

# Some Coupled Supersymmetries and Their Associated Bargmann Transforms

Cameron L. Williams

joint work with Bernhard G. Bodmann and Donald J. Kouri  
inspired by John R. Klauder

University of Houston  
Departments of Mathematics and Physics

48th Canadian Operator Symposium  
May 28, 2020

# The Oscillator Algebra

# The Oscillator Algebra

## Definition

Given a Hilbert space  $\mathfrak{H}$  and a closed and densely-defined operator  $a$  on  $\mathfrak{H}$ , the *oscillator algebra* is a Lie algebra generated by  $\{1, a, a^*, a^*a\}$  satisfying

# The Oscillator Algebra

## Definition

Given a Hilbert space  $\mathfrak{H}$  and a closed and densely-defined operator  $a$  on  $\mathfrak{H}$ , the *oscillator algebra* is a Lie algebra generated by  $\{1, a, a^*, a^*a\}$  satisfying

$$[a, a^*] = 1 \quad [a^*a, a] = -a \quad [a^*a, a^*] = +a^*.$$

# The Oscillator Algebra

## Definition

Given a Hilbert space  $\mathfrak{H}$  and a closed and densely-defined operator  $a$  on  $\mathfrak{H}$ , the *oscillator algebra* is a Lie algebra generated by  $\{1, a, a^*, a^*a\}$  satisfying

$$[a, a^*] = 1 \quad [a^*a, a] = -a \quad [a^*a, a^*] = +a^*.$$

It is easy to see that  $\{1, a, a^*, a^*a\}$  is a basis for the oscillator algebra.

# The Oscillator Algebra

## Definition

Given a Hilbert space  $\mathfrak{H}$  and a closed and densely-defined operator  $a$  on  $\mathfrak{H}$ , the *oscillator algebra* is a Lie algebra generated by  $\{1, a, a^*, a^*a\}$  satisfying

$$[a, a^*] = 1 \quad [a^*a, a] = -a \quad [a^*a, a^*] = +a^*.$$

It is easy to see that  $\{1, a, a^*, a^*a\}$  is a basis for the oscillator algebra.

Given a separable Hilbert space  $\mathfrak{H}$  with orthonormal basis  $\{e_n\}_{n=0}^\infty$ , a representation of the oscillator algebra can be found by defining  $ae_0 = 0$  and  $ae_n = \sqrt{n}e_{n-1}$  for  $n > 0$  so that  $a^*e_n = \sqrt{n+1}e_{n+1}$  and  $a^*ae_n = ne_n$ .

# The Oscillator Algebra


## Definition

Given a Hilbert space  $\mathfrak{H}$  and a closed and densely-defined operator  $a$  on  $\mathfrak{H}$ , the *oscillator algebra* is a Lie algebra generated by  $\{1, a, a^*, a^*a\}$  satisfying

$$[a, a^*] = 1 \quad [a^*a, a] = -a \quad [a^*a, a^*] = +a^*.$$

It is easy to see that  $\{1, a, a^*, a^*a\}$  is a basis for the oscillator algebra.

Given a separable Hilbert space  $\mathfrak{H}$  with orthonormal basis  $\{e_n\}_{n=0}^\infty$ , a representation of the oscillator algebra can be found by defining  $ae_0 = 0$  and  $ae_n = \sqrt{n}e_{n-1}$  for  $n > 0$  so that  $a^*e_n = \sqrt{n+1}e_{n+1}$  and  $a^*ae_n = ne_n$ .

$a^*$  and  $a$  are called the creation and annihilation operators, resp. 

A well-known, concrete representation of the oscillator algebra exists on  $L^2(\mathbb{R}, dx)$ .



A well-known, concrete representation of the oscillator algebra exists on  $L^2(\mathbb{R}, dx)$ .

Define  $af(x) = \frac{1}{\sqrt{2}}(f'(x) + xf(x))$  on functions  $f \in L^2(\mathbb{R}, dx)$  such that  $f', xf \in L^2(\mathbb{R}, dx)$ .

A well-known, concrete representation of the oscillator algebra exists on  $L^2(\mathbb{R}, dx)$ .

Define  $a f(x) = \frac{1}{\sqrt{2}}(f'(x) + x f(x))$  on functions  $f \in L^2(\mathbb{R}, dx)$  such that  $f', x f \in L^2(\mathbb{R}, dx)$ .

On such functions,  $a^*$  has the representation

$$a^* f(x) = \frac{1}{\sqrt{2}}(-f'(x) + x f(x)).$$

A well-known, concrete representation of the oscillator algebra exists on  $L^2(\mathbb{R}, dx)$ .

Define  $a f(x) = \frac{1}{\sqrt{2}}(f'(x) + x f(x))$  on functions  $f \in L^2(\mathbb{R}, dx)$  such that  $f', x f \in L^2(\mathbb{R}, dx)$ .

On such functions,  $a^*$  has the representation

$$a^* f(x) = \frac{1}{\sqrt{2}}(-f'(x) + x f(x)).$$

For nice functions  $f \in L^2(\mathbb{R}, dx)$  (e.g.  $f'', x f', x^2 f \in L^2(\mathbb{R}, dx)$ ),  $a^* a$  then has the representation

$$a^* a f(x) = -\frac{1}{2} f''(x) + \frac{1}{2} x^2 f(x) - \frac{1}{2} f(x).$$

A well-known, concrete representation of the oscillator algebra exists on  $L^2(\mathbb{R}, dx)$ .

