

New Lax-Oleinik Type Operators in the Weak KAM Theory

—A joint work with Kaizhi Wang

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In this talk, we introduce some modifications of the Lax-Oleinik operators in the context of weak KAM theory, both in the autonomous and non-autonomous setting. These modified operators enjoy better properties of convergence: In the non autonomous case, they do converge, unlike the genuine L-O operators, and in the autonomous case, they converge more quickly. For example, when the Mather set is a quasiperiodic torus the new L-O operators converge faster than the genuine ones in the sense of order.

1. Background

Let M be a closed and connected smooth manifold.

Standard assumptions in Mather theory:

Consider a C^2 Lagrangian $L : TM \times \mathbb{R}^1 \rightarrow \mathbb{R}^1$, $(x, v, t) \mapsto L(x, v, t)$. We suppose that L satisfies the following conditions introduced by Mather:

- (H1) **Periodicity.** L is 1-periodic in the \mathbb{R}^1 factor, i.e., $L(x, v, t) = L(x, v, t + 1)$ for all $(x, v, t) \in TM \times \mathbb{R}^1$.
- (H2) **Positive Definiteness.** For each $x \in M$ and each $t \in \mathbb{R}^1$, the restriction of L to $T_x M \times t$ is strictly convex in the sense that its Hessian second derivative is everywhere positive definite.

(H3) **Superlinear Growth.** $\lim_{\|v\|_x \rightarrow +\infty} \frac{L(x,v,t)}{\|v\|_x} = +\infty$ uniformly on $x \in M$, $t \in \mathbb{R}^1$, where $\|\cdot\|_x$ denotes the norm on $T_x M$ induced by a Riemannian metric. By the compactness of M , this condition is independent of the choice of the Riemannian metric.

(H4) **Completeness of the Euler-Lagrange Flow.** Every solution of the Euler-Lagrange equation, which in local coordinates is:

$$\frac{d}{dt} \frac{\partial L}{\partial v}(x, \dot{x}, t) = \frac{\partial L}{\partial x}(x, \dot{x}, t),$$

are defined on all of \mathbb{R}^1 .

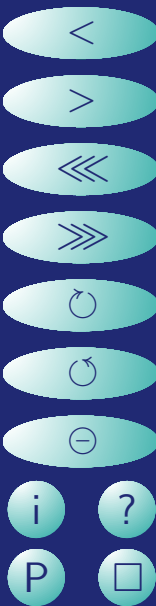
The corresponding Hamiltonian equations read

$$\dot{x} = \frac{\partial H}{\partial p}, \quad \dot{p} = -\frac{\partial H}{\partial x},$$

where $H(x, p) = p\dot{x} - L(x, \dot{x}, t)$ and $p = \frac{\partial L}{\partial \dot{x}}$. The corresponding Hamilton-Jacobi equation is

$$u_t + H(x, u_x, t) = c, \tag{1}$$

where $c = c(L)$ is the Mañé critical value of the Lagrangian L .



For all $t_2 \geq t_1$ and $x, y \in M$, let

$$F_{t_1, t_2}(x, y) = \inf_{\gamma} \int_{t_1}^{t_2} L(\gamma(s), \dot{\gamma}(s), s) ds,$$

where the infimum is taken over the continuous and piecewise C^1 paths $\gamma : [t_1, t_2] \rightarrow M$ such that $\gamma(t_1) = x$ and $\gamma(t_2) = y$. Define the action potential and extended Peierls barrier as follows. For each $(t_1, t_2) \in \mathbb{S}^1 \times \mathbb{S}^1$, let

$$\Phi_{\tau_1, \tau_2}(x, y) = \inf F_{t_1, t_2}(x, y)$$

for all $(x, y) \in M \times M$, where the infimum is taken on the set of $(t_1, t_2) \in \mathbb{R}^2$ such that $\tau_1 = [t_1]$, $\tau_2 = [t_2]$ and $t_2 \geq t_1 + 1$. For each $(\tau_1, \tau_2) \in \mathbb{S}^1 \times \mathbb{S}^1$, let

$$h_{\tau_1, \tau_2}(x, y) = \liminf_{t_2 - t_1 \rightarrow +\infty} F_{t_1, t_2}(x, y)$$

for all $(x, y) \in M \times M$, where the liminf is restricted to the set of $(t_1, t_2) \in \mathbb{R}^2$ such that $\tau_1 = [t_1]$, $\tau_2 = [t_2]$.

