## Solitons in Bryant's $G_2$ -Laplacian flow.

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## The Lie group $G_2$ and 3-forms in 7 dimensions

In 1886 Engel suggested to Killing a way to construct a 14-dimensional simple Lie group using *generic* 3-forms in 7 variables; details not settled until 1907.

A complex k-form on  $\mathbb{C}^n$  is **generic** if its  $GL(n,\mathbb{C})$ -orbit is *open* in  $\Lambda^k(\mathbb{C}^n)^*$ . Given any k-form  $\phi$  on  $\mathbb{C}^n$  define a Lie subgroup of  $GL(n,\mathbb{C})$  by

$$G_{\phi} := \{ A \in GL(n, \mathbb{C}) : A^* \phi = \phi \}.$$

For a generic k-form  $\phi$  we need dim  $GL(n,\mathbb{C})$  – dim  $G_{\phi}=\dim \Lambda^k(\mathbb{C}^n)^*$ .

For k=3, n=7 the dimension of  $G_{\phi}$  for a generic form must be

$$dim G_{\phi} = dim GL(7,\mathbb{C}) - dim \Lambda^3(\mathbb{C}^7)^* = 49 - 35 = 14$$

which is the dimension of  $G_2$ !

## Generic 3-forms in $\mathbb{R}^7$ , $G_2$ and the octonions

In 1900 Engel showed:

- there is exactly one  $GL(7,\mathbb{C})$  orbit of generic 3-forms in  $\mathbb{C}^7$
- For every generic 3-form  $\phi$  the isotropy group  $G_{\phi}$  is isomorphic to  $G_2(\mathbb{C})$ .

In 1907 Reichel (Engel's student) considered generic real 3-forms in  $\mathbb{R}^7$ .

- There are 2 types of generic (real) 3-forms in  $\mathbb{R}^7$ .
  - □ Case 1:  $\exists$  an invariant symmetric bilinear form of signature (7,0);  $G_{\phi}$  is isomorphic to the compact real simple Lie group  $G_2 \subset SO(7)$ . We will call such a generic real 3-form  $\varphi$  positive and write  $\varphi \in \mathcal{P}^3(\mathbb{R}^7)$ .
  - $\square$  Case 2:  $\exists$  invariant symmetric bilinear form of signature (4,3).

Define a vector cross-product and a 3-form on  $\mathbb{R}^7=\mathrm{Im}(\mathbb{O})$  using octonionic multiplication and the norm

$$u \times v := \operatorname{Im}(uv),$$
  
$$\varphi_0(u, v, w) := \langle u \times v, w \rangle = \langle uv, w \rangle.$$

Then  $\varphi_0 \in \mathcal{P}^3(\mathbb{R}^7)$  so  $G_2 = Aut(\mathbb{O}) = \{A \in GL(7,\mathbb{R}) | A^*\varphi_0 = \varphi_0\}.$ 

### Positive 3-forms & G<sub>2</sub>-structures on 7-manifolds

For an oriented smooth 7-manifold M and  $p \in M$ 

$$\mathcal{P}_{p}(M) := \{ \varphi \in \Lambda^{3} T_{p}^{*} M \, | \, \iota^{*} \varphi_{0} = \varphi \text{ for } \iota : T_{p} M \to \mathbb{R}^{7} \}$$

where  $\iota$  is any orientation preserving isomorphism.  $\mathcal{P}(M)$  denotes the bundle over M with fibre  $\mathcal{P}_{\mathcal{P}}(M)$ .

A 3-form  $\varphi$  on M is *positive* if  $\varphi$  is a section of  $\mathcal{P}(M)$ , i.e.  $\varphi_p \in \mathcal{P}_p(M) \ \forall p$ .

Each positive 3-form on M defines a reduction of the frame bundle  $\mathcal{F}M$  to a principal subbundle of  $\mathcal{F}M$  with fibre  $G_2$ , i.e. a  $G_2$ -structure on M that induces the given orientation on M.

Positive 3-forms on  $M \leftrightarrow \text{(oriented)} G_2\text{-structures on } M$ .

