# Quantum symmetric states on free product C*-algebras 

Claus Köstler University College Cork Joint work with<br>Ken Dykema \& John Williams<br>arXiv:1305.7293

Focus Program on Noncommutative Distributions in Free Probability
Workshop on Combinatorial and Random Matrix Aspects of Noncommutative Distributions and Free Probability

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## Introduction and Motivation

## Classical Probability



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exchangeability $\longleftrightarrow$ independence $\longleftrightarrow$| symmetries of |
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Free Probability

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$\rightsquigarrow$ New direction of research in free probability

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Random variables $X_{1}, X_{2}, \ldots \subset L^{\infty}(\mathbb{A})$ are exchangeable if

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\varphi\left(X_{i_{1}} \cdots X_{i_{n}}\right)=\varphi\left(X_{\pi\left(i_{1}\right)} \cdots X_{\pi\left(i_{n}\right)}\right)
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$\equiv$ exchangeable sequence of $0-1$-valued random variables

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- ...dynamical co-symmetry: there is a *-homomorphism $\alpha: L^{\infty}(\mathbb{A}) \rightarrow C\left(\mathbb{S}_{k}\right) \otimes L^{\infty}(\mathbb{A})$ such that

$$
\begin{array}{rlrl}
\text { id } \otimes \varphi \circ \alpha(\bullet) & =\varphi(\bullet) \mathbb{1} & \quad \text { (invariance } \\
\alpha\left(X_{i}\right) & =\sum_{j=1}^{k} e(\pi)_{i j} \otimes X_{j}
\end{array}
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for all permutation matrices $e(\pi)$ with $e(\pi)_{i j}=\delta_{m(i)}$

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Exchangeability as a co-symmetry can be 'quantized': $\diamond$ Replace $X_{k}$ 's by operators $x_{1}, x_{2}, \ldots$

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Key idea (K. \& Speicher 2008):
Exchangeability as a co-symmetry can be 'quantized':
$\diamond$ Replace $X_{k}$ 's by operators $x_{1}, x_{2}, \ldots$
$\diamond$ Replace $C\left(\mathbb{S}_{k}\right)$ by quantum permutation group $A_{s}(k)$

## Quantum Permutation Groups

Definition and Theorem (Wang 1998)
The quantum permutation group $A_{s}(k)$ is the universal unital
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The abelianization of $A_{s}(k)$ is $C\left(\mathbb{S}_{k}\right)$, the continuous functions on the symmetric group $\mathbb{S}_{k}$

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Consider a $\mathrm{W}^{*}$-probability space $(\mathcal{A}, \varphi)$. A sequence of operators $x_{1}, x_{2}, \ldots \subset \mathcal{A}$ is quantum exchangeable if its distribution is invariant under the coaction of quantum permutations:

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\varphi\left(x_{i_{1}} \cdots x_{i_{n}}\right) \mathbb{1}_{A_{s}(k)}=\sum_{j_{1}, \ldots, j_{n}=1}^{k} e_{i_{1} j_{1}} \cdots e_{i_{n} j_{n}} \varphi\left(x_{j_{1}} \cdots x_{j_{n}}\right)
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Remark
quantum exchangeability $\nRightarrow$ exchangeability

## A free analogue of de Finetti's theorem

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The following are equivalent for an infinite sequence of noncommutative random variables $x_{1}, x_{2}, \ldots$ in a $W^{*}$-probability space $(\mathcal{A}, \varphi)$ :

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(a) the sequence is quantum exchangeable
(b) the sequence is identically distributed and freely independent with amalgamation over $\mathcal{T}$

Here $\mathcal{T}$ denotes the tail von Neumann algebra

$$
\mathcal{T}=\bigcap_{n \in \mathbb{N}} v \mathrm{~N}\left(x_{k} \mid k \geq n\right)
$$