Standard assumptions in the weak KAM theory:

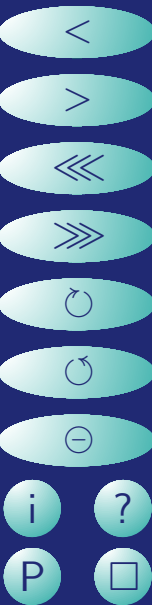
Let $L_a : TM \rightarrow \mathbb{R}^1$, $(x, v) \mapsto L_a(x, v)$ be a C^2 Lagrangian satisfying the following two conditions:

(H2') **Positive Definiteness.** For each $(x, v) \in TM$, the Hessian second derivative $\frac{\partial^2 L_a}{\partial v^2}(x, v)$ is positive definite.

(H3') **Superlinear Growth.** $\lim_{\|v\|_x \rightarrow +\infty} \frac{L_a(x, v)}{\|v\|_x} = +\infty$ uniformly on $x \in M$.

The corresponding Hamilton-Jacobi equation is

$$H_a(x, u_x) = c(L_a). \quad (2)$$



Definition (Lax-Oleinik semigroup) For each $u \in C(M, \mathbb{R}^1)$ and each $t \geq 0$, let

$$T_t^a u(x) = \inf_{\gamma} \left\{ u(\gamma(0)) + \int_0^t L_a(\gamma(s), \dot{\gamma}(s)) ds \right\} \quad (3)$$

for all $x \in M$, and

$$T_t u(x) = \inf_{\gamma} \left\{ u(\gamma(0)) + \int_0^t L(\gamma(s), \dot{\gamma}(s), s) ds \right\} \quad (4)$$

for all $x \in M$, where the infimums are taken among the continuous and piecewise C^1 paths $\gamma : [0, t] \rightarrow M$ with $\gamma(t) = x$. In view of (3) and (4), for each $t \geq 0$, T_t^a and T_t are operators from $C(M, \mathbb{R}^1)$ to itself. It is not difficult to check that $\{T_t^a\}_{t \geq 0}$ and $\{T_n\}_{n \in \mathbb{N}}$ are one-parameter semigroups of operators. $\{T_t^a\}_{t \geq 0}$ and $\{T_n\}_{n \in \mathbb{N}}$ are called the L-O semigroup associated with L_a and L , respectively.

Definition (weak KAM solution—time-independent case) A weak KAM solution of the Hamilton-Jacobi equation (2) is a function $u : M \rightarrow \mathbb{R}^1$ such that

(1) u is dominated by L , i.e.,

$$u(x) - u(y) \leq \Phi_{0,0}(y, x), \quad \forall x, y \in M.$$

(2) For every $x \in M$ there exists a curve $\gamma : (-\infty, 0] \rightarrow M$ with $\gamma(0) = x$ such that

$$u(x) - u(\gamma(t)) = \int_t^0 L_a(\gamma(s), \dot{\gamma}(s)) ds, \quad \forall t \in (-\infty, 0].$$

Theorem (Fathi)

- (1) For every $u \in C(M, \mathbb{R}^1)$, the uniform limit $\lim_{t \rightarrow +\infty} T_t^a u = \bar{u}$ exists and \bar{u} is a weak KAM solution of (2).
- (2) The weak KAM solutions and viscosity solutions are the same.

Fathi (1998) raised the question as to **whether the analogous result holds in the time-periodic case**. This would be the convergence of $T_n u$, $\forall u \in C(M, \mathbb{R}^1)$, as $n \rightarrow +\infty$, $n \in \mathbb{N}$.

Fathi and Mather (2000) gave a negative answer to the question.

2. New Lax-Oleinik type operators

2.1 Time-periodic case

Definition (new L-O operator—time-periodic case) For each $\tau \in [0, 1]$, each $n \in \mathbb{N}$ and each $u \in C(M, \mathbb{R}^1)$, let

$$\tilde{T}_n^\tau u(x) = \inf_{\substack{k \in \mathbb{N} \\ n \leq k \leq 2n}} \inf_{\gamma} \left\{ u(\gamma(0)) + \int_0^{\tau+k} L(\gamma(s), \dot{\gamma}(s), s) ds \right\}$$

for all $x \in M$, where the second infimum is taken among the continuous and piecewise C^1 paths $\gamma : [0, \tau + k] \rightarrow M$ with $\gamma(\tau + k) = x$.