Existence of a positive 3-form on  $M^7$  is a topological question: M needs to be orientable and spinnable.

## G<sub>2</sub> as a Riemannian holonomy group

Possible holonomy groups of Riemannian manifolds<sup>1</sup> are extremely limited:

- 5 possible infinite families and
- 2 exceptional cases, the Lie groups  $G_2$  and  $Spin_7$  (in dims 7 and 8)

Both exceptional cases and 3 of the infinite families constitute the special holonomy metrics. (The other 2 are generic Riemannian/Kähler metrics).

Special (exceptional) holonomy metrics are necessarily Einstein (Ricci-flat). Proving existence of metrics of special holonomy (locally; complete metrics on noncompact and on compact spaces) took many years (1955–1997), and involved many deep developments in geometry and geometric analysis:

- Yau's proof of the Calabi conjecture (1978) settled affirmatively the cases with holonomy SU(n) and Sp(n). Yau used analytic methods to prove existence of solutions to a complex Monge-Ampère equation a fully nonlinear scalar elliptic equation.
- For the two exceptional holonomy cases we can no longer reduce to a scalar equation as in the SU(n) case. The best one can do involves systems of nonlinear first-order PDEs.

<sup>&</sup>lt;sup>1</sup>Berger 1955: simply connected irreducible non-locally-symmetric case

# 1st-order PDE system for $G_2$ holonomy metrics

**Important fact:** The holonomy group  $Hol_g(M)$  determines the parallel tensors on (M,g). In particular  $Hol_g(M) \subseteq G_2 \subset SO(7)$  implies

 $M^7$  admits a g-parallel positive 3-form  $\varphi$ .

Converse: How to get a  $G_2$ -holonomy metric from a  $G_2$ -structure?

#### **Theorem**

Let  $(M, \varphi, g_{\varphi})$  be a  $G_2$ -structure; the following are equivalent

- **1.** Hol $(g_{\varphi}) \subseteq G_2$  and  $\varphi$  is the induced 3-form
- **2.**  $d\varphi = d^*\varphi = 0$ , where  $d^*$  is defined using Hodge star \* w.r.t.  $g_{\varphi}$ .

Call such a  $G_2$ -structure a torsion-free  $G_2$  structure.

**2** is *nonlinear* in  $\varphi$ :  $g_{\varphi}$  depends nonlinearly on  $\varphi$  and  $d^*$  depends on  $g_{\varphi}$ . By writing equation for 3-form  $\varphi$  (not metric g directly) and allowing  $\operatorname{Hol}(g_{\varphi}) \subseteq G_2$  we obtain *differential* (not integro-differential) equations.

**2** is a 1st-order system of 49=(35+21-7) equations on the 35 coeffs of  $\varphi$ !

## **Bryant's Laplacian flow**

There is a natural flow on *closed*  $G_2$ -structures. Solve

$$\frac{d\varphi_t}{dt} = \Delta_{\varphi_t} \varphi_t \tag{LF}$$

with initial condition  $\varphi_0$  satisfying  $d\varphi_0 = 0$ . (Then  $d\varphi_t = 0$  for all t.)

■ Induced metric  $g_t$  evolves under (LF) by

$$\frac{dg_t}{dt} = -2\text{Ric}(g_t) + \text{ terms quadratic in torsion of } \varphi_t$$

- Stationary points of (LF) are exactly torsion-free  $G_2$ -structures.
- (LF) is the (upward) gradient flow for Hitchin's volume functional

$$\operatorname{vol}(\varphi) := \frac{1}{7} \int_{M} \varphi \wedge *\varphi$$

when restricted to cohomology class of  $\varphi_0$ .

Critical points of  $vol(\varphi)$  in  $[\varphi]$  are maxima (strict modulo diffeos).

- On a compact manifold  $vol(\varphi_t)$  is increasing along (LF).
- Bryant-Xu : Short-time existence & uniqueness of solutions to (LF)
- Lotay–Wei: Torsion-free  $G_2$ -structures are stable under (LF).

