

The Geometry of Anticoherent Spin States

Jason Crann

Department of Mathematics & Statistics, University of Guelph

Canadian Quantum Information Student Conference, 2009



Outline

- 1 Anticoherence
 - Definition
 - The Majorana Representation
- 2 Spherical Designs
 - Symmetric Designs
 - Connection to Anticoherence
- 3 Application
 - Symmetric State Entanglement

Coherent States

- The most "classical" spin states.
- Can be identified with a classical spin direction: $|\psi_{\mathbf{n}}\rangle \sim \mathbf{n}$.
- For a spin- s system in the coherent state $|\psi_{\mathbf{n}}\rangle$,

$$\langle \mathbf{S} \rangle = (\langle S_x \rangle, \langle S_y \rangle, \langle S_z \rangle) = s\hbar \mathbf{n}.$$

Anticoherent States

- Cannot be associated with a classical spin direction.
- Point "nowhere" in the mean: $\langle \mathbf{S} \rangle = \langle \psi | \mathbf{S} | \psi \rangle = \mathbf{0} \Leftrightarrow \langle \mathbf{n} \cdot \mathbf{S} \rangle$ being independent of \mathbf{n} .

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Definition

A state $|\psi\rangle$ is anticoherent if $\langle \mathbf{n} \cdot \mathbf{S} \rangle$ is independent of \mathbf{n} .

Example: $|\psi\rangle = |s = 1, m_z = 0\rangle$ satisfies $\langle \mathbf{S} \rangle = \mathbf{0}$.

Higher Orders of Anticoherence

- Two labs L_1 and L_2 :
 - L_1 measures $\mathbf{n}_1 \cdot \mathbf{S}$ many times on a system prepared in the anticoherent state $|s = 1, m_z = 0\rangle$.
 - L_2 measures $\mathbf{n}_2 \cdot \mathbf{S}$ many times on the same type of system prepared in the same state.
- Directional dependence of $\langle (\mathbf{n} \cdot \mathbf{S})^2 \rangle$.

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Definition

A state $|\psi\rangle$ is anticoherent to order t if $\langle(\mathbf{n} \cdot \mathbf{S})^k\rangle$ is independent of \mathbf{n} for $k = 1, \dots, t$.

THE MAJORANA REPRESENTATION

Spin-1/2

$$|\psi\rangle = a_{\uparrow}|\uparrow\rangle + a_{\downarrow}|\downarrow\rangle$$

ratio $r = \frac{a_{\downarrow}}{a_{\uparrow}}$

root of $M(z) = a_{\uparrow}z - a_{\downarrow}$

Stereographic projection:

$$\mathbf{V}(r) = \left(\frac{2\Re(r)}{1+|r|^2}, \frac{2\Im(r)}{1+|r|^2}, \frac{1-|r|^2}{1+|r|^2} \right)$$

Spin-1/2

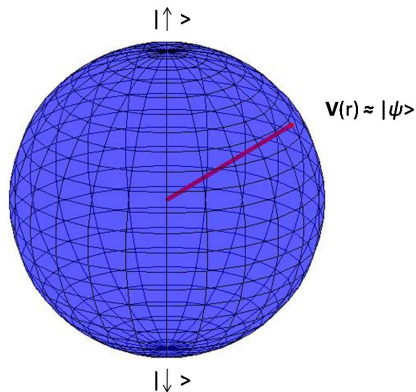
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Spin-1

$$|\psi\rangle = a_1|1\rangle + a_0|0\rangle + a_{-1}|-1\rangle$$



$$M(z) = a_1 z^2 - \sqrt{2} a_0 z + a_{-1}$$

two roots: r_1, r_2



two points on S^2 : $\mathbf{V}(r_1), \mathbf{V}(r_2)$

Spin-1

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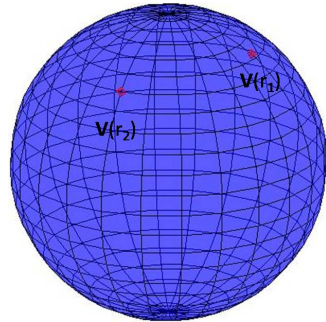


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States \Leftrightarrow points on \mathcal{S}^2

$$\text{Spin} - s : \quad |\psi\rangle = \sum_{m=-s}^s a_m |m\rangle \in \mathcal{H}_{2s+1}$$

$$\Updownarrow$$

$$M(z) = \sum_{m=-s}^s a_m (-1)^{m-s} \sqrt{\binom{2s}{s+m}} z^{s+m} \quad \text{degree} = 2s$$

$$\Updownarrow$$

$$\{r_i\}_{i=1}^{2s}$$

$$\Updownarrow$$

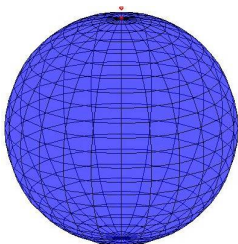
2s points on \mathcal{S}^2 via stereographic projection

