# Large Deviations 2 

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## Interacting Particle Systems

- Kipnis, Rezakhanlou, Quastel, Jensen,


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- Local Equilibria


## Interacting Particle Systems

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- Conserved Quantities.
- Keep track of motion of many particles.
- Local Equilibria
- Averaging has to be done.
$\square$ Model. Symmetric Simple exclusion on $Z^{d}$.


## $\square$ Model. Symmetric Simple exclusion on $Z^{d}$.

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& (\mathcal{L} f)(\eta)=\sum_{u, v} p(v-u) \eta(u)(1-\eta(v))\left[f\left(\eta^{u, v}\right)-f(\eta)\right] \\
& \eta \rightarrow \eta^{u, v} \Leftrightarrow(1,0) \leftrightarrow(0,1) .
\end{aligned}
$$

$$
\begin{array}{r}
(\mathcal{A} f)\left(\left\{u_{j}\right\}\right)= \\
\sum_{i} \sum_{w} p(w)\left(1-\eta\left(u_{i}+w\right)\right) \\
{\left[f\left(\left\{u_{j}\right\}, i, u_{i}+w\right)-f\left(\left\{u_{j}\right\}\right)\right]}
\end{array}
$$

$\square \mathbf{Z}_{N}^{d}$, periodic. Number of particles $\rho N^{d}$
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$-\mathbf{Z}_{N}^{d}$, periodic. Number of particles $\rho N^{d}$
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$$
\begin{gathered}
\mathcal{L}_{N}=N^{2} \mathcal{L} ; \mathcal{A}_{N}=N^{2} \mathcal{A} \\
R=\frac{1}{N^{d}} \sum_{i} \delta_{\frac{u_{i}(\cdot)}{N}}
\end{gathered}
$$

## Initial Condition

$$
\frac{1}{N^{d}} \sum_{i} \delta_{\frac{u_{i}(0)}{N}}=\frac{1}{N^{d}} \sum_{u} \eta_{0}(u) \delta_{\frac{u}{N}}=r_{N}
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- A measure on the Torus $\mathcal{T}^{d}$.
$\square r_{N} \rightarrow \rho_{0}(x) d x$
$\square 0 \leq \rho_{0}(x) \leq 1$
$\int_{\mathcal{T}^{d}} \rho_{0}(x) d x=\rho$


## How does the density evolve?

$$
\square \frac{1}{N^{d}} \sum_{u} \delta_{\frac{u}{N}} \eta_{t}(u) \rightarrow \rho(t, x) d x
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& \square \frac{1}{N^{d}} \sum_{u} \delta_{\frac{u}{N}} \eta_{t}(u) \rightarrow \rho(t, x) d x \\
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- Converges in probability. Deviations are possible.
- Large Deviation probability.
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- How can we achieve a given profile $\rho(t, x)$ with $\rho(0, x)=\rho_{0}(x)$ ?
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- Large Deviation probability.
$\square$ How can we achieve a given profile $\rho(t, x)$ with $\rho(0, x)=\rho_{0}(x)$ ?
- The rates do not have to be equal. Introduce a bias.
$\square q:[0, T] \times \mathcal{T}^{d} \times Z^{d} \rightarrow R$,

$$
\begin{array}{r}
\left(\mathcal{L}_{N} f\right)(\eta)=N^{2} \sum_{u, v}\left(p(v-u)+\frac{q\left(t, \frac{u}{N}, v-u\right)}{N}\right) \\
\eta(u)(1-\eta(v))\left[f\left(\eta^{u, v}\right)-f(\eta)\right]
\end{array}
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- Large Deviation probability.
$\square$ How can we achieve a given profile $\rho(t, x)$ with $\rho(0, x)=\rho_{0}(x)$ ?
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\eta(u)(1-\eta(v))\left[f\left(\eta^{u, v}\right)-f(\eta)\right] \\
-q(t, x, z)+q(t, x,-z)=0 .
\end{gathered}
$$

$\square$ Effect: $b(t, x)=\sum_{w} w q(t, x, w)$

$$
\rho_{t}(t, x)=\frac{1}{2} \Delta \rho-\nabla \cdot(b(t, x) \rho(t, x)(1-\rho(t, x))
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Entropy Cost $H\left(\lambda_{2}, \lambda_{1}\right)=\lambda_{2} \log \frac{\lambda_{2}}{\lambda_{1}}-\left(\lambda_{2}-\lambda_{1}\right)$

$$
N^{2} E\left[\int_{0}^{T} \sum_{u, w}\left[\eta(t, u)(1-\eta(t, u+w)) H\left(p+\frac{q}{N}, p\right)\right] d t\right]
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Minimize over $q$, fixing $b$.
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Minimize over $q$, fixing $b$.
Replace $\eta(t, u)[1-\eta(t, u+w)]$ by
$\rho\left(t, \frac{x}{N}\right)\left(1-\rho\left(t, \frac{x}{N}\right)\right)$.