## Quantum symmetric states

## Definition (Dykema \& K \& Williams)

Let $A$ be a unital $C^{*}$-algebra, $(\mathfrak{B}, \phi)$ a $C^{*}$-probability space and $\lambda_{i}: \mathfrak{A} \rightarrow \mathfrak{B}$ unital ${ }^{*}$-homomorphisms $(i \in \mathbb{N})$. The sequence $\left(\lambda_{i}\right)_{i}$ is said to be quantum exchangeable w.r.t. $\psi$ or the state $\psi$ is said to be quantum symmetric w.r.t. $\left(\lambda_{i}\right)_{i}$ if

$$
\begin{array}{r}
\psi\left(\lambda_{i_{1}}\left(a_{1}\right) \cdots \lambda_{i_{n}}\left(a_{n}\right)\right) \mathbb{1}_{A_{s}(k)} \\
=\sum_{j_{1}, \ldots, j_{n}=1}^{k} e_{i_{1} j_{1}} \cdots e_{i_{n} j_{n}} \psi\left(\lambda_{j_{1}}\left(a_{1}\right) \cdots \lambda_{j_{n}}\left(a_{n}\right)\right) \mathbb{1}_{A_{s}(k)}
\end{array}
$$

for all $n \in \mathbb{N}$, all $i_{1}, \ldots i_{n} \in\{1, \ldots, k\}$, all $a_{1}, \ldots a_{n} \in A$, all $k \times k$-matrices $\left(e_{i j}\right)_{i j}$ satisfying defining relations for $A_{s}(k)$.

## Standing notation - Goal

Throughout $A$ is a unital $C^{*}$-algebra and $\mathfrak{A}:=*_{1}^{\infty} A$ denotes the universal, unital free product of infinitely many copies of $A$.

QSS $(A)$ denotes the set of all quantum symmetric states on $\mathfrak{A}$.
TQSS $(A)$ denotes the set of all tracial quantum symmetric states on $\mathfrak{A}$.

## Proposition

For any unital $C^{*}$-algebra $A$ the sets $\operatorname{QSS}(A)$ and $\operatorname{TQSS}(A)$ are compact, convex subsets of the Banach space dual of $\mathfrak{A}$ in the weak*-topology.

Goal
Study and characterize $\operatorname{QSS}(A)$ and $\operatorname{TQSS}(A)$ as far as possible.

## What is the free probability counterpart to the following 'classical' result?

## Theorem (Størmer 1969)

Let $A$ be a unital $C^{*}$-algebra and $\otimes_{1}^{\infty} \mathrm{A}$ the infinite minimal tensor product of $A$. Then the set of all symmetric states on $\otimes_{1}^{\infty} A$ is a Choquet simplex and an extreme symmetric state on $\otimes_{1}^{\infty} A$ is an infinite tensor product state of the form $\otimes_{1}^{\infty} \psi$, with $\psi \in \mathrm{S}(A)$.

## Remark

So, for a given unital $C^{*}$-algebra, one has a bijective correspondence between symmetric states on $\bigotimes_{1}^{\infty} A$ and probability measures on $S(A)$.

## Quantum symmetric states arising from freeness

## Proposition

Let $\mathfrak{B}$ be a unital C*-algebra and $^{\mathfrak{D}} \subseteq \mathfrak{B}$ be unital C*-subalgebra $^{*}$ with conditional expectation $E: \mathfrak{B} \rightarrow \mathfrak{D}$ (i.e. a projection of norm one onto $\mathfrak{D}$ ). Suppose that

$$
\pi_{i}: A \rightarrow \mathfrak{B} \quad(i \in \mathbb{N})
$$

are ${ }^{*}$-homorphisms such that $E \circ \pi_{i}$ is the same for all $i$ and that $\left(\pi_{i}(A)\right)_{i=1}^{\infty}$ is free with respect to $E$. Let $\pi=*_{i=1}^{\infty} \pi_{i}: \mathfrak{A} \rightarrow \mathfrak{B}$ be the resulting free product *-homomorphism. For a state $\rho$ on $\mathfrak{D}$ consider the state $\psi=\rho \circ E \circ \pi$. Then $\psi$ is quantum symmetric.

## Remark

The proof uses Speicher's free $\mathfrak{D}$-valued cumulants and the defining properties of the projections $e_{i j}$ from Wang's quantum permutation groups.