## Solitons in the Laplacian flow

 $G_2$ -structure  $\varphi$ , vector field X,  $\lambda \in \mathbb{R}$  satisfying

$$\begin{cases} d\varphi = 0, \\ \Delta_{\varphi}\varphi = \lambda\varphi + \mathcal{L}_{X}\varphi. \end{cases}$$

⇔ self-similar solution of Laplacian flow

$$\varphi_t = k(t)^3 f^* \varphi, \qquad \frac{df}{dt} = k(t)^{-2} X, \qquad k(t) = \frac{3 + 2\lambda t}{3}$$

$$\lambda > 0$$
: expanders (immortal solutions, i.e. exist up to  $t = +\infty$ )

 $\lambda=$  0: steady solitons (eternal solutions, i.e. exist for all time  $t\in\mathbb{R}$ )

$$\lambda <$$
 0: shrinkers (ancient solutions, i.e. exist backwards to  $t = -\infty$ )

- Non-steady soliton  $\Rightarrow \varphi$  exact
- Solitons on a *compact* manifold are stationary or expanders
- Scaling behaviour:  $(\varphi, X)$  is a  $\lambda$ -soliton  $\Leftrightarrow (k^3 \varphi, k^{-2} X)$  is a  $k^{-2} \lambda$ -soliton.

### Our motivation and overview of results

**Motivation:** in most geometric flows solitons provide the models for singularity formation. So we look for (symmetric) solitons of Laplacian flow.

**Goal:** Find asymptotically conical (AC)  $G_2$  solitons with cohomogeneity one: SU(3)-invariant ones on  $\Lambda^2_+\mathbb{C}P^2$  and Sp(2)-invariant ones on  $\Lambda^2_+\mathbb{S}^4$ .

### Theorem (A)

 $\exists$  a 1-parameter family of steady solitons on  $\Lambda_+^2 \mathbb{C}P^2$  asymptotic with rate -1 to torsion-free cone (deformations of the Bryant-Salamon AC  $G_2$ -metric).

• AC steady solitons a new feature (compared to Ricci/Kähler-Ricci flow).

#### Theorem (B)

 $\exists$  an explicit AC shrinker with rate -2 on  $\Lambda^2_+\mathbb{S}^4$  and  $\Lambda^2_+\mathbb{C}P^2$ .

• Shrinkers are rare! Possible models for formation of conical singularities.

### Theorem (C)

 $\exists$  a 1-parameter family of complete expanders on  $\Lambda^2_+\mathbb{S}^4$  and on  $\Lambda^2_+\mathbb{C}P^2$ .

• Models for how Laplacian flow can smooth out certain conical singularities.

# Closed invariant $G_2$ -structures on $\Lambda^2_+ M^4 \setminus M$

For  $M=\mathbb{C}P^2$  or  $\mathbb{S}^4$ ,  $\Lambda_+^2M$  has a cohomogeneity one action by G=SU(3) or Sp(2).  $\Lambda_+^2M\setminus M$  is diffeomorphic to  $\mathbb{R}_+\times \Sigma$ , for  $\Sigma=SU(3)/T^2$  or  $\mathbb{C}P^3$ .

There are *G*-invariant forms  $\omega_1, \omega_2, \omega_3 \in \Omega^2(\Sigma)$  and  $\alpha \in \Omega^3(\Sigma)$  such that any closed *G*-invariant  $G_2$ -structure on  $\mathbb{R}_+ \times \Sigma$  with  $\|\frac{\partial}{\partial t}\| = 1$  can be written as

$$\varphi = (f_1^2\omega_1 + f_2^2\omega_2 + f_3^2\omega_3) \wedge dt + f_1f_2f_3\alpha, \qquad f_i : \mathbb{R}_+ \to \mathbb{R}_+$$

with

$$\frac{d(f_1 f_2 f_3)}{dt} = \frac{1}{2} (f_1^2 + f_2^2 + f_3^2).$$

$$(\#) \Rightarrow \frac{d}{dt} (f_1 f_2 f_3)^{1/3} \ge \frac{1}{2}.$$

For Sp(2)-invariance in addition require  $f_2 = f_3$ . Structure equations for  $\omega_i$ ,  $\alpha$  the same in both cases  $\Rightarrow \Lambda_+^2 \mathbb{S}^4$  case can be treated as a special case of  $\Lambda_+^2 \mathbb{C}P^2$  case where  $f_2 = f_3$ .

