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Remark: In general, we cannot expect to have a single global generator for a foliation.



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Holonomy Groupoid







any path $p \mathop{\rightarrow} q$ on L can be lifted to nearby leafs



ol:
$$(\Sigma, p) \rightarrow (\Omega, q)$$

$$\operatorname{hol} \in \operatorname{Diff}^{\omega}(\Sigma \to \Omega)$$

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(see the course of Patrick...)
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Elementary germs - and some words about classical normal forms... (over \mathbb{C})

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Elementary germs - and some words about classical normal forms... (over \mathbb{C})

A germ of vector field ∂ at $p \in M$ defines a derivation of the local ring $(\mathcal{O}, \mathbf{m}) = (\mathcal{O}_p, \mathbf{m}_p)$. Namely, in local coordinates $x = (x_1, \dots, x_n)$ we can write

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(i.e. that $\partial \in \operatorname{End}_{\mathbb{C}}(\mathcal{O})$ stabilizes the maximal ideal)

Flow-box Theorem Then, there exists local analytic coordinates $(f, g_1, \ldots, g_{n-1})$ such that

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Let us assume that u = 1 to simplify.

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Therefore $f, \Phi(g_1), \ldots, \Phi(g_n)$ is the required new coordinate system.

Then, (by Leibniz' rule) $\partial(\boldsymbol{m}^{k+1}) \subset \boldsymbol{m}^{k+1}$ for each $k \in \mathbb{N}$, and ∂ induces an sequence of endomorphism $\{\partial^k\}_k$ on the jet spaces

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Moreover, ∂_s and ∂_n are derivations of $\hat{\mathcal{O}} = \lim_{k \to \infty} J^k$ (see Jean Martinet - Exposé Bourbaki'81).

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Definition. A germ of vector field ∂ is *elementary* if:

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is generated (over \mathbb{C}) by the monomials $x^k = x_1^{k_1} \dots x_n^{k_n}$ such that $\langle k, \lambda \rangle = \alpha$.

The set of diagonal vector fields

$$L(\mu) = \sum_{i=1}^{n} \mu_i x_i \frac{\partial}{\partial x_i}, \qquad \mu \in \mathbb{C}^n$$

forms an abelian Lie \mathbb{C} -subalgebra, i.e. $[L(\mu), L(\lambda)] = 0$.

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where k ranges over the subset $\mathbb{Z}^n \setminus \{0\}$ such that $\langle \lambda, k \rangle = 0$. These are the **resonant monomials**.

$$\partial = (x + \dots) \frac{\partial}{\partial x} - (y + \dots) \frac{\partial}{\partial y}$$

Then, $\operatorname{Spec}(\partial|_{J^1}) = \{1, -1\}$ and the resonant monomials are $(xy)^k$, $k \in \mathbb{Z}$.

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$$\partial = \underbrace{\left(x\frac{\partial}{\partial x} - y\frac{\partial}{\partial y}\right)}_{\partial_s} + \underbrace{\sum_{k \ge 1} (xy)^k \left(a_k x\frac{\partial}{\partial x} + b_k y\frac{\partial}{\partial y}\right)}_{\partial_n}$$

where u = xy is the generator of the subring ker (∂_s) . By further reductions, we can write

$$(1+F)\left(\left(x\frac{\partial}{\partial x}-y\frac{\partial}{\partial y}\right)+\frac{u^n}{1+\rho u^n}\left(x\frac{\partial}{\partial x}+y\frac{\partial}{\partial y}\right)\right) \quad \text{or} \qquad (1+F)\left(x\frac{\partial}{\partial x}-y\frac{\partial}{\partial y}\right)$$

for some $F \in \mathbb{C}[[u]]$ of order ≥ 1 , $n \geq 1$ and $\rho \in \mathbb{C}$.

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The corresponding differential system is given by

$$\left(\frac{1}{u^{n+1}} + \rho \frac{1}{u}\right) du = \frac{dv}{v}$$

and, by direct integration,

$$I = \frac{1}{n u^n} + \rho \ln u - \ln v$$

This is a first integral of the vector field (namely, $\partial I = 0$). It is an element of $\mathbb{R}_{an,exp}$.

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and the first integral is simply $I = x^{\mu}y^{\lambda}$.





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Over \mathbb{C}^2 : There are several **rigidity phenomena**

E.g. Some analytic invariants are topologically determined (for instance, linearizability).