Spin-1: special examples

Basis states: $|1\rangle, |0\rangle, |-1\rangle$

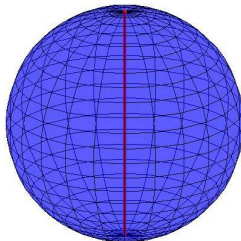
$|1\rangle$

spin up - coherent



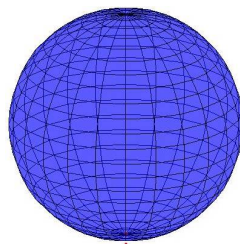
$|0\rangle$

anticoherent



$|-1\rangle$

spin down - coherent



Majorana Picture

- Coherent states: one point with multiplicity $2s$.
- Anticoherent states: points are "nicely" spread over the sphere.

SPHERICAL DESIGNS

Equivalent Definitions

Definition

Let $t \in \mathbb{N}$. A finite (non-empty) subset $X \subset \mathcal{S}^2$ is a spherical t -design if

$$\frac{1}{|X|} \sum_{\mathbf{x} \in X} f(\mathbf{x}) = \frac{1}{\mu(\mathcal{S}^2)} \int_{\mathcal{S}^2} f(\mathbf{x}) d\mu(\mathbf{x})$$

holds for any polynomial $f(\mathbf{x})$ of $\deg \leq t$.

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Observation

$X \subset \mathcal{S}^2$ is a spherical t -design if the k th moments of X are invariant under $\mathcal{O}(3)$ for $k = 1, \dots, t$.

A spherical 1-design is a set whose vector sum is zero.

Spherical Designs vs. Anticoherent States

Definition

A state $|\psi\rangle$ is anticoherent to order t if $\langle(\mathbf{n} \cdot \mathbf{S})^k\rangle$ is independent of \mathbf{n} for $k = 1, \dots, t$.

\Updownarrow via Majorana Representation ?

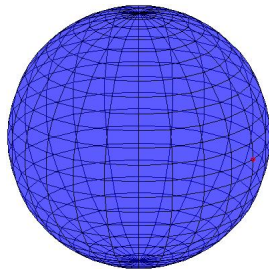
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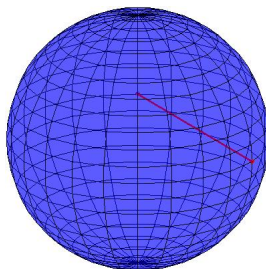
Orbits

Can generate spherical designs using orbits of finite subgroups of $\mathcal{O}(3)$. Take $G = \{I, R_{2\pi/3}, R_{4\pi/3}\}$ and $\mathbf{x} = (1, 0, 0)$.

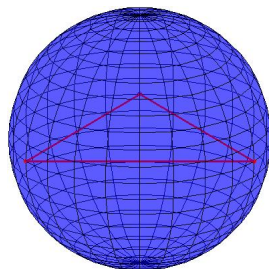
$$\{I\mathbf{x}\}$$



$$\{I\mathbf{x}, R_{2\pi/3}\mathbf{x}\}$$



$$\{I\mathbf{x}, R_{2\pi/3}\mathbf{x}, R_{4\pi/3}\mathbf{x}\}$$



Main Result

Theorem

Let $\mathbf{x} \in S^2$ and let G be a (non-empty) finite subgroup of $\mathcal{O}(3)$. If $G\mathbf{x}$ is a spherical t -design, then the state $|\psi\rangle$ whose Majorana representation is $G\mathbf{x}$ is anticoherent to order t .