- The quantity reduces to

$$
\frac{N^{d}}{2} \int_{0}^{T} \int_{\mathcal{F} T^{d}}\|b(t, x)\|^{2} \rho(t, x)(1-\rho(t, x)) d t d x
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\rho_{t}=\frac{1}{2} \Delta-\nabla \cdot(b(t, x) \rho(t, x)(1-\rho(t, x)))
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$$
\begin{aligned}
& \rho_{t}=\frac{1}{2} \Delta-\nabla \cdot(b(t, x) \rho(t, x)(1-\rho(t, x))) \\
& \mathcal{I}(\rho)=\frac{1}{2} \int_{0}^{T}\left\|\rho_{t}-\frac{1}{2} \Delta \rho\right\|_{-1, \rho(t,)(1-\rho(\cdot))}^{2} d t
\end{aligned}
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\end{aligned}
$$

$\square$ This does not track individual particles.

- Track a particle in equilibrium. Bernoulli Density $\rho$. Diffuses. Exclude $1-d$ nearest neighbor.
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$\square$ Diffuses. $S(\rho)$. $S(\rho) \rightarrow I$ as $\rho \rightarrow 0$ and $S(\rho) \rightarrow 0$ as $\rho \rightarrow 1$.
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\mathcal{D}_{t}=\frac{1}{2} \nabla \cdot S(\rho(t, x)) \nabla+[S(\rho)-I] \frac{\nabla \rho}{2 \rho} \cdot \nabla
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\mathcal{D}_{t}=\frac{1}{2} \nabla \cdot S(\rho(t, x)) \nabla+[S(\rho)-I] \frac{\nabla \rho}{2 \rho} \cdot \nabla \\
\rho_{t}=D_{t}^{*} \rho
\end{gathered}
$$

$\square R_{n} \rightarrow P$.
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\mathcal{D}_{b, t}= & \frac{1}{2} \nabla \cdot S(\rho(t, x)) \nabla+[S(\rho)-I] \frac{\nabla \rho}{2 \rho} \cdot \nabla \\
& +b(t, x)(1-\rho(t, x)) \cdot \nabla
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- $R_{n} \rightarrow P$.
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$\mathcal{J}(b) . P_{b}$.

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- Probability $R_{n} \sim Q$
$\square Q$ determines $q(t, x), b(t, x)$ compatible with it.
$\square I(Q)=\mathcal{J}(b)+H\left(Q \mid P_{b}\right)$
$\square$ Can put in initial randomness. $\rho_{0}$ can be different.


## Totally asymmetric case

$\square d=1 . p(1)=1 . N t \rightarrow t$.

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$\square$ Solution is not unique.
$\square$ Entropy condition.

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- where $h^{\prime}(\rho)(1-2 \rho)=g^{\prime}(\rho)$
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- As a distribution,

$$
[h(\rho)]_{t}+[g(\rho)]_{x} \leq 0
$$

- One strictly convex function is enough.
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$\square h(p)=-[p \log p+(1-p) \log (1-p)]$
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- $h(p)=-[p \log p+(1-p) \log (1-p)]$
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$$
I(\rho(\cdot, \cdot))=\int_{\mathcal{S}} \int_{0}^{T}\left[[h(\rho)]_{t}+[g(\rho)]_{x}\right]^{+} d x d t
$$

## Diffusions in a random environment

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\frac{1}{2} a(x) D_{x}^{2}+b(x) D_{x}
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## Diffusions in a random environment

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\begin{gathered}
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\end{gathered}
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$\square a, b$ random, stationary in $x$.
$\square P\left[\frac{x(t)}{t} \simeq a\right] ?$

