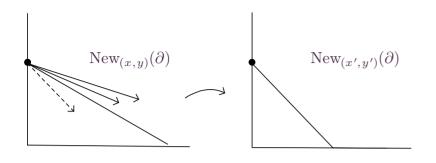


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is has a main edge  $\mathfrak{e}'$  of slope  $k < \beta/\alpha$  (because the action of  $y \to y + \xi x^k$  on  $\operatorname{New}_{(x,y)}(\partial)$  is effective).

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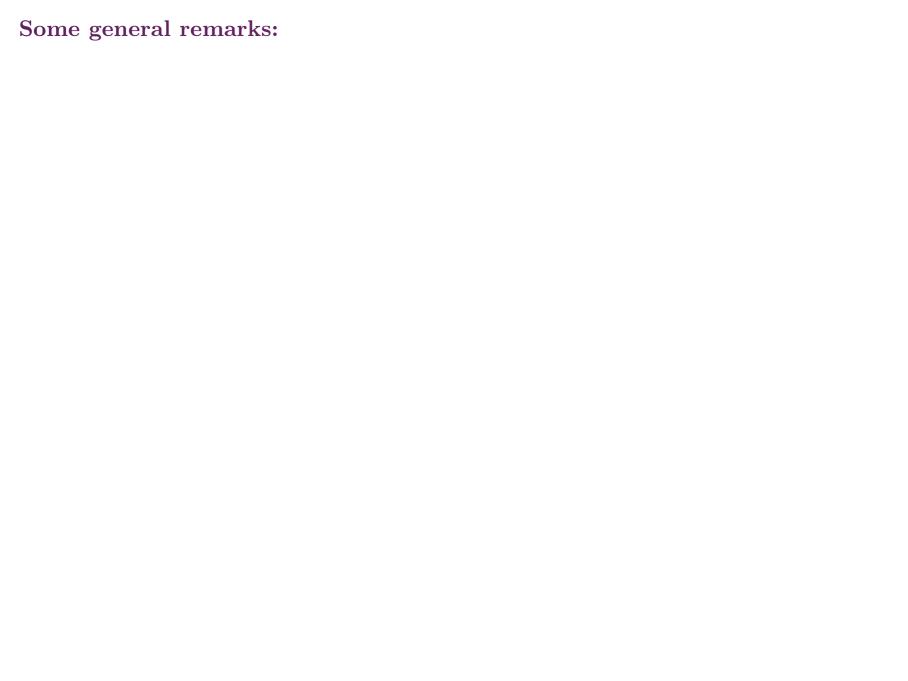
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But this contradicts the fact that the **inverse transformation**  $y = y_1 - \xi x^k$  eliminates the main edge.



## Some general remarks:

1) We cannot expect to obtain a **fully convergent Tchirnhaussen preparation** (or, more generally, a maximal contact hypersurface which would allow to use induction in the dimension)

Recall that, in the classical case of a germ of singular hypersurface S, this corresponds to choose a local equation of the form

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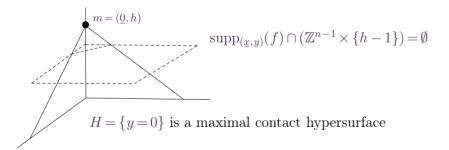
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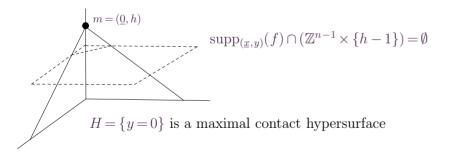
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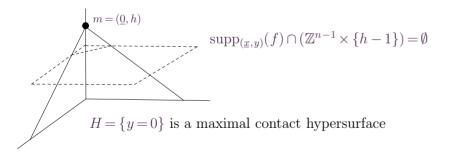
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$$\delta = (\operatorname{ad}_{\partial/\partial y})^h \partial = \left( \left\lceil \frac{\partial}{\partial y}, \cdot \right\rceil \right)^h \partial = \left\lceil (h+1)! \, b \, y \, \frac{\partial}{\partial y} + h! \, a \, x \, \frac{\partial}{\partial x} \right\rceil + (\text{terms of higher order})$$