Some Curious Geometry

Spin-2

- $G = \text{Sym}(T)$
- $G\mathbf{x}$ is a spherical 2-design for all $\mathbf{x} \in \mathcal{S}^2$.
- Gives anticoherent states of order 2.
- Example:

$$|\psi\rangle = \frac{1}{\sqrt{3}}(\sqrt{2}|1\rangle - |-2\rangle)$$

Some Curious Geometry

Spin-4

- $G = \text{Sym}(C)$
- $G\mathbf{x}$ is a spherical 3-design for all $\mathbf{x} \in \mathcal{S}^2$.
- Gives anticoherent states of order 3.
- Example:

$$\frac{1}{2\sqrt{6}}(\sqrt{5}|4\rangle + \sqrt{14}|0\rangle + \sqrt{5}|-4\rangle)$$

Some Curious Geometry

Spin-10

- $G = \text{Sym}(D)$
- $G\mathbf{x}$ is a spherical 5-design for all $\mathbf{x} \in \mathcal{S}^2$.
- Gives anticoherent states of order 5.
- Example:

$$a|10\rangle + b|5\rangle + |0\rangle - b|-5\rangle + a|-10\rangle$$

SYMMETRIC STATE ENTANGLEMENT

Symmetric States

Symmetric states on n qubits:

$$|\psi_{sym}\rangle = N \sum_{\sigma \in \mathcal{S}_n} |r_{\sigma(1)}\rangle |r_{\sigma(2)}\rangle \cdots |r_{\sigma(n)}\rangle$$

Example:

$$|GHZ\rangle = \frac{1}{\sqrt{2}}(|000\rangle + |111\rangle) = N \sum_{\sigma \in \mathcal{S}_3} |r_{\sigma(1)}\rangle |r_{\sigma(2)}\rangle |r_{\sigma(3)}\rangle$$

where

$$\begin{aligned} |r_1\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \\ |r_2\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + e^{i2\pi/3}|1\rangle) \\ |r_3\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + e^{i4\pi/3}|1\rangle) \end{aligned}$$

Geometric Measure of Entanglement

$$E_G(|\psi\rangle) = 1 - \max_{|\phi\rangle \in \text{Prod}} |\langle\psi|\phi\rangle|^2$$

- $E_G \geq 0$ with equality iff $|\psi\rangle$ is a product state.
- Maximum product state $|\phi\rangle$ can be found in symmetric subspace.
- Maximum product state $|\phi\rangle$ is *necessarily* symmetric.
[Hübener, Kleinmann, et. al, 2009]

Maximum Entanglement

- For symmetric states $|\psi\rangle$ and symmetric product states $|\phi\rangle = |a\rangle|a\rangle \cdots |a\rangle$,

$$|\langle\psi|\phi\rangle|^2 \propto \prod_{i=1}^n |\langle r_i|a\rangle|^2 = \prod_{i=1}^n \frac{\|\mathbf{V}(r_i) - \mathbf{V}(a^*)\|^2}{4}$$

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- For fixed n , symmetric states $|\psi\rangle$ which maximize $E_G(|\psi\rangle)$ satisfy

$$\min_{|\psi\rangle \in \text{Sym}(\mathcal{H})} \max_{\mathbf{x} \in S^2} \prod_{i=1}^n \|\mathbf{V}(r_i) - \mathbf{x}\|^2$$

3-qubits

$$|GHZ\rangle = \frac{1}{\sqrt{2}}(|000\rangle + |111\rangle)$$

- "Maximally entangled" via E_G . [Tamaryan, Wei, Park, 2009]

$$|r_1\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

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roots: $1, e^{i2\pi/3}, e^{i4\pi/3}$

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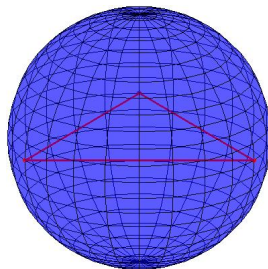
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roots: $1, e^{i2\pi/3}, e^{i4\pi/3}$

-of unity!

$|GHZ\rangle$



Spherical 1-design

Final Conjecture

- Symmetric state with highest geometric entanglement is that whose n points are mutually maximally distant ($n > 3$).
[Markham, 2007]

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Conjecture

Fix $n \in \mathbb{N}$. Let τ_n be the largest value of t for which an n -point spherical t -design exists. Symmetric states that are spherical τ_n -designs maximize E_G .

For Further Reading I



J. Zimba.

"Anticoherent" Spin States via the Majorana Representation.
Electronic Journal of Theoretical Physics, No.10, 143-156,
2006.



R. Hübener, M. Kleinmann, T.C. Wei, C. González-Guillén and
O. Gühne.

The geometric measure of entanglement for symmetric states
arXiv:0905.4822v2, 2009.



S. Tamaryan, T.C. Wei, D. Park.

Maximally entangled three-qubit states via geometric measure
of entanglement
arXiv:0905.3791v2, 2009.

For Further Reading II



D. Markham

Symmetric State Entanglement in the Majorana Representation

<http://www.pps.jussieu.fr/~djhm>

Thank you

- Rajesh Pereira
- David Kribs
- NSERC



- Fields Institute



- Thank you all!