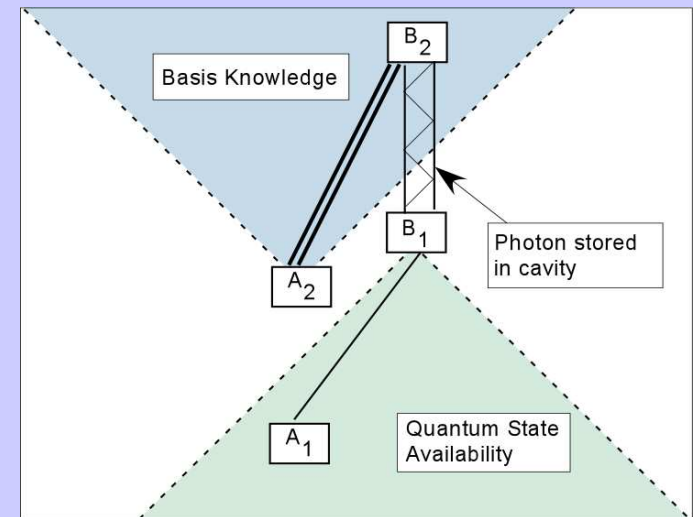
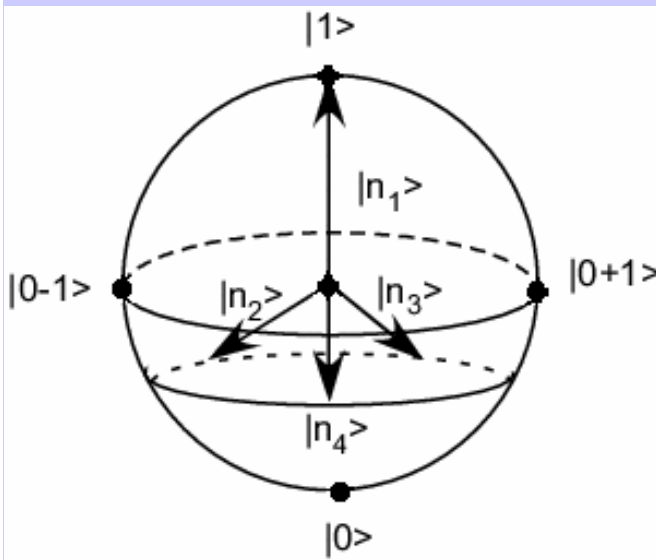


Quantum Communication via Entanglement: Quantum Orienteering & *Relativistic* Quantum Cryptography