## **Closed invariant** $G_2$ **cones**

Helpful to analyse invariant  $G_2$ -structures on  $\mathbb{R}_+ \times \Sigma$  in terms of scale and homothety class of invariant metrics on  $\Sigma$ :

scale 
$$g := \sqrt[3]{f_1 f_2 f_3} = \sqrt[6]{\operatorname{vol}(\Sigma)}$$
  
homethety class  $\frac{f_1}{g}, \frac{f_2}{g}, \frac{f_3}{g}$ 

 $\varphi$  closed and homothety class constant implies g linear and  $\varphi$  conical:

$$d\varphi = 0 \implies \frac{dg}{dt} = \frac{1}{6} \left( \frac{f_1^2}{g^2} + \frac{f_2^2}{g^2} + \frac{f_3^2}{g^2} \right) \implies f_i = c_i t$$

with

$$6c_1c_2c_3 = c_1^2 + c_2^2 + c_3^2. (*)$$

Note: any positive triple  $(c_1, c_2, c_3)$  can be uniquely rescaled to satisfy (\*)  $\rightsquigarrow$  2-parameter family of closed conical  $G_2$ -structures on  $\mathbb{R}_+ \times SU(3)/T^2$ .

In other words, given homothety class on  $\Sigma$ , there is a unique choice of "cone angle" that makes it a closed cone.

### **Evolution equations**

On the face of it, the soliton condition for

$$\varphi = (f_1^2\omega_1 + f_2^2\omega_2 + f_3^2\omega_3) \wedge dt + f_1f_2f_3\alpha, \qquad X = u\frac{\partial}{\partial t}$$

is 2nd-order ODE system for  $(f_1, f_2, f_3, u)$  (with some constraints).

Can rewrite as a 1st-order system in 5 variables: the 3  $f_i$  and 2 variables determining the torsion 2-form  $\tau$  of  $\varphi$ .

- $\Rightarrow$  there is a 4-parameter family of local SU(3)-invariant solitons and a 2-parameter family of  $Sp_2$ -invariant solitons.
- **Q:** Which of these invariant solitons extends to a complete AC solution? *Tendency:* if  $\frac{f_1}{g}$ ,  $\frac{f_2}{g}$ ,  $\frac{f_3}{g}$  remain bounded as  $t \to \infty$  then asymptotic to closed cone.

**Rough strategy** for finding AC solitons on  $\Lambda_+^2 M = M \sqcup \mathbb{R}_+ \times \Sigma$ .

- **1.** Solutions on  $(0, \epsilon) \times \Sigma$  that extend smoothly across M at t = 0?
- **2.** Solutions for large t asymptotic to prescribed closed cone  $(c_1, c_2, c_3)$ ?
- **3.** Do solutions from **1** and **2** fit together?

# Initial value problem near zero section of $\Lambda^2_+ M^4$

Understand solutions near zero section of  $\Lambda_+^2 M$  à la Eschenburg-Wang.

$$\varphi = (f_1^2 \omega_1 + f_2^2 \omega_2 + f_3^2 \omega_3) \wedge dt + f_1 f_2 f_3 \alpha$$

on  $\mathbb{R}_+ \times \Sigma$  extends to smooth  $G_2$ -structure on  $\Lambda^2_+ M$  iff

 $f_1$  is odd with  $f_1'(0)=1$ , and  $f_2$  and  $f_3$  are even with  $m:=f_2(0)=f_3(0)\neq 0$ .

Resulting singular initial value problem has formal power series solutions that are convergent. (It is a **regular singular point** of 1st-order ODE system).

#### **Proposition**

For each  $\lambda \in \mathbb{R}$ , there is

- a 2-parameter family  $\varphi_{m,c}$  of solutions defined for small t that extend smoothly to a  $\lambda$ -soliton on (nhd of zero section in)  $\Lambda^2_+\mathbb{C}P^2$ ;
- 1-parameter subfamily  $\varphi_m = \varphi_{m,0}$  also defines  $\lambda$ -solitons on  $\Lambda^2_+\mathbb{S}^4$ .

