

Zero discord

Aharon Brodutch

Macquarie University

August 2009



Local quantum vs classical statistics

- A classical probability distribution (mixed state)

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\rho_i$$

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow$

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|}$$

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_x \Rightarrow$$

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_x \Rightarrow \boxed{\rho_f = \frac{1}{2} [|\uparrow\rangle \langle\uparrow| + |\downarrow\rangle \langle\downarrow|]}$$

- But

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_x \Rightarrow \boxed{\rho_f = \frac{1}{2} [|\uparrow\rangle \langle\uparrow| + |\downarrow\rangle \langle\downarrow|]}$$

- But

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|}$$

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_x \Rightarrow \boxed{\rho_f = \frac{1}{2} [|\uparrow\rangle \langle\uparrow| + |\downarrow\rangle \langle\downarrow|]}$$

- But

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_z \Rightarrow$$

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_x \Rightarrow \boxed{\rho_f = \frac{1}{2} [|\uparrow\rangle \langle\uparrow| + |\downarrow\rangle \langle\downarrow|]}$$

- But

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_z \Rightarrow \boxed{\rho_i = \rho_f}$$

Local quantum vs classical statistics

- A classical probability distribution (mixed state)

$$\boxed{\rho_i} \Rightarrow \text{measurement} \Rightarrow \boxed{\rho_f = \rho_i}$$

- A quantum state

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_x \Rightarrow \boxed{\rho_f = \frac{1}{2} [|\uparrow\rangle \langle\uparrow| + |\downarrow\rangle \langle\downarrow|]}$$

- But

$$\boxed{\rho_i = \alpha |\uparrow\rangle_z \langle\uparrow| + \beta |\downarrow\rangle_z \langle\downarrow|} \Rightarrow \sigma_z \Rightarrow \boxed{\rho_i = \rho_f}$$



Quantum measurements on a bipartite system

A classical bipartite system



Entropy: $H(A) = \sum_i p_i^A \log_2 p_i^A$, $H(A, B) = \sum_i p_i^{A,B} \log_2 p_i^{A,B}$

Quantum measurements on a bipartite system

A classical bipartite system

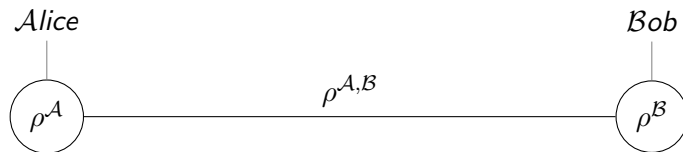


Entropy: $H(A) = \sum_i p_i^A \log_2 p_i^A$, $H(A, B) = \sum_i p_i^{A,B} \log_2 p_i^{A,B}$

Conditional Entropy: $H(A|B) = \sum_i p(B = b_i) H(A|B = b_i)$

Quantum measurements on a bipartite system

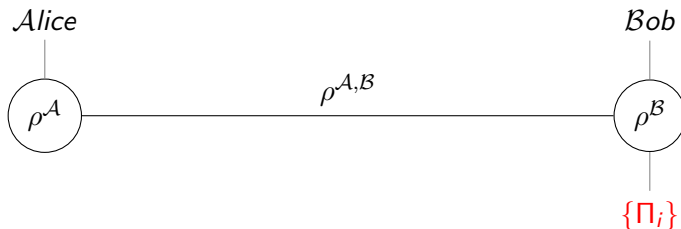
A quantum bipartite system



Entropy: $H(A) = -\rho^A \log_2 \rho^A$, $H(A, B) = -\rho^{A,B} \log_2 \rho^{A,B}$

Quantum measurements on a bipartite system

A quantum bipartite system



Entropy: $H(A) = \rho^A \log_2 \rho^A$, $H(A, B) = \rho^{A,B} \log_2 \rho^{A,B}$

Conditional Entropy: $H(A|B_{\{\Pi_i^B\}}) = \sum_i p(\Pi_i^B) H(A|\Pi_i^B)$

Mutual information

Classical Mutual information, two equivalent expressions

$$\mathcal{I}(A : B) = H(A) + H(B) - H(A, B) \quad (1)$$

$$\mathfrak{J}(A : B) = H(A) - H(A|B) \quad (2)$$

¹H. Ollivier and W. H. Zurek PRL 88 (2002)