Paul Kwiat



Outline

1. Quantum “Orienteering”
2. “Relativistic” Quantum Cryptography

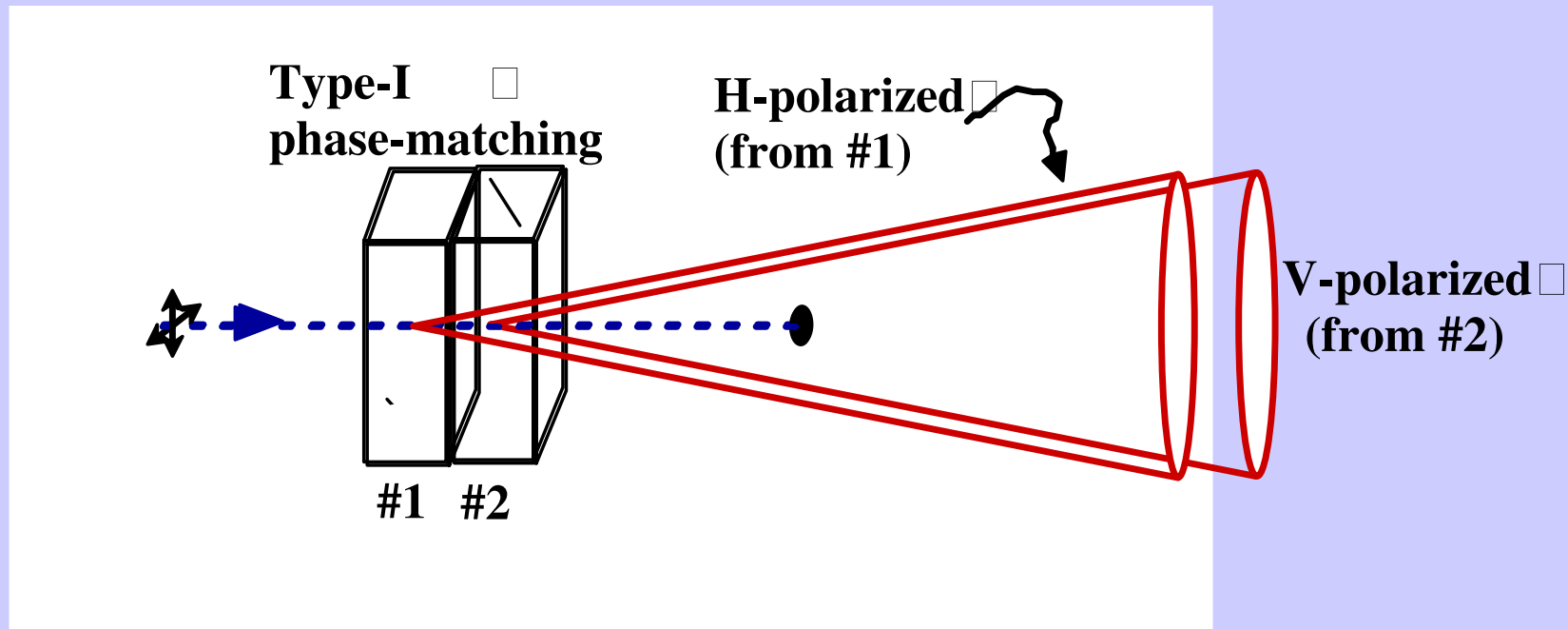
PGK Group

Graduate Students: **Joe Altepeter**, Julio Barreiro, Onur Hosten, **Evan Jeffrey**,
Nicholas Peters, Radhika Rangarajan, **Aaron VanDevender**, Joseph Yasi
Undergraduates: Kyle Arnold, Gleb Akselrod, Rachel Hillmer, Kevin Uskali
Associated Theory Post-Doc: Tzu-Cheih Wei



Postdocs welcome!

Two-crystal Polarization-Entangled Source:



$$|\psi\rangle = \frac{1}{\sqrt{2}} (|H\rangle_1 |H\rangle_2 + e^{i\varphi} |V\rangle_1 |V\rangle_2)$$

