Damping-induced self-recovery phenomenon in mechanical systems with an unactuated cycle variable

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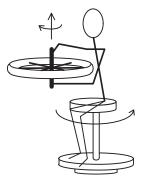
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Angular Momentum Conservation



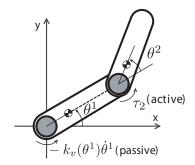


Thought Experiment



$$I_{\rm s}\omega_{\rm s} + I_{\rm w}\omega_{\rm w} = 0.$$

Horizontally Planar 2-Link Arm



$$I_{\rm i}\omega_{\rm i} + I_{\rm o}\omega_{\rm o} = 0.$$

Horizontally Planar 2-Link Arm: With or Without Damping

without damping

with damping

Horizontally Planar 2-Link Arm: Global Self-Recovery

Self-recovery is global, remembering the winding number.

Mechanical System with an Unactuated Cyclic Variable

- Configuration space Q = open subset of \mathbb{R}^n .
- Lagrangian $L(\mathbf{q},\dot{\mathbf{q}}) = \frac{1}{2}m_{ij}\dot{q}^i\dot{q}^j V(\mathbf{q})$ with cyclic variable q^1

$$\frac{\partial L}{\partial q^1} = 0$$

• Equations of Motion (EL equations with forces):

$$\begin{array}{ll} \displaystyle \frac{d}{dt} \frac{\partial L}{\partial \dot{q}^1} & = -k_v(q^1) \dot{q}^1 \\ \displaystyle \frac{d}{dt} \frac{\partial L}{\partial \dot{q}^a} - \frac{\partial L}{\partial q^a} = u_a, \quad a = 2, \, ..., \, n \end{array}$$

where

- $-k_v(q^1)\dot{q}^1$ is a viscous damping force
- u_2, \ldots, u_n are control forces

Mechanical System with an Unactuated Cyclic Variable

• Lagrangian $L(\mathbf{q}, \dot{\mathbf{q}}) = \frac{1}{2}m_{ij}\dot{q}^i\dot{q}^j - V(\mathbf{q})$ with cyclic variable q^1

$$\frac{\partial L}{\partial q^1} = 0.$$

• Equations of Motion (EL equations with forces):

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}^1} = -k_v(q^1)\dot{q}^1 \frac{d}{dt} \frac{\partial L}{\partial \dot{q}^a} - \frac{\partial L}{\partial q^a} = u_a, \quad a = 2, ..., n$$

• Without damping $(k_v = 0)$

$$\frac{d}{dt}\frac{\partial L}{\partial \dot{q}^1} = 0 \Rightarrow \frac{\partial L}{\partial \dot{q}^1} = \text{conserved}.$$

Mechanical System with an Unactuated Cyclic Variable

• Lagrangian $L(\mathbf{q}, \dot{\mathbf{q}}) = \frac{1}{2}m_{ij}\dot{q}^i\dot{q}^j - V(\mathbf{q})$ with cyclic variable q^1 such that

$$\frac{\partial L}{\partial q^1} = 0.$$

• Equations of Motion (EL equations with forces):

$$\begin{array}{ll} \displaystyle \frac{d}{dt} \frac{\partial L}{\partial \dot{q}^1} & = -k_v(q^1) \dot{q}^1 \\ \displaystyle \frac{d}{dt} \frac{\partial L}{\partial \dot{q}^a} - \frac{\partial L}{\partial q^a} = u_a, \quad a = 2, \, ..., \, n \end{array}$$

New conserved quantity with damping

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}^{1}} + \int_{0}^{q^{1}} k_{v}(x) dx \right) = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}^{1}} + k_{v}(q^{1}) \dot{q}^{1} = 0$$

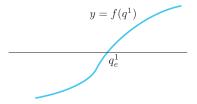
$$\Rightarrow \underbrace{\frac{\partial L}{\partial \dot{q}^{1}}}_{\text{damping-added momentum}} + \int_{0}^{q^{1}} k_{v}(x) dx = \text{conserved}$$

•
$$\int_0^{q^1(t)} k_v(x) dx = \int_0^t k_v(x) \dot{x} dt = (-)$$
 impulse due to friction.