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We can recover the (orbital) analytic class of the saddle from the analytic class of one of these maps (once we fix the ratio μ/λ)

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Definition: Two germs of vector fields $\partial, \tilde{\partial}$ are **orbitally analytic equivalent** if there exists a unit $u \in \mathbb{C}\{x\}$ such that ∂ is analytically conjugated to $u \tilde{\partial}$.
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The problem is reasonably well-understood for **elementary singularities in dimension two** (modulo some very hard *small divisor problems*) see e.g. Dulac,Ecalle,Ilyashenko,Martinet,Ramis,Yoccoz and Perez Marco,... works.

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This problem is much less understood for vector fields higher dimensions.

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Example: (Cerveau-Moussu 1988) The cuspidal singularity

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The cusp $\Gamma = \{f = 0\}$ is an invariant curve.

Γ





There are two **distinct** corner transition maps.









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Blow-up 3: $x \rightarrow x$, $y \rightarrow xy$

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(3:1)
(1:6)
(1:2)

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The foliation is now organized in a neighborhood of the exceptional divisor..



Can we recover the analytic moduli from the transverse behaviour?



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(Moussu) The vanishing holonomy $\operatorname{Hol}(\mathcal{F}, L) = \langle f, g \in \operatorname{Diff}(\mathbb{C}, 0) | f^2 = g^3 = \operatorname{id} \rangle$ characterizes the analytic class of the germ of foliation.

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(in fact, $p \in \operatorname{Nilp}(M, \mathcal{F}) \iff \partial(\boldsymbol{m}_p) \subset \boldsymbol{m}_p$ and $\partial_1 \in \operatorname{End}_{\mathbb{C}}(\boldsymbol{m}_p / \boldsymbol{m}_p^2)$ is a nilpotent endomorphism, for ∂ some arbitrarily chosen local generator).
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Alternatively,

$$p \in \operatorname{Nilp}(M, \mathcal{F}) \iff \forall k \in \mathbb{N} \exists n \in \mathbb{N} : (\partial_k)^n = 0$$

where $\partial_k: J^k \to J^k$ is the induced derivation on the k^{th} jet.

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We further say that \mathcal{F} is **tightly adapted** to D if there exists an index i such that

 $\partial(\langle f^i\rangle)\not\subset \langle f^{i+1}\rangle$

In other words, for $E = (x_1 \dots x_k = 0)$,

$$\partial = \sum_{i=1}^{k} a_i \left(x_i \frac{\partial}{\partial x_i} \right) + \sum_{i=k+1}^{n} a_i \frac{\partial}{\partial x_i}$$

with $a_1, \ldots, a_n \in \mathbb{C}\{x\}$ such that $\langle a_1, \ldots, a_n \rangle \not\subset \langle x_i \rangle$, for each $i = 1, \ldots, k$.

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 $b \neq 0$: The generic point on the divisor is non-singular







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We can never get rid of saddle points...

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We can never get rid of a node if $\rho \notin \mathbb{Q}$.

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After m directional blowing-ups: $x \mathop{\rightarrow} x, \, y \mathop{\rightarrow} x \, y$

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First integral
$$h = (x^m y) \exp\left(\frac{1}{kx^k}\right)$$



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is tangent to the Whitney umbrella $W = y^2 - zx^2$.

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with $\beta \notin \frac{1}{2}\mathbb{Z}_{>0}$, $\lambda \in \mathbb{C}^*$.



Formal expansion of the "handle"

$$y = \tau(z) = \sum \tau_n z^n, \quad \tau_n \sim \lambda \, (n!)^2$$

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We cannot take the handle as a blowing-up center because it is non-analytic.

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The charts of a weighted-blowing up
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We interpret $(y_{1,..}, y_n)$ as an **orbifold chart** on \widetilde{M} . Namely the affine space \mathbb{C}^n equipped with an action of the cyclic group $\mathbb{Z}/\omega_1\mathbb{Z}$, defined by

$$y_1 \to \xi y_1, \qquad \text{For } 2 \leqslant k \leqslant n \colon y_k \longrightarrow \xi^{-\omega_k} y_k$$

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The glueing of these charts equipps \widetilde{M} with the structure of an orbifold.

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An **orbifold** is a pair (M, \mathcal{U}) where M is paracompact Hausdorff topological space and \mathcal{U} is a maximal orbifold atlas on M.

A sub-variety $Y \subset M$ is a sub-orbifold if for each point $p \in Y$ there exists a local chart (U, G, ϕ) such that $\phi^{-1}(Y)$ is a G-invariant submanifold of U.