Maximally entangled state

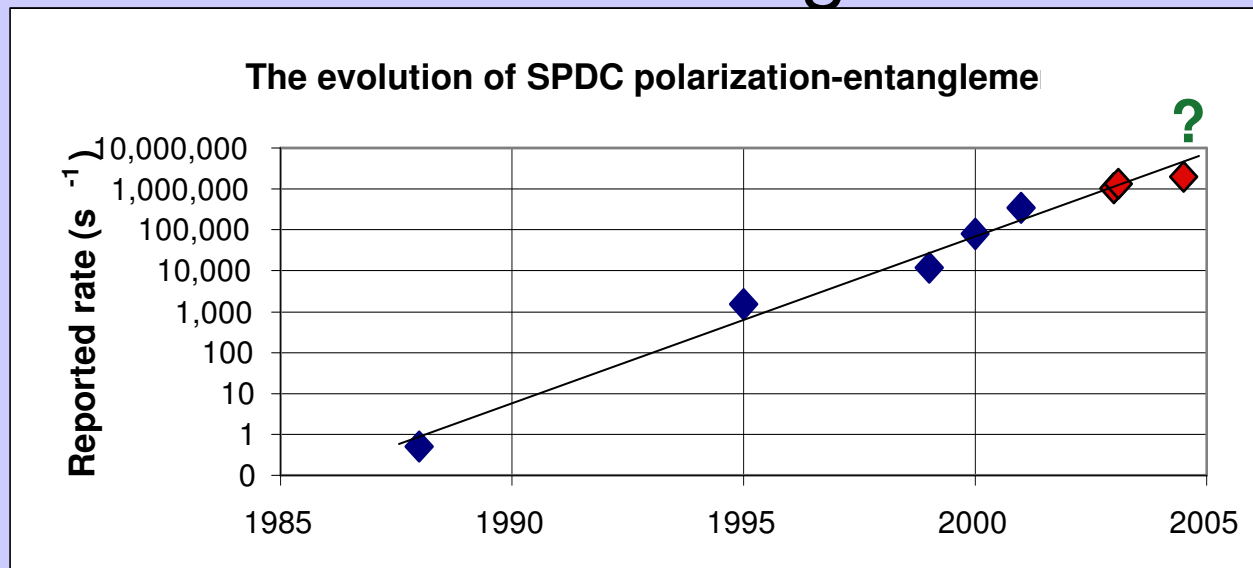
Tune pump polarization:

à **Nonmax. entangled states**

Add decoherence to arms

à **(Partially) mixed states**

Moore's law for entanglement



We have observed polarization-entangled pairs
@ **2,000,000 s⁻¹**, with **F ~98%**!

Next main limitation: detector saturation

Bell-Ineq. Tests

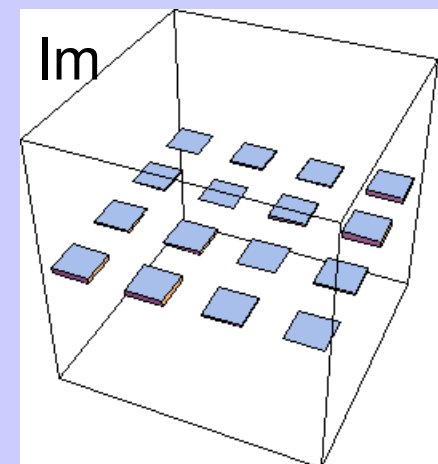
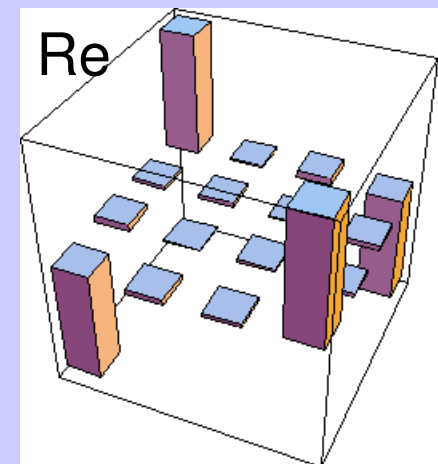
New source: $|S_{\text{expt}}| = 2.7260 \pm 0.0008$ (216 σ in 0.8 s)

($S_{\text{LHV}} \leq 2$) $|S_{\text{expt}}| = 2.7392 \pm 0.00008$ (2417 σ in 2 min)

Optimized $|S_{\text{QM, max}}| = 2\sqrt{2} = 2.828$

Bell test: $|S_{\text{expt}}| = 2.826 \pm 0.005$ 165 σ

$$\Phi^{(-)} \sim |HH\rangle - |VV\rangle$$



(Total counting
time = 10 s)

Opt. Exp. **13**, 8951 (2005)

Classical Orienteering

Directions can be transmitted classically by sending a spinning object.