Mutual information

Classical Mutual information, two equivalent expressions

$$\mathcal{I}(A : B) = H(A) + H(B) - H(A, B) \quad (1)$$

$$\mathfrak{J}(A : B) = H(A) - H(A|B) \quad (2)$$

For a quantum system the expression for 2 should be revised

$$\mathfrak{J}(\mathcal{A} : \mathcal{B})_{\{\Pi_j^{\mathcal{B}}\}} = H(\mathcal{A}) - H(\mathcal{A}|\{\Pi_j^{\mathcal{B}}\}) \quad (3)$$

¹H. Ollivier and W. H. Zurek PRL 88 (2002)

Mutual information

Classical Mutual information, two equivalent expressions

$$\mathcal{I}(A : B) = H(A) + H(B) - H(A, B) \quad (1)$$

$$\mathfrak{J}(A : B) = H(A) - H(A|B) \quad (2)$$

For a quantum system the expression for 2 should be revised

$$\mathfrak{J}(\mathcal{A} : \mathcal{B})_{\{\Pi_j^{\mathcal{B}}\}} = H(\mathcal{A}) - H(\mathcal{A}|\{\Pi_j^{\mathcal{B}}\}) \quad (3)$$

Discord¹

$$\delta(\mathcal{A} : \mathcal{B})_{\{\Pi_j^{\mathcal{B}}\}} = \mathcal{I} - \mathfrak{J}(\mathcal{A} : \mathcal{B})_{\{\Pi_j^{\mathcal{B}}\}}$$

¹H. Ollivier and W. H. Zurek PRL 88 (2002)

Discord

- Discord is the difference in the MI before and after a measurement
- The expression (3) can be used to measure classical correlation²
- Usually some sort of optimization over all possible local measurements is used to get the Discord

$$\delta(\mathcal{A} : \mathcal{B})_{\{\Pi_j^{\mathcal{B}}\}} = \min_{\{\Pi_j^{\mathcal{B}}\}} [H(\mathcal{B}) - H(\mathcal{A}, \mathcal{B}) + H(\mathcal{A} | \{\Pi_j^{\mathcal{B}}\})] \quad (4)$$

- Separable systems can have non zero discord.
- Discord is not symmetric

²L. Henderson and V. Vedral, J Phys 34 (2001).

Examples of discord

$$1 \quad \frac{1}{2} [|00\rangle \langle 00| + |1+\rangle \langle 1+|]$$

Examples of discord

$$\textcircled{1} \quad \frac{1}{2} [|00\rangle \langle 00| + |1+\rangle \langle 1+|] \quad \delta(\mathcal{A} : \mathcal{B}) \approx 0.2 ; \delta(\mathcal{B} : \mathcal{A})_{0,1} = 0$$

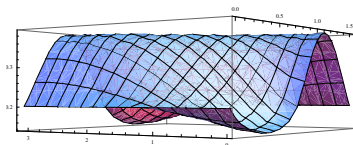
Examples of discord

$$① \quad \frac{1}{2} [|00\rangle \langle 00| + |1+\rangle \langle 1+|] \quad \delta(\mathcal{A} : \mathcal{B}) \approx 0.2 ; \delta(\mathcal{B} : \mathcal{A})_{0,1} = 0$$

$$② \quad \frac{1}{2} [|00\rangle \langle 00| + |++\rangle \langle ++|]$$

Examples of discord

- 1 $\frac{1}{2}[|00\rangle\langle 00| + |1+\rangle\langle 1+|]$ $\delta(\mathcal{A} : \mathcal{B}) \approx 0.2$; $\delta(\mathcal{B} : \mathcal{A})_{0,1} = 0$
- 2 $\frac{1}{2}[|00\rangle\langle 00| + |++\rangle\langle ++|]$ $\delta(\mathcal{A} : \mathcal{B}) = \delta(\mathcal{B} : \mathcal{A}) \approx 0.14$