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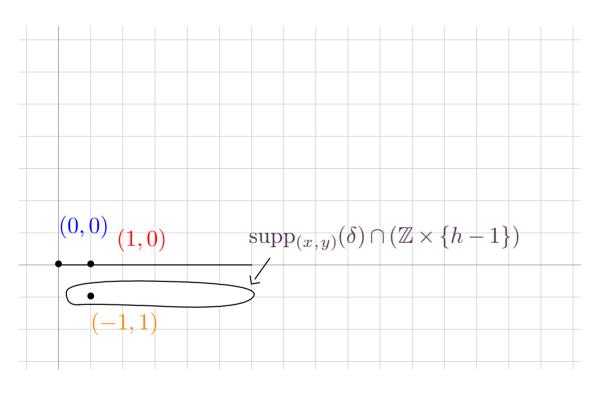
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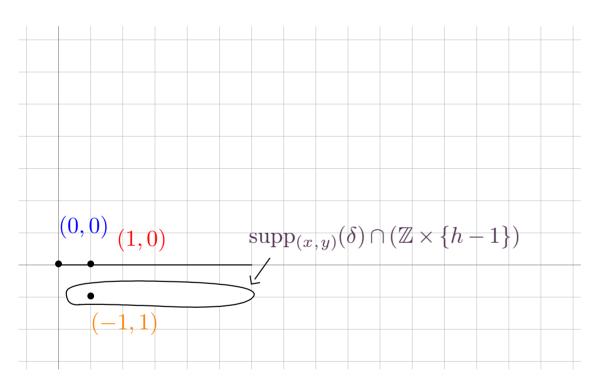
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$$\delta(y-f) \subset \langle y-f \rangle$$

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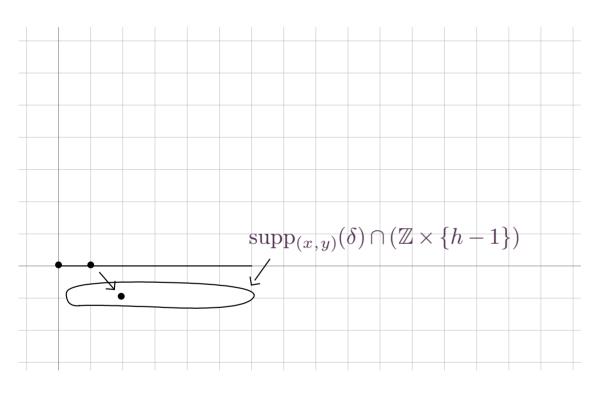


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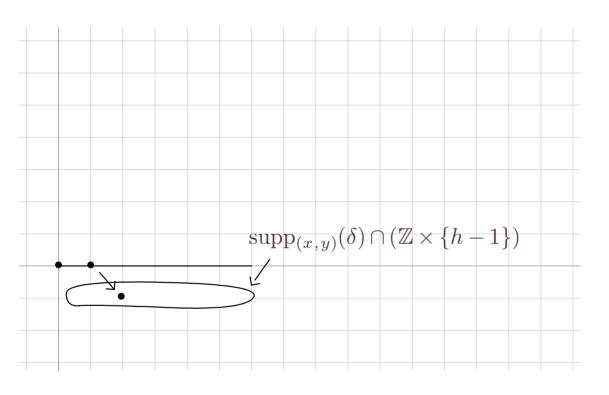


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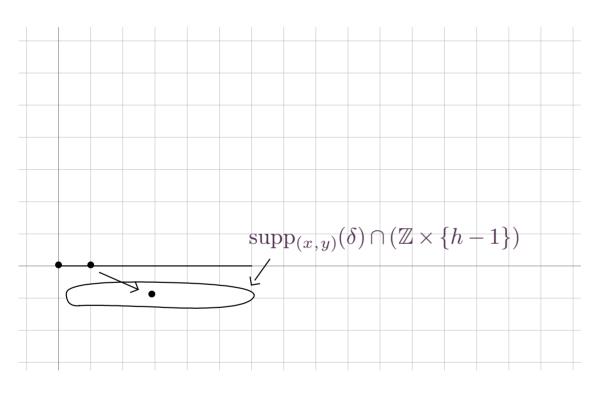