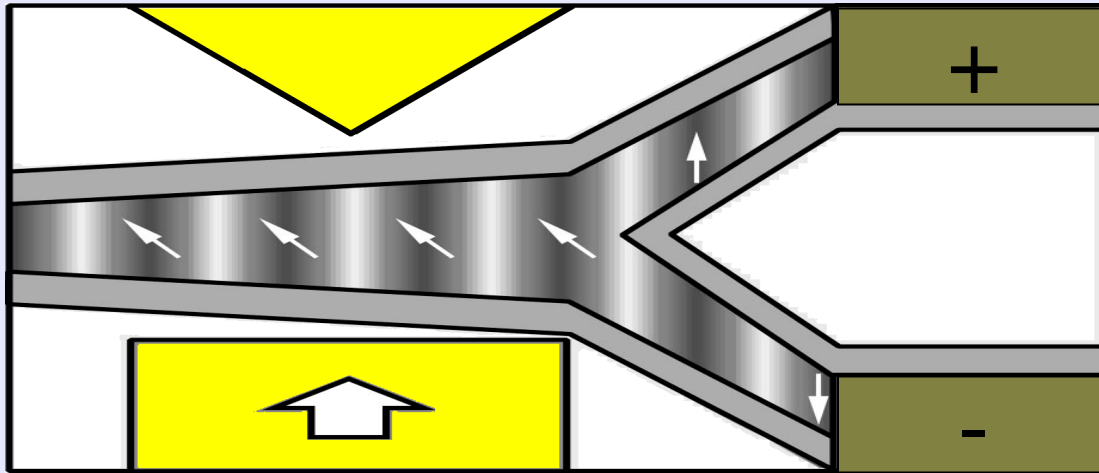


The indicated direction can be measured to arbitrary precision.

Quantum Orienteering

- Instead of gyroscopes, send individual spin-1/2 particles.
- Assume a particle is sent in direction \hat{n} .
- There is a finite amount of information that can be extracted, depending on Bob's measurement.

Orienteering with a single spin



Measure in the z-basis

Guess “up” or “down” based on outcome

$$\Rightarrow F_{1 \text{ spin}} = 2/3 \quad (1/2 \text{ without measurement})$$

Extendable to multiple state copies*:

For two spins, measure, e.g., z and x $\Rightarrow F_{\text{LOCC}} = 73.6\%$

*Massar, PRA **62**, 040101 (2000)

Orienteering with 2 identical spins

For 2 or more spins, *local* measurements are not optimal.

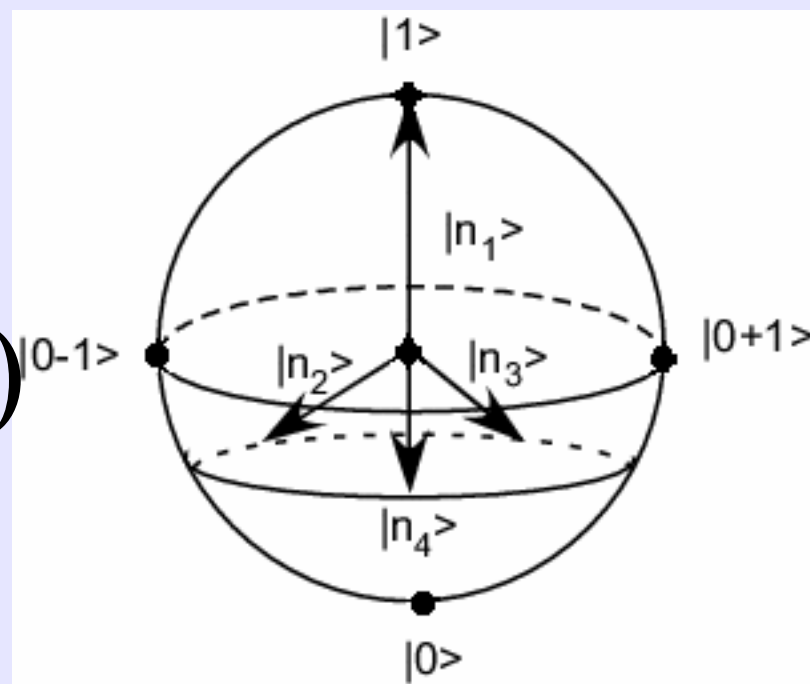
If Alice sends two *identical* spins $|\hat{n}, \hat{n}\rangle$ the best measurements* have eigenvectors:

$$|\psi_k\rangle = \frac{\sqrt{3}}{2} |\hat{n}_k, \hat{n}_k\rangle + \frac{1}{2} |\psi^{(-)}\rangle$$

$$|\psi^{(-)}\rangle \equiv \frac{1}{\sqrt{2}} (|\hat{n}_k, -\hat{n}_k\rangle - |-\hat{n}_k, \hat{n}_k\rangle)$$

$$\Rightarrow F = 3/4$$

$$(\text{For } N \text{ qubits: } F = \frac{N+1}{N+2})$$



*Massar & Popescu, PRL **74**, 1259 (1995)

Optimal orienteering with 2 spins

Surprisingly, even this is not optimal*...

Alice should flip the second spin and transmit $|\hat{n}, -\hat{n}\rangle$

The optimal measurement states are then**:

$$|\psi_k\rangle = \frac{\sqrt{3}}{2} \frac{|\hat{n}_k, -\hat{n}_k\rangle + |-\hat{n}_k, \hat{n}_k\rangle}{\sqrt{2}} + \frac{1}{2} |\psi^{(-)}\rangle$$

$$\Rightarrow F = 0.789$$

Why does it help to use anti-parallel spins? They reside in the full Hilbert space, i.e., singlet and triplet; the state $|\hat{n}, \hat{n}\rangle$ does not occupy the singlet sector of the space.

*Gisin & Popescu, PRL **83**, 432 (1999)

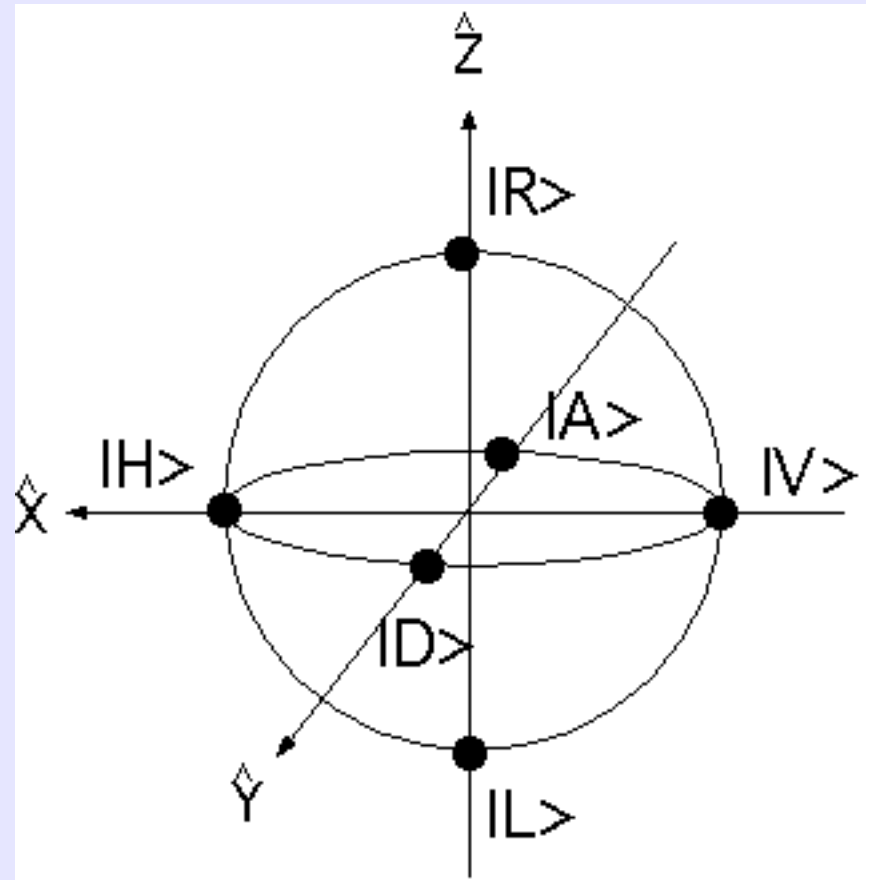
**Terry Rudolph, private comm.