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 $X = \operatorname{Spec} \mathbb{C}[x, y]^G$ (ring of invariants)

$$\mathbb{C}[x, y]^G = \mathbb{C}[x^2, xy, y^2]$$
$$X = \operatorname{spec} \mathbb{C}[u, v, w] / (v^2 - uw)$$

 \boldsymbol{X} is the quadratic cone.

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given by $\Phi(t, \bar{x}) = t^{\omega} \bar{x}$. The exceptional divisor is the **boundary**

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This is a non-trivial topological restriction.

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More generally, all automorphisms obtained by integrating the Lie algebra (over $\mathbb C$) generated by

$$\left\{x\frac{\partial}{\partial x}, y\frac{\partial}{\partial y}, x^l\frac{\partial}{\partial y}, y^m\frac{\partial}{\partial x} \mid m \ge 1, l \ge \beta\right\}$$

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Example:
$$\partial = x \frac{\partial}{\partial x} + n y \frac{\partial}{\partial y}$$
, $n \in \mathbb{Z}_{>0}$.

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Transformation of the logarithmic basis

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The solution curves of ∂ are precisely the orbits of the torus action $t \cdot (x, y) = (tx, t^n y)$.

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In the x-chart: $x \rightarrow x^2, y \rightarrow x^3 y$

$$\partial = xy\left(x\frac{\partial}{\partial x} - 3y\frac{\partial}{\partial y}\right) + 3xy^{-1}\left(y\frac{\partial}{\partial y}\right) = x\left(xy\frac{\partial}{\partial x} + 3(1-y^2)\frac{\partial}{\partial y}\right)$$

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The divisor $\{x=0\}$ is contained in the nilpotent locus. We factor out x and write

$$\partial_1 = x y \frac{\partial}{\partial x} + 3(1-y^2) \frac{\partial}{\partial y}$$

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and, factoring out y, we obtain

$$\partial_2 = 2(1-x^3)\frac{\partial}{\partial x} - x^2y\frac{\partial}{\partial y}$$

In the $y\text{-chart: }x \mathop{\rightarrow} y^2 x, \, y \mathop{\rightarrow} y^3$

$$\partial = 2x^{-1}y\left(x\frac{\partial}{\partial x}\right) + 3x^2y^{-1}\left(y\frac{\partial}{\partial y}\right)$$

$$\rightarrow \qquad \partial = 2x^{-1}y\left(x\frac{\partial}{\partial x}\right) + x^2y\left(y\frac{\partial}{\partial y} - 2x\frac{\partial}{\partial x}\right) = y\left(2(1-x^3)\frac{\partial}{\partial x} + x^2y\frac{\partial}{\partial y}\right)$$

and, factoring out y, we obtain

$$\partial_2 = 2(1-x^3)\frac{\partial}{\partial x} - x^2 y \frac{\partial}{\partial y}$$



Local symmetries of the foliated orbifold

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$$\pi_1(L) = \{\gamma, \eta, \rho | \gamma^2 = \eta^3 = 1, \rho = \gamma \eta \}$$

$$\partial_1 = x y \frac{\partial}{\partial x} + 3(1 - y^2) \frac{\partial}{\partial y} \quad \bigcirc \quad \mathbb{Z}/2\mathbb{Z}$$

$$\partial_1 = x y \frac{\partial}{\partial x} + 3(1 - y^2) \frac{\partial}{\partial y} \quad \bigcirc \quad \mathbb{Z}/2\mathbb{Z}$$

$$g \cdot x = -x, \quad g \cdot y \to -y$$
$$g \cdot \partial_1 = -\partial_1$$



$$\partial_1 = x y \frac{\partial}{\partial x} + 3(1 - y^2) \frac{\partial}{\partial y} \quad \bigcirc \quad \mathbb{Z}/2\mathbb{Z}$$

 $g \cdot x = -x, \quad g \cdot y \to -y$

$$g \cdot \partial_1 = -\partial_1$$



Other chart

$$\partial_2 = 2(1 - x^3)\frac{\partial}{\partial x} - x^2 y \frac{\partial}{\partial y}$$
$$g \cdot x = \xi^{-2} x, \quad g \cdot y = \xi y, \qquad (\xi^3 = \mathrm{id})$$
$$g \cdot \partial_2 = \xi^2 \partial_2$$



Elimination of nilpotent points in dimension two