Define  $af(x) = \frac{1}{\sqrt{2}}(f'(x) + xf(x))$  on functions  $f \in L^2(\mathbb{R}, dx)$  such that  $f', xf \in L^2(\mathbb{R}, dx)$ .

On such functions,  $a^*$  has the representation

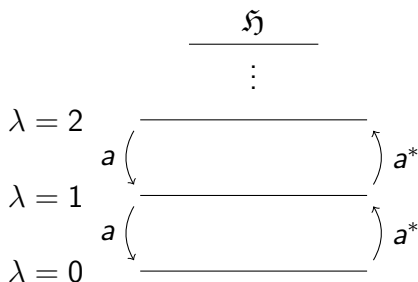
$$a^*f(x) = \frac{1}{\sqrt{2}}(-f'(x) + xf(x)).$$

For nice functions  $f \in L^2(\mathbb{R}, dx)$  (e.g.  $f'', xf', x^2f \in L^2(\mathbb{R}, dx)$ ),  $a^*a$  then has the representation

$$a^*af(x) = -\frac{1}{2}f''(x) + \frac{1}{2}x^2f(x) - \frac{1}{2}f(x).$$

$a^*a$  has eigenvalues  $0, 1, 2, \dots$

Diagrammatically, the action of  $a$  and  $a^*$  can be realized as



The rungs in this ladder represent different eigenfunctions of  $a^*a$  with increasing eigenvalue  $(0, 1, 2, \dots)$ .

# Holomorphic Function Spaces

# Holomorphic Function Spaces

A *holomorphic function space* is a vector space of entire functions equipped with an inner product with respect to which it forms a Hilbert space.

# Holomorphic Function Spaces

A *holomorphic function space* is a vector space of entire functions equipped with an inner product with respect to which it forms a Hilbert space.

Examples include the Bergman space  $A^2(\mathbb{D})$  for some bounded domain  $\mathbb{D} \subset \mathbb{C}$  and the Segal-Bargmann space. The Segal-Bargmann space is the focus of the next section.



# Segal-Bargmann Space

# Segal-Bargmann Space

The connections between the  $L^2(\mathbb{R}, dx)$  representation of the oscillator algebra and  $x$  and  $\frac{d}{dx}$  suggest another potential representation.

# Segal-Bargmann Space

The connections between the  $L^2(\mathbb{R}, dx)$  representation of the oscillator algebra and  $x$  and  $\frac{d}{dx}$  suggest another potential representation.

Note that  $[\frac{d}{dx}, x] = 1$ , however  $x$  is a symmetric operator whereas  $\frac{d}{dx}$  is an anti-symmetric operator, so  $x$  and  $\frac{d}{dx}$  do not—and cannot—form an  $a, a^*$  pair.

# Segal-Bargmann Space

The connections between the  $L^2(\mathbb{R}, dx)$  representation of the oscillator algebra and  $x$  and  $\frac{d}{dx}$  suggest another potential representation.

Note that  $[\frac{d}{dx}, x] = 1$ , however  $x$  is a symmetric operator whereas  $\frac{d}{dx}$  is an anti-symmetric operator, so  $x$  and  $\frac{d}{dx}$  do not—and cannot—form an  $a, a^*$  pair.

Define the measure  $d\rho(z, \bar{z}) = \frac{1}{\pi} e^{-z\bar{z}} dA(z)$  on  $\mathbb{C}$ , where  $dA(x + iy) = dx dy$ .

# Segal-Bargmann Space

The connections between the  $L^2(\mathbb{R}, dx)$  representation of the oscillator algebra and  $x$  and  $\frac{d}{dx}$  suggest another potential representation.

Note that  $[\frac{d}{dx}, x] = 1$ , however  $x$  is a symmetric operator whereas  $\frac{d}{dx}$  is an anti-symmetric operator, so  $x$  and  $\frac{d}{dx}$  do not—and cannot—form an  $a, a^*$  pair.

Define the measure  $d\rho(z, \bar{z}) = \frac{1}{\pi} e^{-z\bar{z}} dA(z)$  on  $\mathbb{C}$ , where  $dA(x + iy) = dx dy$ . On the space of entire functions with finite  $L^2(\mathbb{C}, d\rho)$  norm, the operator  $a f(z) = \frac{d}{dz} f(z)$  has adjoint  $a^*$  given by  $a^* f(z) = z f(z)$ , i.e.

# Segal-Bargmann Space

The connections between the  $L^2(\mathbb{R}, dx)$  representation of the oscillator algebra and  $x$  and  $\frac{d}{dx}$  suggest another potential representation.

Note that  $[\frac{d}{dx}, x] = 1$ , however  $x$  is a symmetric operator whereas  $\frac{d}{dx}$  is an anti-symmetric operator, so  $x$  and  $\frac{d}{dx}$  do not—and cannot—form an  $a, a^*$  pair.

Define the measure  $d\rho(z, \bar{z}) = \frac{1}{\pi} e^{-z\bar{z}} dA(z)$  on  $\mathbb{C}$ , where  $dA(x + iy) = dx dy$ . On the space of entire functions with finite  $L^2(\mathbb{C}, d\rho)$  norm, the operator  $af(z) = \frac{d}{dz}f(z)$  has adjoint  $a^*$  given by  $a^*f(z) = zf(z)$ , i.e.

$$\int_{\mathbb{C}} \frac{d}{dz} f(z) \overline{g(z)} d\rho(z, \bar{z}) = \int_{\mathbb{C}} f(z) \overline{zg(z)} d\rho(z, \bar{z}).$$

# The Bargmann Transform

# The Bargmann Transform

An isomorphism between the  $L^2(\mathbb{R}, dx)$  and  $L^2(\mathbb{C}, d\rho)$  representations of the oscillator algebra exists.