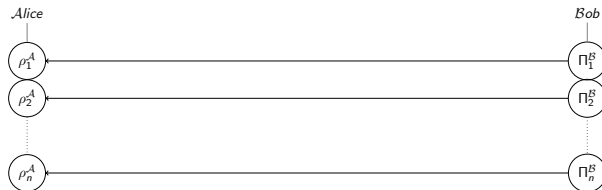


Using discord

- Measure of classical correlations.
- Quantum Maxwell's demons - Zurek (2003), Oppenheim, Horodecki, Horodecki, and Horodecki (2002).
- Correlations in DQC1 - Datta, Shaji and Caves (2008)
- The most interesting results involve **zero discord**.
- Zero discord defines a Pointer basis.
- CP evolution - Jordan, Shaji and Sudarshan (2007)

A zero discord systems can be as a sum of (orthogonal) systems of the type

$$\rho_{\mathcal{A},\mathcal{B}} = \sum_j a_j \rho_j^{\mathcal{A}} \otimes \Pi^{\mathcal{B}} \quad (5)$$



Zero discord between a system and the environment is essential for CP evolution.³

³Shabani and Lidar PRL 102 (2009)

Calculating zero discord

- Calculating discord requires optimization.

Calculating zero discord

- Calculating discord requires optimization. **It's hard work**



Calculating zero discord

- Calculating discord requires optimization. **It's hard work**
- The marginals are diagonal in the zero discord basis.

$$\rho^{\mathcal{B}} = \begin{pmatrix} a_1 & 0 & \dots & 0 \\ 0 & a_2 & \dots & 0 \\ \vdots & \dots & \dots & 0 \\ 0 & \dots & \dots & a_d \end{pmatrix} \quad (6)$$

Calculating zero discord

Given a density matrix ρ^{AB}

- 1 Find the marginal $\rho^B = \text{tr}_A \rho^{AB}$
- 2 Find the eigenstates $|j\rangle^B$
- 3 Use the eigenstates to build a projector basis $\Pi_j^B = |j\rangle^B \langle j|$
- 4 Calculate the discord in this basis

$$\delta(\mathcal{A} : \mathcal{B})_{\{\Pi_j^B\}} = H(\mathcal{B}) - H(\mathcal{A}, \mathcal{B}) + H(\mathcal{A} | \{\Pi_j^B\})$$
- 5 If and only if $\delta(\mathcal{A} : \mathcal{B})_{\Pi_j^B} = 0$, the state is a zero discord state.

Calculating zero discord

Given a density matrix ρ^{AB}

- 1 Find the marginal $\rho^B = \text{tr}_A \rho^{AB}$
- 2 Find the eigenstates $|j\rangle^B$
- 3 Use the eigenstates to build a projector basis $\Pi_j^B = |j\rangle^B \langle j|$
- 4 Calculate the discord in this basis

$$\delta(\mathcal{A} : \mathcal{B})_{\{\Pi_j^B\}} = H(\mathcal{B}) - H(\mathcal{A}, \mathcal{B}) + H(\mathcal{A} | \{\Pi_j^B\})$$
- 5 If and only if $\delta(\mathcal{A} : \mathcal{B})_{\Pi_j^B} = 0$, the state is a zero discord state.

Does not work for maximally mixed marginals :(

Locally distinguishable states

Non-locality without entanglement⁴

⁴Bennett et al PRA 59 (1999)

Locally distinguishable states

Non-locality without entanglement⁴

There is no simple relation between discord and distinguishability.

⁴Bennett et al PRA 59 (1999)

Locally distinguishable states

Non-locality without entanglement⁴

There is no simple relation between discord and distinguishability.

Examples:

State	Discord	Distinguishability
Equal mixture of locally orthogonal states	DZero	Local
Equal mixture of the 9 Bennett states	DZero	Non local
$ 00\rangle\langle 00 $ and $ ++\rangle\langle ++ $	Non zero	Impossible
A singlet and a $ 11\rangle$ state	Non zero	Local

⁴Bennett et al PRA 59 (1999)

Conclusions

- 1 Discord is a measure of quantum correlations.
- 2 Zero discord can be verified easily using the local diagonal basis.
- 3 There is no simple way to define distinguishability using discord.
- 4 Completely mixed marginals are strange.

The end

Questions?