- Look at the Periodic case. $[0,1]$
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$\square$ There is a unique periodic invariant density $\phi(x)$. $\mathcal{L}_{a, b}^{*} \phi=0$.
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- Look at the Periodic case. $[0,1]$
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$\square P\left[\frac{x(t)}{t} \sim \int_{0}^{1} b(x) \phi(x) d x\right]=1$
$\square P\left[\frac{x(t)}{t} \sim m\right]=e^{-t I(m)+o(t)}$
- Look at the Periodic case. $[0,1]$
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$\square P\left[\frac{x(t)}{t} \sim \int_{0}^{1} b(x) \phi(x) d x\right]=1$
- $P\left[\frac{x(t)}{t} \sim m\right]=e^{-t I(m)+o(t)}$
- Change $b(x)$ to $c(x)$.
$\square$ Invariant density $\phi_{c}(x)$
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Now $Q\left[\frac{x(t)}{t} \sim \int_{0}^{1} c(x) \phi_{c}(x) d x\right] \sim 1$
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- Entropy cost is proportional to $t$ an the constant is

$$
\mathcal{J}(c)=\frac{1}{2} \int_{0}^{1} \frac{(c(x)-b(x))^{2}}{a(x)} \phi_{c}(x)
$$

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- Now $Q\left[\frac{x(t)}{t} \sim \int_{0}^{1} c(x) \phi_{c}(x) d x\right] \sim 1$
- Entropy cost is proportional to $t$ an the constant is

$$
\begin{gathered}
\mathcal{J}(c)=\frac{1}{2} \int_{0}^{1} \frac{(c(x)-b(x))^{2}}{a(x)} \phi_{c}(x) \\
I(m)=\inf _{c: \int_{0}^{1} c(x) \phi_{c}(x) d x=m} \mathcal{J}(c)
\end{gathered}
$$

- Random Stationary Case.
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- Invariant measure?
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- Another process $\phi(x, \omega) \geq 0$ such that $\phi, a, b$ are jointly stationary $E^{P}[\phi(x, \omega)]=1$, and
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$\square$ It may not exist. If it does it is unique.
- Random Stationary Case.
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- Another process $\phi(x, \omega) \geq 0$ such that $\phi, a, b$ are jointly stationary $E^{P}[\phi(x, \omega)]=1$, and
- $\mathcal{L}_{a, b}^{*} \phi=0$
$\square$ It may not exist. If it does it is unique.
$\square \frac{x(t)}{t} \rightarrow E[\phi(0, \omega) b(0, \omega)]$.
- Large Deviations


## - Large Deviations

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$\square P\left[\frac{x(t)}{t} \sim m\right]=\exp [-t I(m)+o(t)] ?$
$\square$ Find pairs $c, \phi_{c}$ such that $E\left[\phi_{c}\right]=1, E\left[c \phi_{c}\right]=m$

$$
\mathcal{L}_{a, c}^{*} \phi_{c}=0
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- Large Deviations
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\begin{gathered}
\mathcal{L}_{a, c}^{*} \phi_{c}=0 \\
\mathcal{J}(c)=\frac{1}{2} E\left[\frac{(c-b)^{2}}{a} \phi_{c}\right]
\end{gathered}
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\begin{gathered}
\mathcal{L}_{a, c}^{*} \phi_{c}=0 \\
\mathcal{J}(c)=\frac{1}{2} E\left[\frac{(c-b)^{2}}{a} \phi_{c}\right] \\
I(m)=\inf _{c: E\left[c \phi_{c}\right]=m} \mathcal{J}(c)
\end{gathered}
$$

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$\square$ Need the limit as $T \rightarrow \infty$ of $v_{T}(0,0)$.

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- Let

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h(x, b, \omega)=\sup _{v}[b v-f(x, v, \omega)]
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## $v_{\epsilon}(0,0, \omega)=$

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\sup _{b=b(s, x, \omega)} E\left[\theta x(1)-\int_{0}^{1} h\left(\frac{x(s)}{\epsilon}, b(s, x(s), \omega), \omega\right) d s\right]
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With $T=\epsilon^{-1}$, and rescaling space and time
$v_{T}(0,0, \omega)=$
$\sup _{b=b(s, x, \omega)} E\left[\theta \frac{x(T)}{T}-\frac{1}{T} \int_{0}^{T} h(x(s), b(s, x(s), \omega), \omega) d s\right]$

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$$
v(0,0)=\sup _{\left(b, P_{b}\right)} E^{P_{b}}[\theta b(\omega)-h(b(\omega), \omega)]
$$

## Last Slide

## THE END