# The Bargmann Transform

An isomorphism between the  $L^2(\mathbb{R}, dx)$  and  $L^2(\mathbb{C}, d\rho)$  representations of the oscillator algebra exists.

The unitary map  $\mathcal{B} : L^2(\mathbb{R}, dx) \rightarrow L^2(\mathbb{C}, d\rho)$  defined by

$$\mathcal{B}f(z) = \frac{1}{\pi^{1/4}} \int_{\mathbb{R}} e^{-z^2/2 + \sqrt{2}xz - x^2/2} f(x) dx$$

# The Bargmann Transform

An isomorphism between the  $L^2(\mathbb{R}, dx)$  and  $L^2(\mathbb{C}, d\rho)$  representations of the oscillator algebra exists.

The unitary map  $\mathcal{B} : L^2(\mathbb{R}, dx) \rightarrow L^2(\mathbb{C}, d\rho)$  defined by

$$\mathcal{B}f(z) = \frac{1}{\pi^{1/4}} \int_{\mathbb{R}} e^{-z^2/2 + \sqrt{2}xz - x^2/2} f(x) dx$$

satisfies

$$\mathcal{B} \left( \frac{1}{\sqrt{2}} \left( \frac{d}{dx} f + xf \right) \right) (z) = \frac{\partial}{\partial z} \mathcal{B}f(z)$$
$$\mathcal{B} \left( \frac{1}{\sqrt{2}} \left( -\frac{d}{dx} f + xf \right) \right) (z) = z \mathcal{B}f(z).$$

# The Bargmann Transform

An isomorphism between the  $L^2(\mathbb{R}, dx)$  and  $L^2(\mathbb{C}, d\rho)$  representations of the oscillator algebra exists.

The unitary map  $\mathcal{B} : L^2(\mathbb{R}, dx) \rightarrow L^2(\mathbb{C}, d\rho)$  defined by

$$\mathcal{B}f(z) = \frac{1}{\pi^{1/4}} \int_{\mathbb{R}} e^{-z^2/2 + \sqrt{2}xz - x^2/2} f(x) dx$$

satisfies

$$\mathcal{B} \left( \frac{1}{\sqrt{2}} \left( \frac{d}{dx} f + xf \right) \right) (z) = \frac{\partial}{\partial z} \mathcal{B}f(z)$$
$$\mathcal{B} \left( \frac{1}{\sqrt{2}} \left( -\frac{d}{dx} f + xf \right) \right) (z) = z \mathcal{B}f(z).$$

The Segal-Bargmann space is denoted by  $\mathcal{O}L^2(\mathbb{C}, d\rho)$ .

The previous relations can be summarized with a pair of commutative diagrams:

The previous relations can be summarized with a pair of commutative diagrams:

$$\begin{array}{ccc} L^2(\mathbb{R}, dx) & \xrightarrow{a} & L^2(\mathbb{R}, dx) \\ \mathcal{B} \downarrow & & \downarrow \mathcal{B} \\ \mathcal{O}L^2(\mathbb{C}, d\rho) & \xrightarrow{a} & \mathcal{O}L^2(\mathbb{C}, d\rho) \end{array}$$

The previous relations can be summarized with a pair of commutative diagrams:

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{a} & L^2(\mathbb{R}, dx) \\
 \mathcal{B} \downarrow & & \downarrow \mathcal{B} \\
 \mathcal{O}L^2(\mathbb{C}, d\rho) & \xrightarrow{a} & \mathcal{O}L^2(\mathbb{C}, d\rho)
 \end{array}$$

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{a^*} & L^2(\mathbb{R}, dx) \\
 \mathcal{B} \downarrow & & \downarrow \mathcal{B} \\
 \mathcal{O}L^2(\mathbb{C}, d\rho) & \xrightarrow{a^*} & \mathcal{O}L^2(\mathbb{C}, d\rho)
 \end{array}$$

# Coupled Supersymmetry

# Coupled Supersymmetry

Let  $\mathfrak{H}_1, \mathfrak{H}_2$  be Hilbert spaces. Let  $a : \mathfrak{H}_1 \rightarrow \mathfrak{H}_2$  be closed and densely-defined; similarly, let  $b : \mathfrak{H}_2 \rightarrow \mathfrak{H}_1$  be closed and densely-defined; and let  $\gamma, \delta \in \mathbb{R}$ .



# Coupled Supersymmetry

Let  $\mathfrak{H}_1, \mathfrak{H}_2$  be Hilbert spaces. Let  $a : \mathfrak{H}_1 \rightarrow \mathfrak{H}_2$  be closed and densely-defined; similarly, let  $b : \mathfrak{H}_2 \rightarrow \mathfrak{H}_1$  be closed and densely-defined; and let  $\gamma, \delta \in \mathbb{R}$ .