## Tail algebras of symmetric states

Let $\psi$ be a state on $\mathfrak{A}=*_{1}^{\infty} A$. Passing to the GNS representation $\left(\mathcal{H}_{\psi}, \pi_{\psi}, \Omega_{\psi}\right)$ of $(\mathfrak{A}, \psi)$, put

$$
\mathcal{M}_{\psi}=\pi_{\psi}(\mathfrak{A})^{\prime \prime} \quad \hat{\psi}:=\left\langle\Omega_{\psi}, \bullet \Omega_{\psi}\right\rangle
$$

The tail algebra of $\psi$ is the von Neumann subalgebra

$$
\mathcal{T}_{\psi}=\bigcap_{n=1}^{\infty} W^{*}\left(\bigcup_{i=n}^{\infty} \pi_{\psi}\left(\lambda_{i}(A)\right)\right) \subset \mathcal{M}_{\psi}
$$

## Proposition (Dykema \& K \& Williams)

Suppose $\psi$ is symmetric. Then there exists a normal $\hat{\psi}$-preserving conditional expection $E_{\psi}$ from $\mathcal{M}_{\psi}$ onto $\mathcal{T}_{\psi}$.

## 'Quantum symmetric noncommutative distributions' imply freeness with amalgamation

## Theorem (Dykema \& K \& Williams)

Let $\psi$ be a quantum symmetric state on $\mathfrak{A}=*_{1}^{\infty} A$ and put

$$
\mathcal{B}_{i}:=W^{*}\left(\pi_{\psi}\left(\lambda_{i}(A)\right) \cup \mathcal{T}_{\psi}\right) .
$$

Then $\left(\mathcal{B}_{i}\right)_{i=1}^{\infty}$ is free with respect to $E_{\psi}$.

## Remark

- Our proof is modeled along the proof of K \& Speicher (2009), but now starts in an $\mathrm{C}^{*}$-algebraic setting and does not assume the faithfulness of states.
- Curran's approach (2009) considers the more delicate situation of quantum exchangeability of finite sequences in a *-algebraic setting and obtains the result for infinite sequences as a limiting



## What tail algebras can appear?

In Størmer's setting of symmetric states on $\bigotimes_{1}^{\infty} A$, only abelian tail algebras can arise. But in our setting of quantum symmetric states one has:

Theorem (Dykema \& K '2012)
Let $\mathcal{N}$ be a countable generated von Neumann algebra. Then there exists a unital $C^{*}$-algebra $A$ and a quantum symmetric state $\psi$ on $\mathfrak{A}=*_{1}^{\infty} A$ such that $\mathcal{T}_{\psi} \simeq \mathcal{N}$.

## Remark

Størmer's approach does not use tail algebras; the machinery of ergodic decomposition of states and Choquet theory is available.

## Description of quantum symmetric states $\operatorname{QSS}(A)$

For a unital $C^{*}$-algebra $A$, let $\mathcal{V}(A)$ be the set (of all equivalence classes) of quintuples ( $\mathcal{B}, \mathcal{D}, E, \sigma, \rho$ ) such that
(i) $\mathcal{B}$ is a von Neumann algebra,
(ii) $\mathcal{D}$ is a unital von Neumann subalgebra of $\mathcal{B}$,
(iii) $E: \mathcal{B} \rightarrow \mathcal{D}$ is a normal conditional expectation onto $\mathcal{D}$,
(iv) $\sigma: A \rightarrow \mathcal{B}$ is a unital *-homomorphism,
(v) $\rho$ is a normal state on $\mathcal{D}$,
(vi) the GNS representation of $\rho \circ E$ is a faithful represent. of $\mathcal{B}$;
(vii) $\mathcal{B}=W^{*}(\sigma(A) \cup \mathcal{D})$,
(viii) $\mathcal{D}$ is the smallest unital von Neumann subalgebra of $\mathcal{B}$ that satisfies

$$
E\left(d_{0} \sigma\left(a_{1}\right) d_{1} \cdots \sigma\left(a_{n}\right) d_{n}\right) \in \mathcal{D}
$$