For each $\tau \in [0, 1]$ and each $n \in \mathbb{N}$, \tilde{T}_n^τ is an operator from $C(M, \mathbb{R}^1)$ to itself. We call \tilde{T}_n^τ *the new L-O operator associated with L* . For each $n \in \mathbb{N}$ and each $u \in C(M, \mathbb{R}^1)$, let $U_n^u(x, \tau) = \tilde{T}_n^\tau u(x)$ for all $(x, \tau) \in M \times [0, 1]$. Then U_n^u is a continuous function on $M \times [0, 1]$.

Definition (weak KAM solution—time-periodic case) A weak KAM solution of the Hamilton-Jacobi equation (1) is a function $u : M \times \mathbb{S}^1 \rightarrow \mathbb{R}^1$ such that

(1) u is dominated by L , i.e.,

$$u(x, \tau) - u(y, s) \leq \Phi_{s, \tau}(y, x), \quad \forall (x, \tau), (y, s) \in M \times \mathbb{S}^1.$$

(2) For every $(x, \tau) \in M \times \mathbb{S}^1$ there exists a curve $\gamma : (-\infty, \tilde{\tau}] \rightarrow M$ with $\gamma(\tilde{\tau}) = x$ and $[\tilde{\tau}] = \tau$ such that

$$u(x, \tau) - u(\gamma(t), [t]) = \int_t^{\tilde{\tau}} L(\gamma(s), \dot{\gamma}(s), s) ds, \quad \forall t \in (-\infty, \tilde{\tau}].$$

Now we come to the main result:

Theorem (Wang and Yan, 2010) For each $u \in C(M, \mathbb{R}^1)$, the uniform limit $\lim_{n \rightarrow +\infty} U_n^u$ exists and

$$\lim_{n \rightarrow +\infty} U_n^u(x, \tau) = \inf_{y \in M} (u(y) + h_{0, [\tau]}(y, x))$$

for all $(x, \tau) \in M \times [0, 1]$, where $[\tau] = \tau \bmod 1$, and h denotes the extended Peierls barrier. Furthermore, let $\bar{u}(x, [\tau]) = \inf_{y \in M} (u(y) + h_{0, [\tau]}(y, x))$. Then $\bar{u} : M \times \mathbb{S}^1 \rightarrow \mathbb{R}^1$ is a weak KAM solution of the Hamilton-Jacobi equation (1)

$$u_s + H(x, u_x, s) = 0.$$

Another important result states as follows.

Theorem (Wang and Yan, 2010) Let $\bar{u} \in C(M \times \mathbb{S}^1, \mathbb{R}^1)$. Then the following three statements are equivalent.

- There exists $u \in C(M, \mathbb{R}^1)$ such that the uniform limit $\lim_{n \rightarrow +\infty} U_n^u = \bar{u}$.
- \bar{u} is a weak KAM solution of (1).
- \bar{u} is a viscosity solution of (1).

2.2 Time-independent case

Definition (new L-O operator—time-independent case) For each $u \in C(M, \mathbb{R}^1)$ and each $t \geq 0$, let

$$\tilde{T}_t^a u(x) = \inf_{t \leq \sigma \leq 2t} \inf_{\gamma} \left\{ u(\gamma(0)) + \int_0^\sigma L_a(\gamma(s), \dot{\gamma}(s)) ds \right\}$$

for all $x \in M$, where the second infimum is taken among the continuous and piecewise C^1 paths $\gamma : [0, \sigma] \rightarrow M$ with $\gamma(\sigma) = x$.

It is easy to check that $\{\tilde{T}_t^a\}_{t \geq 0} : C(M, \mathbb{R}^1) \rightarrow C(M, \mathbb{R}^1)$ is a one-parameter semigroup of operators. We call it *the new L-O semigroup associated with L_a* .

Theorem (Wang and Yan, 2010)

(1) For each $u \in C(M, \mathbb{R}^1)$, the uniform limit $\lim_{t \rightarrow +\infty} \tilde{T}_t^a u$ exists and

$$\lim_{t \rightarrow +\infty} \tilde{T}_t^a u = \lim_{t \rightarrow +\infty} T_t^a u = \bar{u}.$$

(2) For each $t \geq 0$ and each $u \in C(M, \mathbb{R}^1)$, $\|\tilde{T}_t^a u - \bar{u}\|_\infty \leq \|T_t^a u - \bar{u}\|_\infty$.