## Definition

We say that the ordered quadruplet  $\{a, b, \gamma, \delta\}$  forms a *coupled supersymmetry (coupled SUSY)* if

# Coupled Supersymmetry

Let  $\mathfrak{H}_1, \mathfrak{H}_2$  be Hilbert spaces. Let  $a : \mathfrak{H}_1 \rightarrow \mathfrak{H}_2$  be closed and densely-defined; similarly, let  $b : \mathfrak{H}_2 \rightarrow \mathfrak{H}_1$  be closed and densely-defined; and let  $\gamma, \delta \in \mathbb{R}$ .

## Definition

We say that the ordered quadruplet  $\{a, b, \gamma, \delta\}$  forms a *coupled supersymmetry (coupled SUSY)* if

$$a^*a = bb^* + \gamma 1_{\mathfrak{H}_1}$$

$$aa^* = b^*b + \delta 1_{\mathfrak{H}_2}.$$

# Coupled Supersymmetry

Let  $\mathfrak{H}_1, \mathfrak{H}_2$  be Hilbert spaces. Let  $a : \mathfrak{H}_1 \rightarrow \mathfrak{H}_2$  be closed and densely-defined; similarly, let  $b : \mathfrak{H}_2 \rightarrow \mathfrak{H}_1$  be closed and densely-defined; and let  $\gamma, \delta \in \mathbb{R}$ .

## Definition

We say that the ordered quadruplet  $\{a, b, \gamma, \delta\}$  forms a *coupled supersymmetry (coupled SUSY)* if

$$a^*a = bb^* + \gamma 1_{\mathfrak{H}_1}$$

$$aa^* = b^*b + \delta 1_{\mathfrak{H}_2}.$$

The oscillator algebra is a special case of a coupled SUSY by taking  $b = a$ ,  $\gamma = -1$ , and  $\delta = 1$ .

# Coupled Supersymmetry

Let  $\mathfrak{H}_1, \mathfrak{H}_2$  be Hilbert spaces. Let  $a : \mathfrak{H}_1 \rightarrow \mathfrak{H}_2$  be closed and densely-defined; similarly, let  $b : \mathfrak{H}_2 \rightarrow \mathfrak{H}_1$  be closed and densely-defined; and let  $\gamma, \delta \in \mathbb{R}$ .

## Definition

We say that the ordered quadruplet  $\{a, b, \gamma, \delta\}$  forms a *coupled supersymmetry (coupled SUSY)* if

$$a^*a = bb^* + \gamma 1_{\mathfrak{H}_1}$$

$$aa^* = b^*b + \delta 1_{\mathfrak{H}_2}.$$

The oscillator algebra is a special case of a coupled SUSY by taking  $b = a$ ,  $\gamma = -1$ , and  $\delta = 1$ .

(Some details omitted for brevity.)

Much like the oscillator algebra, Lie algebras underlie coupled SUSYs:

Much like the oscillator algebra, Lie algebras underlie coupled SUSYs:

$$[ba, a^* b^*] = 2(\delta - \gamma) \left( a^* a - \frac{\gamma}{2} \right)$$

$$[a^* a, ba] = -(\delta - \gamma) ba$$

$$[a^* a, a^* b^*] = +(\delta - \gamma) a^* b^*$$

Much like the oscillator algebra, Lie algebras underlie coupled SUSYs:

$$[ba, a^* b^*] = 2(\delta - \gamma) \left( a^* a - \frac{\gamma}{2} \right)$$

$$[a^* a, ba] = -(\delta - \gamma) ba$$

$$[a^* a, a^* b^*] = +(\delta - \gamma) a^* b^*$$

$$[ab, b^* a^*] = 2(\delta - \gamma) \left( aa^* - \frac{\delta}{2} \right)$$

$$[aa^*, ab] = -(\delta - \gamma) ab$$

$$[aa^*, b^* a^*] = +(\delta - \gamma) b^* a^*$$

Much like the oscillator algebra, Lie algebras underlie coupled SUSYs:

$$[ba, a^* b^*] = 2(\delta - \gamma) \left( a^* a - \frac{\gamma}{2} \right)$$

$$[a^* a, ba] = -(\delta - \gamma) ba$$

$$[a^* a, a^* b^*] = +(\delta - \gamma) a^* b^*$$

$$[ab, b^* a^*] = 2(\delta - \gamma) \left( aa^* - \frac{\delta}{2} \right)$$

$$[aa^*, ab] = -(\delta - \gamma) ab$$

$$[aa^*, b^* a^*] = +(\delta - \gamma) b^* a^*$$

These Lie algebras are isomorphic to the  $\mathfrak{su}(1, 1)$  Lie algebra.



Much like the oscillator algebra, Lie algebras underlie coupled SUSYs:

$$[ba, a^* b^*] = 2(\delta - \gamma) \left( a^* a - \frac{\gamma}{2} \right)$$

$$[a^* a, ba] = -(\delta - \gamma) ba$$

$$[a^* a, a^* b^*] = +(\delta - \gamma) a^* b^*$$

$$[ab, b^* a^*] = 2(\delta - \gamma) \left( aa^* - \frac{\delta}{2} \right)$$

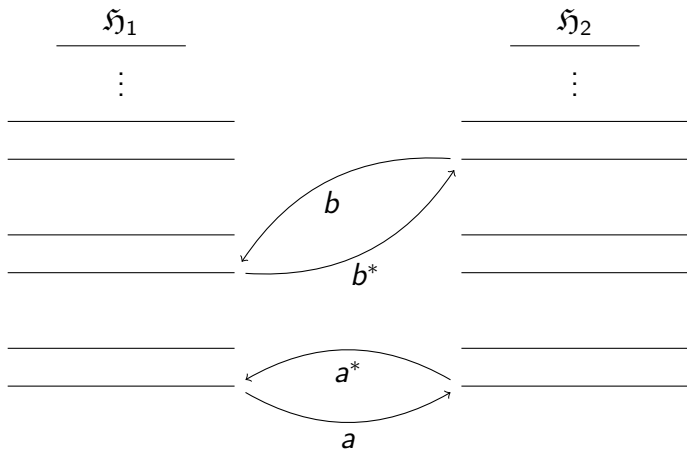
$$[aa^*, ab] = -(\delta - \gamma) ab$$

$$[aa^*, b^* a^*] = +(\delta - \gamma) b^* a^*$$

These Lie algebras are isomorphic to the  $\mathfrak{su}(1, 1)$  Lie algebra.

$a^* a$  has eigenvalues  $m(\delta - \gamma)$ ,  $m(\delta - \gamma) + \delta$ .

Diagrammatically, the coupled SUSY relations can be viewed as



Here the eigenfunctions are eigenfunctions of  $a^*a$  and  $aa^*$ .

A concrete realization of a coupled SUSY is given by the operators  $a$  and  $b$  on  $\mathfrak{H}_1 = L^2(\mathbb{R}, dx) = \mathfrak{H}_2$  defined by

A concrete realization of a coupled SUSY is given by the operators  $a$  and  $b$  on  $\mathfrak{H}_1 = L^2(\mathbb{R}, dx) = \mathfrak{H}_2$  defined by

$$af(x) = \frac{1}{\sqrt{2}} \left( \frac{1}{x^{n-1}} \frac{d}{dx} + x^n \right) f(x)$$

$$bf(x) = \frac{1}{\sqrt{2}} \left( \frac{d}{dx} \frac{1}{x^{n-1}} + x^n \right) f(x)$$

with  $\gamma = -1$  and  $\delta = 2n - 1$  where  $n \in \mathbb{N}$ .

A concrete realization of a coupled SUSY is given by the operators  $a$  and  $b$  on  $\mathfrak{H}_1 = L^2(\mathbb{R}, dx) = \mathfrak{H}_2$  defined by

$$af(x) = \frac{1}{\sqrt{2}} \left( \frac{1}{x^{n-1}} \frac{d}{dx} + x^n \right) f(x)$$
$$bf(x) = \frac{1}{\sqrt{2}} \left( \frac{d}{dx} \frac{1}{x^{n-1}} + x^n \right) f(x)$$

with  $\gamma = -1$  and  $\delta = 2n - 1$  where  $n \in \mathbb{N}$ .

The  $L^2(\mathbb{R}, dx)$  representation of the oscillator algebra arises as a special case by taking  $n = 1$  in which case  $b = a$ ,  $\gamma = -1$ , and  $\delta = 1$ .

Holomorphic representations also exist for  $\gamma = -1$  and  $\delta = 2n - 1$ . Define the holomorphic function spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  to be

Holomorphic representations also exist for  $\gamma = -1$  and  $\delta = 2n - 1$ . Define the holomorphic function spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  to be

$$\mathfrak{F}_1^{(n)} = \overline{\text{span}}\{1, z^{2n-1}, z^{2n}, z^{4n-1}, z^{4n}, \dots\}$$

$$\mathfrak{F}_2^{(n)} = \overline{\text{span}}\{z^{n-1}, z^n, z^{3n-1}, z^{3n}, \dots\}$$

Holomorphic representations also exist for  $\gamma = -1$  and  $\delta = 2n - 1$ . Define the holomorphic function spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  to be

$$\mathfrak{F}_1^{(n)} = \overline{\text{span}}\{1, z^{2n-1}, z^{2n}, z^{4n-1}, z^{4n}, \dots\}$$

$$\mathfrak{F}_2^{(n)} = \overline{\text{span}}\{z^{n-1}, z^n, z^{3n-1}, z^{3n}, \dots\}$$

For  $\mathfrak{F}_1^{(n)}$ , the closure is taken with respect to the norm generated from the inner product with measure  $\rho_1^{(n)}(z, \bar{z}) = \lambda_n (z\bar{z})^{n-\frac{1}{2}} K_{1-\frac{1}{2n}} \left( \frac{(z\bar{z})^n}{n} \right)$ .



Holomorphic representations also exist for  $\gamma = -1$  and  $\delta = 2n - 1$ . Define the holomorphic function spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  to be

$$\mathfrak{F}_1^{(n)} = \overline{\text{span}}\{1, z^{2n-1}, z^{2n}, z^{4n-1}, z^{4n}, \dots\}$$

$$\mathfrak{F}_2^{(n)} = \overline{\text{span}}\{z^{n-1}, z^n, z^{3n-1}, z^{3n}, \dots\}$$

For  $\mathfrak{F}_1^{(n)}$ , the closure is taken with respect to the norm generated from the inner product with measure  $\rho_1^{(n)}(z, \bar{z}) = \lambda_n (z\bar{z})^{n-\frac{1}{2}} K_{1-\frac{1}{2n}} \left( \frac{(z\bar{z})^n}{n} \right)$ .