Two scale-invariant parameters:  $\lambda m^2$  and c.

So up to scale: 2-parameter families of local expanders/shrinkers on  $\Lambda_+^2 \mathbb{C} P^2$  a 1-parameter family of local steady solitons on  $\Lambda_+^2 \mathbb{C} P^2$ 

### **Expanders**

#### Theorem (C)

For  $\lambda > 0$ , each  $\varphi_m$  extends to a complete solution with  $f_2 = f_3$ , and

$$\frac{f_i}{t} \to c_i$$

for  $(c_1, c_2, c_2)$  a closed cone with  $c_1 \leq c_2$ .

This gives 1-parameter families of expanders on both  $\Lambda^2_+\mathbb{S}^4$  and  $\Lambda^2_+\mathbb{C}P^2$ .

- Strong Expectations
- These solitons are all AC, with rate -2
- 1-1 correspondence with closed cones such that  $c_1 < c_2$ : any closed cone with  $c_2 = c_3$  on "one side" of the torsion-free cone  $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$  is the AC end of a unique expander

### Conjecture

For  $\lambda > 0$ , an open subfamily of  $\varphi_{m,c}$  (but not all) extend to complete solutions, defining a 2-parameter family of AC solitons on  $\Lambda_+^2 \mathbb{C} P^2$ .

## Rigidity/stability dichotomy of $\lambda \neq 0$ AC ends

For any  $\lambda \neq 0$ ,  $\exists$  a ! formal power series solution  $\mathcal{P}$  in  $t^{-1}$  determined by the cone and a solution of the ODE system that is smooth in a nhd of  $t=+\infty$  whose Taylor series is  $\mathcal{P}$ .

For  $\lambda < 0$ , for each closed cone  $(c_1, c_2, c_3)$  there is a **unique** solution defined for large t asymptotic to the given cone; so AC shrinker ends are rigid.

Given  $\lambda > 0$  and any closed cone  $(c_1, c_2, c_3)$ , we expect:

- $\exists$  a 2-parameter family of AC ends asymptotic to the given cone.
- Difference between two solutions is of order  $\exp(-\frac{\lambda}{6}t^2)$  \* polynomial.
- If  $c_2=c_3$ , then a 1-parameter subfamily has  $f_2=f_3$ .

Flow lines of this 4=(2+2)-parameter family of solutions fill open subset of 5-dimensional phase space, so **AC expander ends are stable**.

"Explanation" for shrinker/expander dichtomy: ODEs for expanders/ shrinkers asymptotic to given cone has an irregular singularity at  $t = +\infty$ .

- When  $\lambda>0$  other smooth solutions also with Taylor series  $\mathcal P$  at  $t=\infty$  exist because of exponentially small corrections of form  $\exp\left(-\frac{1}{6}\lambda t^2\right)*p(t)$ .
- When  $\lambda < 0$  the exponential terms blow-up as  $t \to \infty$ .

## Shrinkers: consequences of AC end rigidity

#### Heuristic for $\lambda < 0$ :

Invariant shrinkers on  $\mathbb{R}_+ \times SU(3)/T^2$  are flow lines in 5-dim phase space. In 4-dimensional space of flow lines

- 2-dimensional submanifold extends across zero section  $\mathbb{C}P^2 \subset \Lambda^2_+ \mathbb{C}P^2$
- 2-dimensional submanifold has AC behaviour

Expect transverse intersections  $\rightsquigarrow$  finitely many AC shrinkers on  $\Lambda^2_+ \mathbb{C}P^2$ .

Similarly, restricting attention to solutions with  $f_2 = f_3$ :

2-dimensional space of flow lines; 1-dim submanifold extends over special orbit; 1-dim submanifold has AC behaviour. Expect transverse intersections  $\rightsquigarrow$  finitely many AC shrinkers on  $\Lambda_{\perp}^2 \mathbb{S}^4$ .