Damping-Induced Self-Recovery Phenomenon Theorem (Chang and Jeon [2013, ASME J. DSMC]) Let

$$\mu = \frac{\partial L}{\partial \dot{q}^{1}} + \int_{0}^{q^{1}} k_{v}(x) dx = m_{1i}(\mathbf{q}(t)) \dot{q}^{i}(t) + \int_{0}^{q^{1}(t)} k_{v}(x) dx.$$

Let $f(q^1) = \int_0^{q^1} k_v(x) dx - \mu$ such that



Suppose controls $u_a(t)$'s (a = 2, ..., n) are chosen such that $q^a(t)$'s (a = 2, ..., n) are bounded and $\lim_{t \to \infty} \dot{q}^a(t) = 0$ for all a = 2, ..., n. Then, 1. $lim_{t \to \infty} q^1(t) = q_e^1$. 2. If the initial condition is such that $\dot{q}^i(0) = 0$ for all i = 1, ..., n, then $lim_{t \to \infty} q^1(t) = q^1(0)$. Sketch of Proof for Constant k_v with $\mu = 0$

Equation for q^1

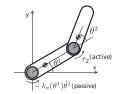
$$0 = m_{1i}\dot{q}^{i}(t) + \int_{0}^{q^{1}} k_{v}dx = m_{1i}\dot{q}^{i}(t) + k_{v}q^{1}$$
$$\Rightarrow \dot{q}^{1} = -\frac{k_{v}}{m_{11}}q^{1} + \left(-\frac{1}{m_{11}}\sum_{a=2}^{n}m_{1a}\dot{q}^{a}\right),$$

where $\dot{q}^i(0) = 0$ for all $i = 1, \dots, n$ and $q^1(0) = 0$. Hence,

$$\lim_{t \to \infty} \dot{q}^a = 0 \quad \forall a = 2, \dots, n \Rightarrow \lim_{t \to \infty} q^1(t) = 0 = q^1(0).$$

Remark: Damping coefficient $k_v(q^1)$ does not have to be a non-negative function. For example, $k_v(q^1) = 1 + 4\cos(q^1)$ shows self-recovery for $\mu = 0$.

Damping-Induced Bound



Suppose

$$\lim_{q^1\to\infty}\int_0^{q^1}k_v(x)dx=\infty,\qquad\qquad \lim_{q^1\to-\infty}\int_0^{q^1}k_v(x)dx=-\infty.$$

If controls $u_a(t)$'s are chosen such that $m_{11}(\mathbf{q}(t))$ is bounded above and below by two positive numbers and $m_{1a}(\mathbf{q}(t))$'s and $\dot{q}^a(t)$'s are bounded where a = 2, ..., n, then $q^1(t)$ is also bounded. Damping-Induced Bound for Horizontally Planar 2-Link Arm

The motion of Link 1 (θ^1) is bounded when $\dot{\theta}^2$ is bounded.

Real Experiment

Several Unactuated Cyclic Variables

Link 2 (θ^2) is actuated and Links 1 and 3 (θ^1, θ^3) are unactuated but under friction.

Self-Recovery Seems to Occur Only for Linear Friction

Cubic friction $F = -kv^3$.

Summary

- Viscous damping force breaks symmetry, so the corresponding momentum is no longer conserved.
- Exists a new conserved quantity called *damping-added momentum*.
- Damping-induced self-recovery is global.
- Damping puts a bound on range of the unactuated variable.
- References:
 - D.E. Chang and S. Jeon, "Damping-induced self recovery phenomenon in mechanical systems with an unactuated cyclic variable," ASME Journal of Dynamic Systems, Measurement, and Control, 135(2), 2013. http://dx.doi.org/10.1115/1.4007556
 - D.E. Chang and S. Jeon, "On the damping-induced self-recovery phenomenon in mechanical systems with several unactuated cyclic variables," J. Nonlinear Science, Submitted. http://arxiv.org/abs/1302.2109