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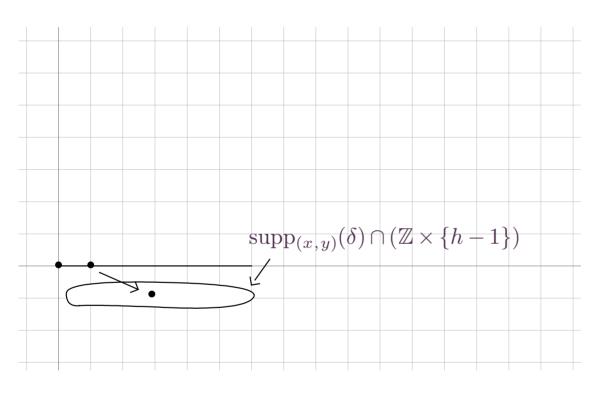


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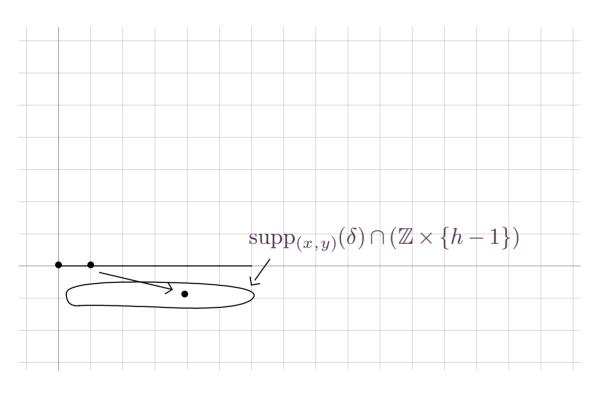


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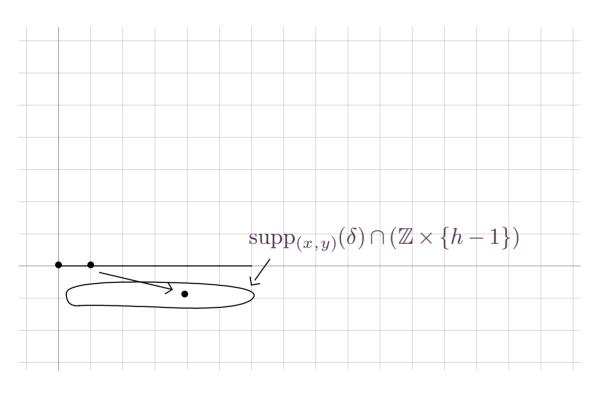


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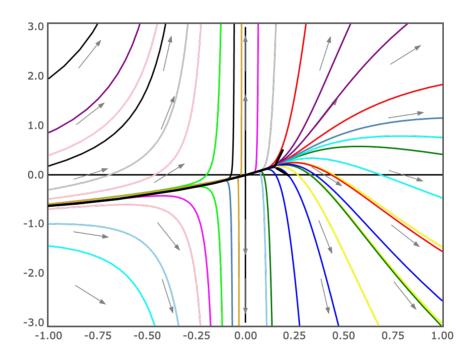


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which is the so-called "center manifold" of the Euler's equation.

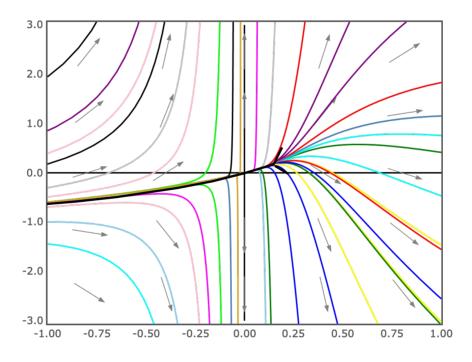
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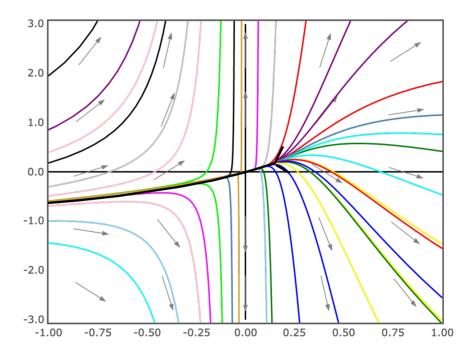
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But which is a  $C^{\infty}$ -curve, lying on the pfaffian extension of  $\mathbb{R}_{an}$ .



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- 2) What to do with the final models in dimension three? (There is no such well developed theory)
- 3) Interesting particular case for the Hilbert's 16<sup>th</sup> problem: The case "2+1".