In fact, can spot one explicit solution!

Theorem B: For  $\lambda = -1$ 

$$f_1 = t$$
,  $f_2^2 = f_3^2 = \frac{9}{4} + \frac{1}{4}t^2$ ,  $u = \frac{t}{3} + \frac{4t}{9+t^2}$ 

is an AC shrinker with rate -2 asymptotic to cone  $(1, \frac{1}{2}, \frac{1}{2})$ .

**Conjecture:** This is the unique  $Sp_2$ -invariant AC shrinker on  $\Lambda^2_+\mathbb{S}^4$ .

## Steady solitons

Significant qualitative differences from  $\lambda \neq 0$ :

Near special orbit  $\mathbb{C}P^2$ , only a 1-parameter family of solutions up to scale. Unique one with  $f_2 = f_3$ : static soliton from Bryant–Salamon AC  $G_2$ -mfd.

#### **Theorem**

No non-stationary steady solitons on  $\Lambda^2_+\mathbb{S}^4$ .

### **Decoupling**

- For  $\lambda = 0$ , the flow can be separated into evolution of scale g and evolution of 4 scale-normalised variables.
- Unique fixed point for the scale-normalised flow is the torsion-free cone; It is a stable fixed point.

### Theorem (A)

There exists a 1-parameter family (up to scale) of AC steady solitons on  $\Lambda^2_+ \mathbb{C}P^2$  all asymptotic to the torsion-free cone over  $SU(3)/T^2$ ; the family includes steady solitons with arbitrarily small torsion.

## Comparison with other flows: the steady case

- All known *steady* solitons in Ricci flow have *sub-Euclidean* volume growth:
- o the Bryant soliton; Appleton's resolutions of (some of) its quotients.
- o Bryant soliton known to appear in a finite-time singularity of RF.
- o known Kähler examples have at most half-dimensional volume growth (Cao, Conlon–Deruelle). *Not* seen in finite-time singular behaviour of KRF.
- Our steady AC  $G_2$  solitons most closely resemble Joyce-Lee-Tsui's (JLT) translating solitons in Lagrangian mean curvature flow (LMCF).
- Joyce conjectures JLT translating solitons can appear in finite-time singularities of LMCF if Floer homology is obstructed.
- $\circ$  Speculate that our steady  $G_2$  solitons can also arise as finite-time singularities of Laplacian flow on a compact 7-manifold.

(Our 2-parameter family of AC  $G_2$  expanders on  $\Lambda^2_+ \mathbb{C}P^2$  resembles JLT's family of exact Maslov-zero LMCF expanders asymptotic to pairs of transverse Lagrangian 3-planes).

## Comparison with other flows: shrinkers

**Ricci flow:** One obvious significant difference: absence of *compact* shrinkers in  $G_2$  flow; associated with positive curvature in RF, whereas scalar curvature is non-positive for closed  $G_2$ -structures.

General theory for *noncompact complete shrinkers* in RF is well-developed:

o their properties are a hybrid of those of positively curved Einstein manifolds and spaces with non-negative Ricci, e.g. at most Euclidean volume growth.

o AC (gradient) shrinkers are extremely rigid—manifestation of parabolic

- backwards uniqueness phenomenon, also seen in MCF.
- $\circ$  AC end behaviour of our (highly symmetric)  $G_2$  shrinkers some indication such strong rigidity also holds for AC  $G_2$  (gradient?) shrinkers.

**LMCF:** self-shrinkers exist and do occur but *not* in the Maslov-zero (graded) setting. **Q:** Is there any natural condition to impose in the  $G_2$  setting that would rule out our AC shrinkers on  $\Lambda_+^2 \mathbb{S}^4$  and  $\Lambda_+^2 \mathbb{C}P^2$ ?

**KRF:** Feldman-Ilmanen-Knopf (FIK) constructed symmetric ALE Kähler shrinkers; simplest FIK shrinker does appear as a finite-time blowup of KRF on 1-point blowup of  $\mathbb{C}P^2$  and is associated with blowing down the point.

Thanks for your attention!