For  $\mathfrak{F}_2^{(n)}$ , the closure is taken with respect to the norm generated from the inner product with measure  $\rho_2^{(n)}(z, \bar{z}) = \lambda_n (z\bar{z})^{n-\frac{1}{2}} K_{-\frac{1}{2n}} \left( \frac{(z\bar{z})^n}{n} \right)$ .

Holomorphic representations also exist for  $\gamma = -1$  and  $\delta = 2n - 1$ . Define the holomorphic function spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  to be

$$\begin{aligned}\mathfrak{F}_1^{(n)} &= \overline{\text{span}}\{1, z^{2n-1}, z^{2n}, z^{4n-1}, z^{4n}, \dots\} \\ \mathfrak{F}_2^{(n)} &= \overline{\text{span}}\{z^{n-1}, z^n, z^{3n-1}, z^{3n}, \dots\}\end{aligned}$$

For  $\mathfrak{F}_1^{(n)}$ , the closure is taken with respect to the norm generated from the inner product with measure  $\rho_1^{(n)}(z, \bar{z}) = \lambda_n (z\bar{z})^{n-\frac{1}{2}} K_{1-\frac{1}{2n}} \left( \frac{(z\bar{z})^n}{n} \right)$ .

For  $\mathfrak{F}_2^{(n)}$ , the closure is taken with respect to the norm generated from the inner product with measure  $\rho_2^{(n)}(z, \bar{z}) = \lambda_n (z\bar{z})^{n-\frac{1}{2}} K_{-\frac{1}{2n}} \left( \frac{(z\bar{z})^n}{n} \right)$ .

With respect to the above measures, it is clear that  $z^m$  and  $z^n$  are orthogonal if  $n \neq m$  since the measures are radial.

For these spaces,  $a_n : \mathfrak{F}_1^{(n)} \rightarrow \mathfrak{F}_2^{(n)}$  given by  $a_n f(z) = \frac{1}{z^{n-1}} \frac{d}{dz} f(z)$

For these spaces,  $a_n : \mathfrak{F}_1^{(n)} \rightarrow \mathfrak{F}_2^{(n)}$  given by  $a_n f(z) = \frac{1}{z^{n-1}} \frac{d}{dz} f(z)$   
has adjoint  $a_n^* : \mathfrak{F}_2^{(n)} \rightarrow \mathfrak{F}_1^{(n)}$  given by  $a_n^* f(z) = z^n f(z)$ .

For these spaces,  $a_n : \mathfrak{F}_1^{(n)} \rightarrow \mathfrak{F}_2^{(n)}$  given by  $a_n f(z) = \frac{1}{z^{n-1}} \frac{d}{dz} f(z)$   
has adjoint  $a_n^* : \mathfrak{F}_2^{(n)} \rightarrow \mathfrak{F}_1^{(n)}$  given by  $a_n^* f(z) = z^n f(z)$ .

And likewise  $b_n : \mathfrak{F}_2^{(n)} \rightarrow \mathfrak{F}_1^{(n)}$  given by  $b_n f(z) = \frac{d}{dz} \frac{1}{z^{n-1}} f(z)$

For these spaces,  $a_n : \mathfrak{F}_1^{(n)} \rightarrow \mathfrak{F}_2^{(n)}$  given by  $a_n f(z) = \frac{1}{z^{n-1}} \frac{d}{dz} f(z)$   
has adjoint  $a_n^* : \mathfrak{F}_2^{(n)} \rightarrow \mathfrak{F}_1^{(n)}$  given by  $a_n^* f(z) = z^n f(z)$ .

And likewise  $b_n : \mathfrak{F}_2^{(n)} \rightarrow \mathfrak{F}_1^{(n)}$  given by  $b_n f(z) = \frac{d}{dz} \frac{1}{z^{n-1}} f(z)$   
has adjoint  $b_n^* : \mathfrak{F}_1^{(n)} \rightarrow \mathfrak{F}_2^{(n)}$  given by  $b_n^* f(z) = z^n f(z)$ .

# Bargmann Transforms for Coupled SUSYs

# Bargmann Transforms for Coupled SUSYs

Bargmann transforms can be defined for these coupled SUSYs connecting the  $L^2(\mathbb{R})$  and holomorphic representations.



# Bargmann Transforms for Coupled SUSYs

Bargmann transforms can be defined for these coupled SUSYs connecting the  $L^2(\mathbb{R})$  and holomorphic representations.

Let  $\mathcal{B}_1^{(n)} : L^2(\mathbb{R}, dx) \rightarrow \mathfrak{F}_1^{(n)}$  and  $\mathcal{B}_2^{(n)} : L^2(\mathbb{R}, dx) \rightarrow \mathfrak{F}_2^{(n)}$  denote the as-of-yet defined Bargmann transforms. The Bargmann transforms can be defined by the following commutative diagrams.