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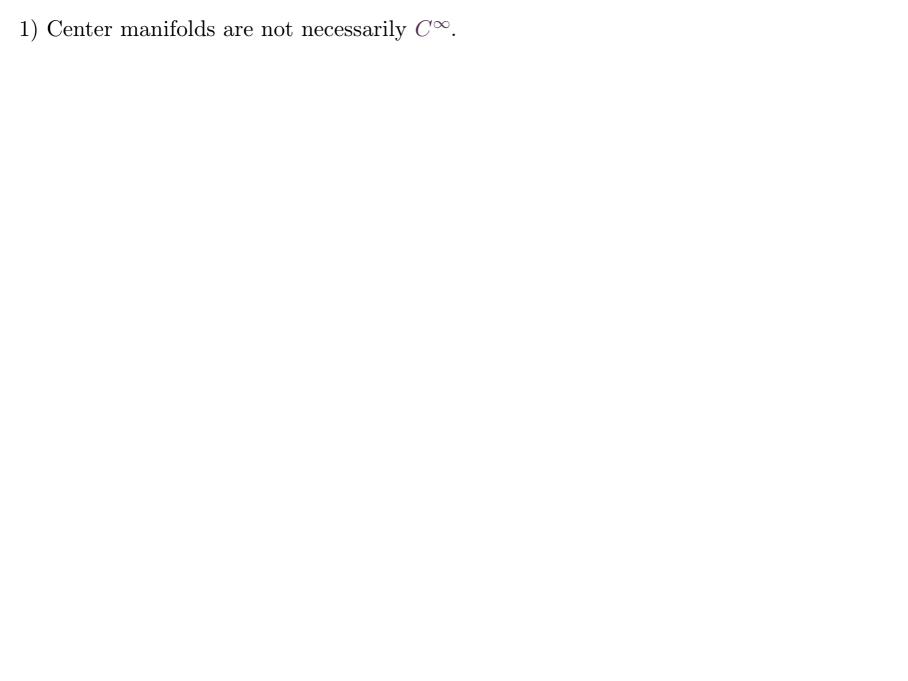
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4) New ideas for dimension greater or equal than four (The Kempf's unstability approach)

Some new phenomena in for final models in dimension three...



Example (van Strien 1979 - further simplyfied by M. Mcquillan)

$$\partial = x y \frac{\partial}{\partial y} + \left(z - \frac{y}{1 - y}\right) \frac{\partial}{\partial z}$$

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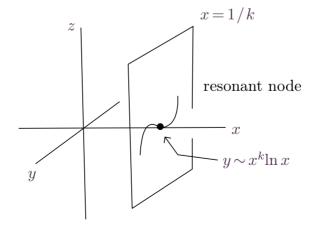
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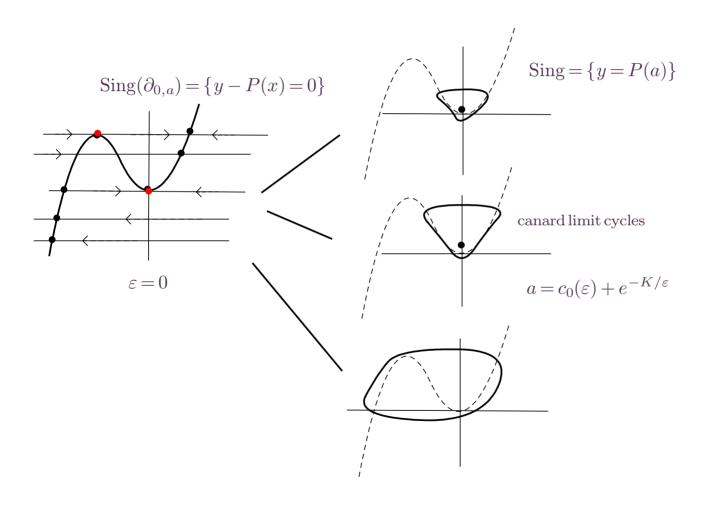
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$$\partial_{\varepsilon,a} = \left(y - \frac{x^2}{2} - \frac{x^3}{3}\right) \frac{\partial}{\partial x} + \varepsilon (a - x) \frac{\partial}{\partial y}, \qquad (x, y) \in \mathbb{R}^2, \varepsilon \in \mathbb{R}_{\geqslant 0}$$

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$$\left(1 - \frac{x^2}{2} - \frac{x^3}{3}\right) \frac{\partial}{\partial x} - \frac{\varepsilon x}{2} \left(y \frac{\partial}{\partial y} - x \frac{\partial}{\partial x} - 2\varepsilon \frac{\partial}{\partial \varepsilon}\right), \quad F = \{d(y^2 \varepsilon) = 0\}$$