3. Rates of convergence of the L-O semigroup (time-independent case) and the family of the new L-O operators.

We believe that there is a deep relation between dynamical properties of the Aubry set (Mather set) and the rates of convergence of the L-O semigroup (time-independent case) and the family of the new L-O operators.

3.1. Results on the rate of convergence of the L-O semigroup $\{T_t^a\}_{t \geq 0}$

1. Iturriaga and Sánchez-Morgado (2009) proved that if the Aubry set consists in a finite number of hyperbolic periodic orbits or hyperbolic fixed points, the L-O semigroup converges exponentially.

2. Wang and Yan (2010) discussed the rate of convergence problem when the Mather set is a quasi-periodic invariant torus of the Euler-Lagrange flow. Consider a class of C^2 superlinear and strictly convex Lagrangians on \mathbb{T}^n

$$L_a^1(x, v) = \frac{1}{2} \langle A(x)(v - \omega), (v - \omega) \rangle + f(x, v - \omega), \quad x \in \mathbb{T}^n, \quad v \in \mathbb{R}^n, \quad (5)$$

where $A(x)$ is an $n \times n$ matrix, $\omega \in \mathbb{S}^{n-1}$ is a given vector, and $f(x, v - \omega) = O(\|v - \omega\|^3)$ as $v - \omega \rightarrow 0$. It is clear that $c(L_a^1) = 0$ and $\tilde{\mathcal{M}}_0 = \tilde{\mathcal{A}}_0 = \tilde{\mathcal{N}}_0 = \cup_{x \in \mathbb{T}^n} (x, \omega)$, which is a quasi-periodic invariant torus with frequency vector ω of the Euler-Lagrange flow associated to L_a^1 . For (5), the authors showed that for each $u \in C(\mathbb{T}^n, \mathbb{R}^1)$, there is a constant $K_2 > 0$ such that

$$\|T_t^a u - \bar{u}\|_\infty \leq \frac{K_2}{t}, \quad \forall t > 0. \quad (6)$$

An example was provided to show that the above result is sharp in the sense of order.

An example.

Consider the following integrable Lagrangian

$$L = \frac{1}{2} \langle v - \omega, v - \omega \rangle.$$

Take

$$u(x) = \begin{cases} \delta - \text{dist}(x, x_0), & \text{dist}(x, x_0) \leq \delta; \\ 0, & \text{otherwise.} \end{cases}$$

$$\lim_{t \rightarrow +\infty} T_t u(x) = \min_{x \in \mathbb{T}^n} u(x) \triangleq u_0.$$

There exist $t_m \rightarrow +\infty$ such that

$$|T_{t_m}^- u(x_0) - u_0(x_0)| \geq \frac{\delta^2}{32t_m}.$$

3.2. Results on the rate of convergence of the new L-O

semigroup $\{\tilde{T}_t^a\}_{t \geq 0}$

Recall the notations for Diophantine vectors: for $\varrho > n - 1$ and $\alpha > 0$, let

$$\mathcal{D}(\varrho, \alpha) = \left\{ \beta \in \mathbb{S}^{n-1} \mid |\langle \beta, k \rangle| \geq \frac{\alpha}{|k|^\varrho}, \forall k \in \mathbb{Z}^n \setminus \{0\} \right\},$$

where $|k| = \sum_{i=1}^n |k_i|$. For (5), Wang and Yan (2010) proved that given any frequency vector $\omega \in \mathcal{D}(\varrho, \alpha)$, for each $u \in C(\mathbb{T}^n, \mathbb{R}^1)$, there is a constant $K_3 > 0$ such that

$$\|\tilde{T}_t^a u - \bar{u}\|_\infty \leq K_3 t^{-(1 + \frac{4}{2\varrho + n})}, \quad \forall t > 0.$$

In view of (6) and (7), we conclude that the new L-O semigroup converges faster than the L-O semigroup in the sense of order when the Aubry set $\tilde{\mathcal{A}}_0$ of the Lagrangian system (5) is a quasi-periodic invariant torus with Diophantine frequency vector $\omega \in \mathcal{D}(\varrho, \alpha)$.

THANK YOU VERY MUCH!

Reference: A New Kind Of Lax-Oleinik Type Operator
With Parameters For Time-Periodic Positive Definite
Lagrangian Systems, accepted by CMP

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