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{a_n} & L^2(\mathbb{R}, dx) \\
 \mathcal{B}_1^{(n)} \downarrow & & \mathcal{B}_2^{(n)} \downarrow \\
 \mathfrak{F}_1^{(n)} & \xrightarrow{a_n} & \mathfrak{F}_2^{(n)}
 \end{array}$$

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{a_n} & L^2(\mathbb{R}, dx) \\
 \mathcal{B}_1^{(n)} \downarrow & & \mathcal{B}_2^{(n)} \downarrow \\
 \mathfrak{F}_1^{(n)} & \xrightarrow{a_n} & \mathfrak{F}_2^{(n)}
 \end{array}$$

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{b_n} & L^2(\mathbb{R}, dx) \\
 \mathcal{B}_2^{(n)} \downarrow & & \mathcal{B}_1^{(n)} \downarrow \\
 \mathfrak{F}_2^{(n)} & \xrightarrow{b_n} & \mathfrak{F}_1^{(n)}
 \end{array}$$

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{a_n^*} & L^2(\mathbb{R}, dx) \\
 \mathcal{B}_2^{(n)} \downarrow & & \mathcal{B}_1^{(n)} \downarrow \\
 \tilde{\mathfrak{F}}_2^{(n)} & \xrightarrow{a_n^*} & \tilde{\mathfrak{F}}_1^{(n)}
 \end{array}$$

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{a_n^*} & L^2(\mathbb{R}, dx) \\
 \mathcal{B}_2^{(n)} \downarrow & & \mathcal{B}_1^{(n)} \downarrow \\
 \mathfrak{F}_2^{(n)} & \xrightarrow{a_n^*} & \mathfrak{F}_1^{(n)}
 \end{array}$$

$$\begin{array}{ccc}
 L^2(\mathbb{R}, dx) & \xrightarrow{b_n^*} & L^2(\mathbb{R}, dx) \\
 \mathcal{B}_1^{(n)} \downarrow & & \mathcal{B}_2^{(n)} \downarrow \\
 \mathfrak{F}_1^{(n)} & \xrightarrow{b_n^*} & \mathfrak{F}_2^{(n)}
 \end{array}$$

The Bargmann transforms can be realized as integral transforms. From the commutative diagrams, the integral kernels satisfy intertwined partial differential equations:

The Bargmann transforms can be realized as integral transforms. From the commutative diagrams, the integral kernels satisfy intertwined partial differential equations:

$$\frac{1}{z^{n-1}} \frac{\partial}{\partial z} A_1^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{\partial}{\partial x} \frac{1}{x^{n-1}} + x^n \right) A_2^{(n)}(z, x)$$

The Bargmann transforms can be realized as integral transforms. From the commutative diagrams, the integral kernels satisfy intertwined partial differential equations:

$$\frac{1}{z^{n-1}} \frac{\partial}{\partial z} A_1^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{\partial}{\partial x} \frac{1}{x^{n-1}} + x^n \right) A_2^{(n)}(z, x)$$

$$\frac{\partial}{\partial z} \frac{1}{z^{n-1}} A_2^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{1}{x^{n-1}} \frac{\partial}{\partial x} + x^n \right) A_1^{(n)}(z, x)$$



The Bargmann transforms can be realized as integral transforms. From the commutative diagrams, the integral kernels satisfy intertwined partial differential equations:

$$\frac{1}{z^{n-1}} \frac{\partial}{\partial z} A_1^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{\partial}{\partial x} \frac{1}{x^{n-1}} + x^n \right) A_2^{(n)}(z, x)$$

$$\frac{\partial}{\partial z} \frac{1}{z^{n-1}} A_2^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{1}{x^{n-1}} \frac{\partial}{\partial x} + x^n \right) A_1^{(n)}(z, x)$$

$$z^n A_1^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{\partial}{\partial x} \frac{1}{x^{n-1}} + x^n \right) A_2^{(n)}(z, x)$$

The Bargmann transforms can be realized as integral transforms. From the commutative diagrams, the integral kernels satisfy intertwined partial differential equations:

$$\frac{1}{z^{n-1}} \frac{\partial}{\partial z} A_1^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{\partial}{\partial x} \frac{1}{x^{n-1}} + x^n \right) A_2^{(n)}(z, x)$$

$$\frac{\partial}{\partial z} \frac{1}{z^{n-1}} A_2^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{1}{x^{n-1}} \frac{\partial}{\partial x} + x^n \right) A_1^{(n)}(z, x)$$

$$z^n A_1^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{\partial}{\partial x} \frac{1}{x^{n-1}} + x^n \right) A_2^{(n)}(z, x)$$

$$z^n A_2^{(n)}(z, x) = \frac{1}{\sqrt{2}} \left( -\frac{1}{x^{n-1}} \frac{\partial}{\partial x} + x^n \right) A_1^{(n)}(z, x)$$

Then (after  $x$  pages of algebra)  $\mathcal{B}_1^{(n)}$  and  $\mathcal{B}_2^{(n)}$  are given by the integral transforms

Then (after  $x$  pages of algebra)  $\mathcal{B}_1^{(n)}$  and  $\mathcal{B}_2^{(n)}$  are given by the integral transforms

$$\mathcal{B}_1^{(n)} f(z) = \tau_n \int_{\mathbb{R}} e^{-z^{2n}/2n} (-zx)^{n-\frac{1}{2}} K_{1-\frac{1}{2n}} \left( \sqrt{2} \frac{(-zx)^n}{n} \right) e^{-x^{2n}/2n} f(x) dx$$

$$\mathcal{B}_2^{(n)} g(z) = \tau_n \int_{\mathbb{R}} e^{-z^{2n}/2n} (-zx)^{n-\frac{1}{2}} K_{-\frac{1}{2n}} \left( \sqrt{2} \frac{(-zx)^n}{n} \right) e^{-x^{2n}/2n} g(x) dx$$

When  $n = 1$ , these reduce exactly to the usual Bargmann transform associated to the oscillator algebra.