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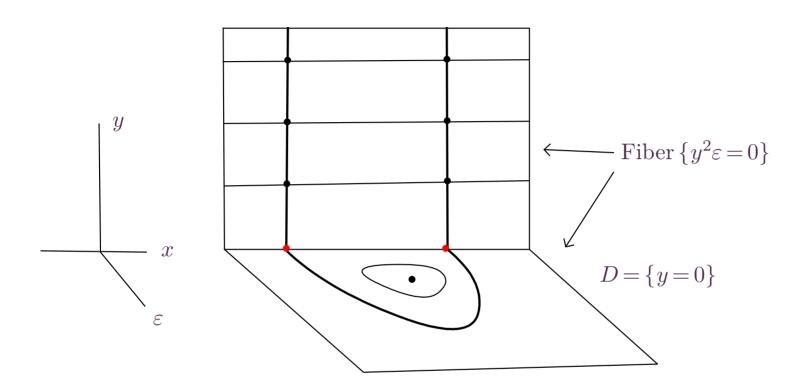
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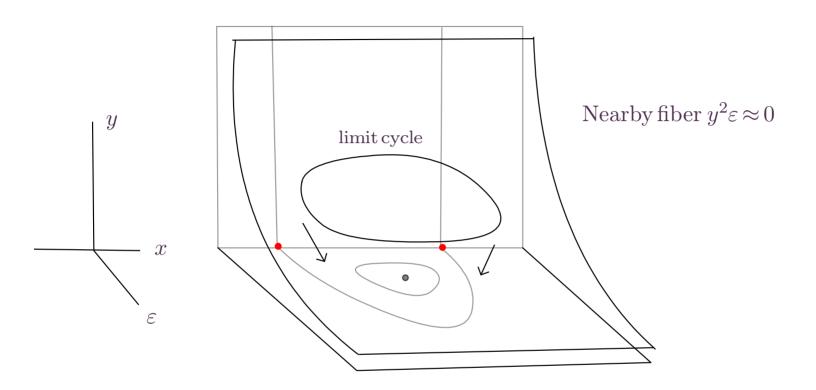
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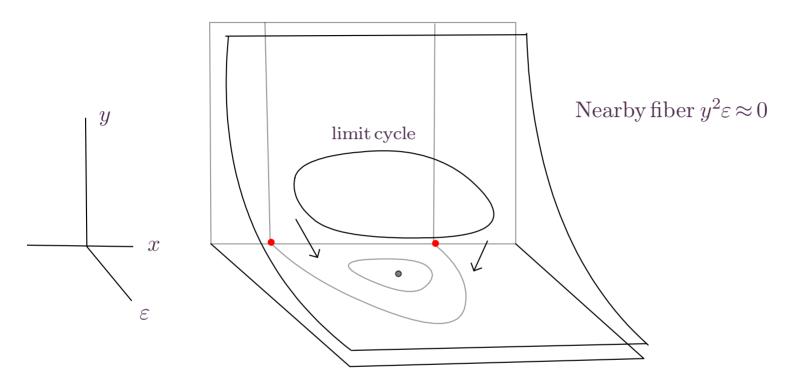
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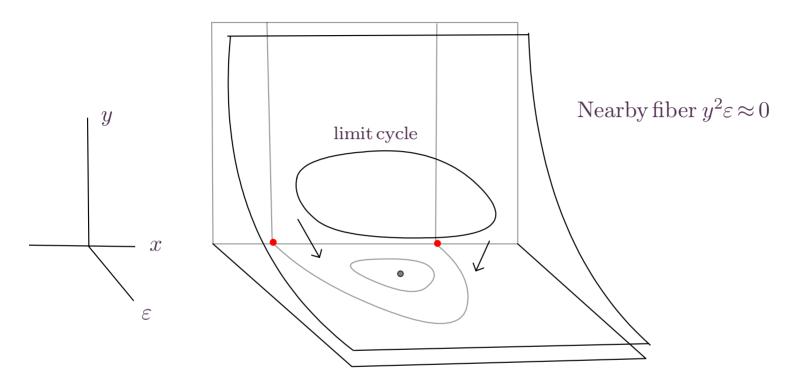
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(c.f. a very nice recent book of Maeschaalk, Dumortier, Roussarie - Canard cycles: from birth to transition).