## Theorem (Dykema \& K \& Williams)

There is a bijection $\mathcal{V}(A) \rightarrow \operatorname{QSS}(A)$ that assigns to $W=(\mathcal{B}, \mathcal{D}, E, \sigma, \rho) \in \mathcal{V}(A)$ the quantum symmetric state $\psi=\psi_{W}$ as follows. Let

$$
(\mathcal{M}, \widetilde{E})=\left(*_{\mathcal{D}}\right)_{i=1}^{\infty}(\mathcal{B}, E)
$$

be the amalgamated free product of von Neumann algebras, let

$$
\pi_{i}: A \xrightarrow{\sigma} \mathcal{B} \xrightarrow{i \text {-th comp }} \mathcal{M}, \quad \pi:=*_{i=1}^{\infty} \pi_{i}: \mathfrak{A} \rightarrow \mathcal{M}
$$

free product *-homomorphism
and set $\psi=\rho \circ \widetilde{E} \circ \pi$. Under this correspondence, the following identifications of objects and resulting constructions can be made:

| from GNS construction | $\mathcal{T}_{\psi}$ | $\mathcal{M}_{\psi}$ | $\pi_{\psi}$ | $\hat{\psi}$ | $E_{\psi}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| from quintuple $W$ | D | $\mathcal{M}$ |  | $\rho \circ \tilde{E}$ | $\tilde{E}$ |  |

## N-pure states on von Neumann algebras

A state $\psi$ on a $C^{*}$-algebra is said to be pure if whenever $\rho$ is a state on this $C^{*}$-algebra with $t \rho \leq \psi$ for some $0<t<1$, it follows $\rho=\psi$.

## Proposition

Let $\rho$ be a normal state on the von Neumann algebra $\mathcal{D}$. TFAE:
(i) the support projection of $\rho$ is a minimal projection in $\mathcal{D}$,
(ii) $\rho$ is pure.

To emphasize this support property, a normal pure state $\rho$ on $\mathcal{D}$ is called an $\mathbf{n}$-pure state.

## Remark

Von Neumann algebras without discrete type I parts possess no n-pure states, but such von Neumann algebras may appear as tail algebra for a quantum symmetric state.

## Extreme quantum symmetric states

Theorem (Dykema \& K \& Williams)
Let $\psi \in \operatorname{QSS}(A)$ and let $W=(\mathcal{B}, \mathcal{D}, E, \sigma, \rho) \in \mathcal{V}(A)$ be the quintuple corresponding to $\psi$ (under the bijection as indicated in the previous theorem).

$$
\psi \in \partial_{e} \operatorname{QSS}(A) \quad \Longleftrightarrow \quad \rho \text { is an n-pure state on } \mathcal{D} \text {. }
$$

## Remark

- Though having an n-pure state is a restriction on the tail algebra and forces it to have a discrete type I part, the tail algebra can still be quite complicated.
- For various examples of quantum symmetric states with 'exotic tail algebras' see our preprint.


## Central quantum symmetric states

Recall that $\mathfrak{A}=*_{1}^{\infty} A$.
Notation
$\operatorname{ZQSS}(A):=\left\{\psi \in \operatorname{QSS}(A) \mid \mathcal{T}_{\psi} \subset \mathcal{Z}\left(\pi_{\psi}(\mathfrak{A})^{\prime \prime}\right)\right\}$
$\operatorname{ZTQSS}(A):=\operatorname{TQSS}(A) \cap \operatorname{ZQSS}(A)$
Theorem (Dykema \& K \& Williams)
$\operatorname{ZQSS}(A)$ and ZTQSS $(A)$ are compact, convex subsets of QSS $(A)$ and both are Choquet simplices, whose extreme points are the free product states and free product tracial states, respectively:

$$
\begin{aligned}
\partial_{e}(\operatorname{ZQSS}(A)) & =\left\{*_{1}^{\infty} \phi \mid \phi \in \mathrm{S}(A)\right\} \\
\partial_{e}(\operatorname{ZTQSS}(A)) & =\left\{*_{1}^{\infty} \tau \mid \tau \in \operatorname{TS}(A)\right\}
\end{aligned}
$$

Remark
$\operatorname{ZQSS}(A)$ is that part of QSS $(A)$ which is in analogy to Størmer's
result on symmetric states on the minimal tenser nondut $\operatorname{cin}^{\infty} \|$.

## A final question ...

It is known that $\operatorname{TS}(A)$, if non-empty, forms a Choquet simplex.
Question
Is $\operatorname{TQSS}(A)$ a Choquet simplex whenever $\operatorname{TS}(A)$ is non-empty?

Thank you for your attention!