The generalized Segal-Bargmann spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  are also reproducing kernel Hilbert spaces. The reproducing kernels are established in the usual way

The generalized Segal-Bargmann spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  are also reproducing kernel Hilbert spaces. The reproducing kernels are established in the usual way

$$K_1^{(n)}(z; w) = \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl} w^{2nl}}{\|z^{2nl}\|^2} + \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+2n-1} w^{2nl+2n-1}}{\|z^{2nl+2n-1}\|^2}$$

$$K_2^{(n)}(z; w) = \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+n-1} w^{2nl+n-1}}{\|z^{2nl+n-1}\|^2} + \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+n} w^{2nl+n}}{\|z^{2nl+n}\|^2}$$

The generalized Segal-Bargmann spaces  $\tilde{\mathfrak{F}}_1^{(n)}$  and  $\tilde{\mathfrak{F}}_2^{(n)}$  are also reproducing kernel Hilbert spaces. The reproducing kernels are established in the usual way

$$K_1^{(n)}(z; w) = \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl} w^{2nl}}{\|z^{2nl}\|^2} + \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+2n-1} w^{2nl+2n-1}}{\|z^{2nl+2n-1}\|^2}$$

$$K_2^{(n)}(z; w) = \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+n-1} w^{2nl+n-1}}{\|z^{2nl+n-1}\|^2} + \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+n} w^{2nl+n}}{\|z^{2nl+n}\|^2}$$

so that if  $f \in \tilde{\mathfrak{F}}_1^{(n)}$ ,  $g \in \tilde{\mathfrak{F}}_2^{(n)}$ , then

The generalized Segal-Bargmann spaces  $\mathfrak{F}_1^{(n)}$  and  $\mathfrak{F}_2^{(n)}$  are also reproducing kernel Hilbert spaces. The reproducing kernels are established in the usual way

$$K_1^{(n)}(z; w) = \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl} w^{2nl}}{\|z^{2nl}\|^2} + \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+2n-1} w^{2nl+2n-1}}{\|z^{2nl+2n-1}\|^2}$$

$$K_2^{(n)}(z; w) = \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+n-1} w^{2nl+n-1}}{\|z^{2nl+n-1}\|^2} + \sum_{l=0}^{\infty} \frac{\bar{z}^{2nl+n} w^{2nl+n}}{\|z^{2nl+n}\|^2}$$

so that if  $f \in \mathfrak{F}_1^{(n)}$ ,  $g \in \mathfrak{F}_2^{(n)}$ , then

$$f(z) = \int_{\mathbb{C}} f(w) \overline{K_1^{(n)}(z; w)} d\rho_1^{(n)}(w, \bar{w})$$

$$g(z) = \int_{\mathbb{C}} g(w) \overline{K_2^{(n)}(z; w)} d\rho_2^{(n)}(w, \bar{w}).$$



# Applications

# Applications

The Bargmann transform can be used for the analysis of the backwards heat equation. In similar vein, the coupled SUSY Bargmann transforms could be used to analyze similar PDEs with the Laplacian replaced by a different second order differential operator.

# Applications

The Bargmann transform can be used for the analysis of the backwards heat equation. In similar vein, the coupled SUSY Bargmann transforms could be used to analyze similar PDEs with the Laplacian replaced by a different second order differential operator.

The Bargmann transform is also effectively a short-time Fourier transform with Gaussian window, i.e. a time-frequency transform. Leaving the Fourier setting, short-time transforms become somewhat mysterious and non-unique. The coupled SUSY Bargmann transforms may be yet another short-time transform, distinct from one Williams, *et al* posited previously.

## Further Questions

What happens if one takes non-integer  $n$ ? Do integrals over  $\mathbb{C}$  need to be replaced by integrals over infinite sectors? Can unitarity of the coupled SUSY Bargmann transforms still be attained?

## Further Questions

What happens if one takes non-integer  $n$ ? Do integrals over  $\mathbb{C}$  need to be replaced by integrals over infinite sectors? Can unitarity of the coupled SUSY Bargmann transforms still be attained?

What play the analogous roles of the translation and modulation operator for the coupled SUSY Bargmann transforms? Is there an analogue of the Heisenberg group for these?

## Further Questions

What happens if one takes non-integer  $n$ ? Do integrals over  $\mathbb{C}$  need to be replaced by integrals over infinite sectors? Can unitarity of the coupled SUSY Bargmann transforms still be attained?

What play the analogous roles of the translation and modulation operator for the coupled SUSY Bargmann transforms? Is there an analogue of the Heisenberg group for these?

Can the distribution theory of the coupled SUSY Bargmann transforms be used to analyze worse distributions, e.g.  $e^{x^2}$ ? If so, can they be employed for the analysis of the backwards heat equation?

## Further Questions

What happens if one takes non-integer  $n$ ? Do integrals over  $\mathbb{C}$  need to be replaced by integrals over infinite sectors? Can unitarity of the coupled SUSY Bargmann transforms still be attained?

What play the analogous roles of the translation and modulation operator for the coupled SUSY Bargmann transforms? Is there an analogue of the Heisenberg group for these?

Can the distribution theory of the coupled SUSY Bargmann transforms be used to analyze worse distributions, e.g.  $e^{x^2}$ ? If so, can they be employed for the analysis of the backwards heat equation?

Are there peculiarities of the Fock spaces associated to coupled SUSYs? What happens if one combines Fock spaces for different values of  $n$ ?

Thank you