Orienteering with photons

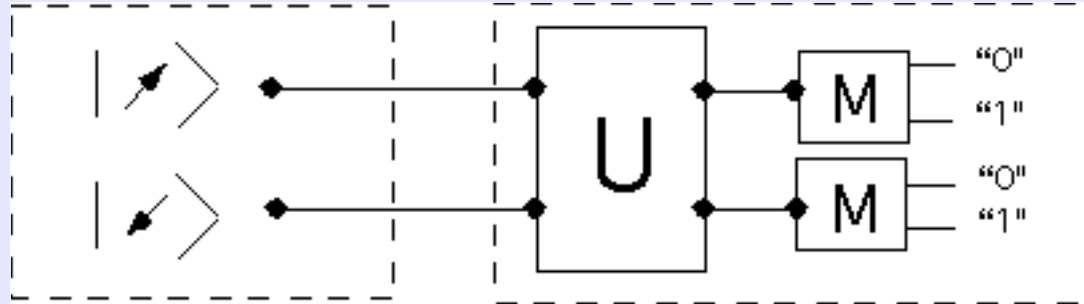
Photon polarization has no natural embedding in space. However, we can associate real-space directions with particular polarization states (directions on the Poincaré sphere).

Note: This requires using a shared reference frame.

Also, we must be wary -- the operation of wave plates and polarizers are k -vector dependent.

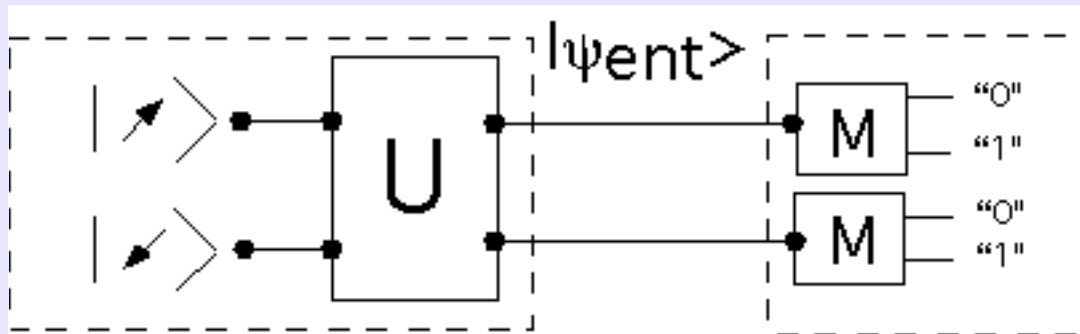


Conceptual considerations



- States are simple to create
- Measurements are difficult (U is a 2-qubit gate)

Because we have a shared reference frame, we can “move” the transformation*:

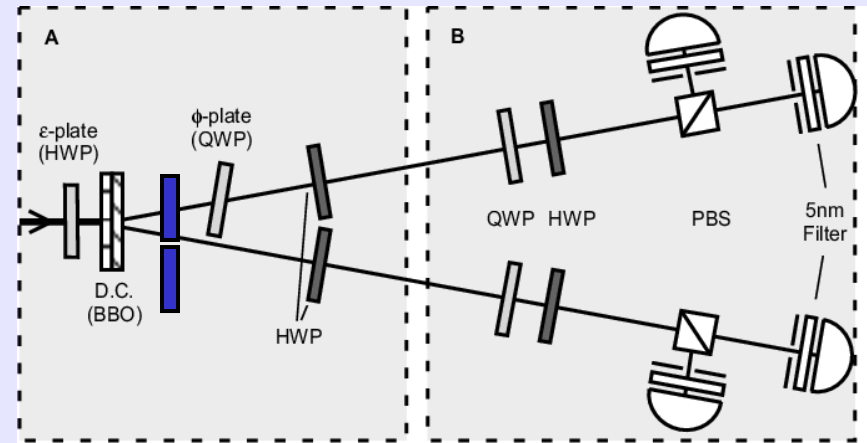


- Separable measurements are easy.
- Precise entangled state synthesis required.

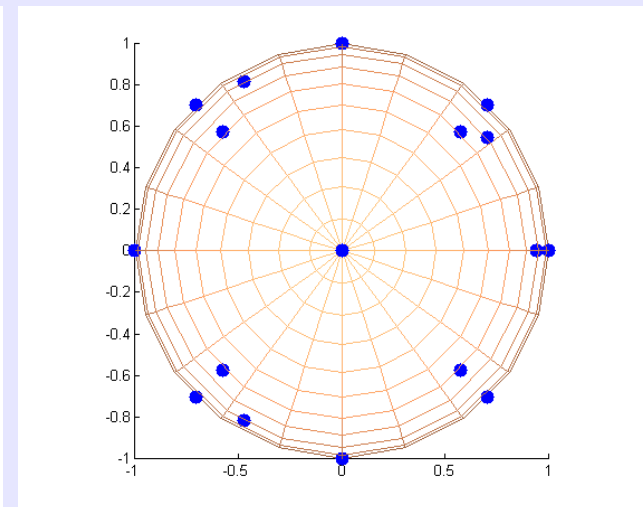
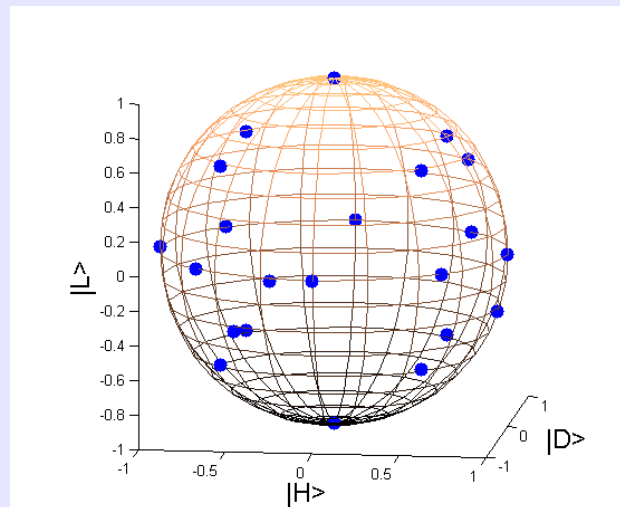
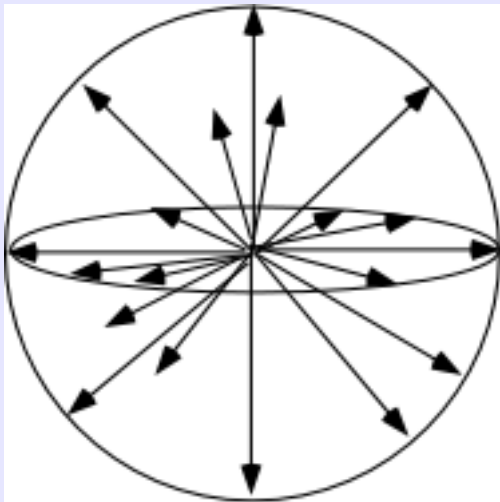
*T. Rudolph, private communication

Experimental configuration

Alice can create a wide variety of entangled states;
Bob can make arbitrary separable measurements.



Directions sent



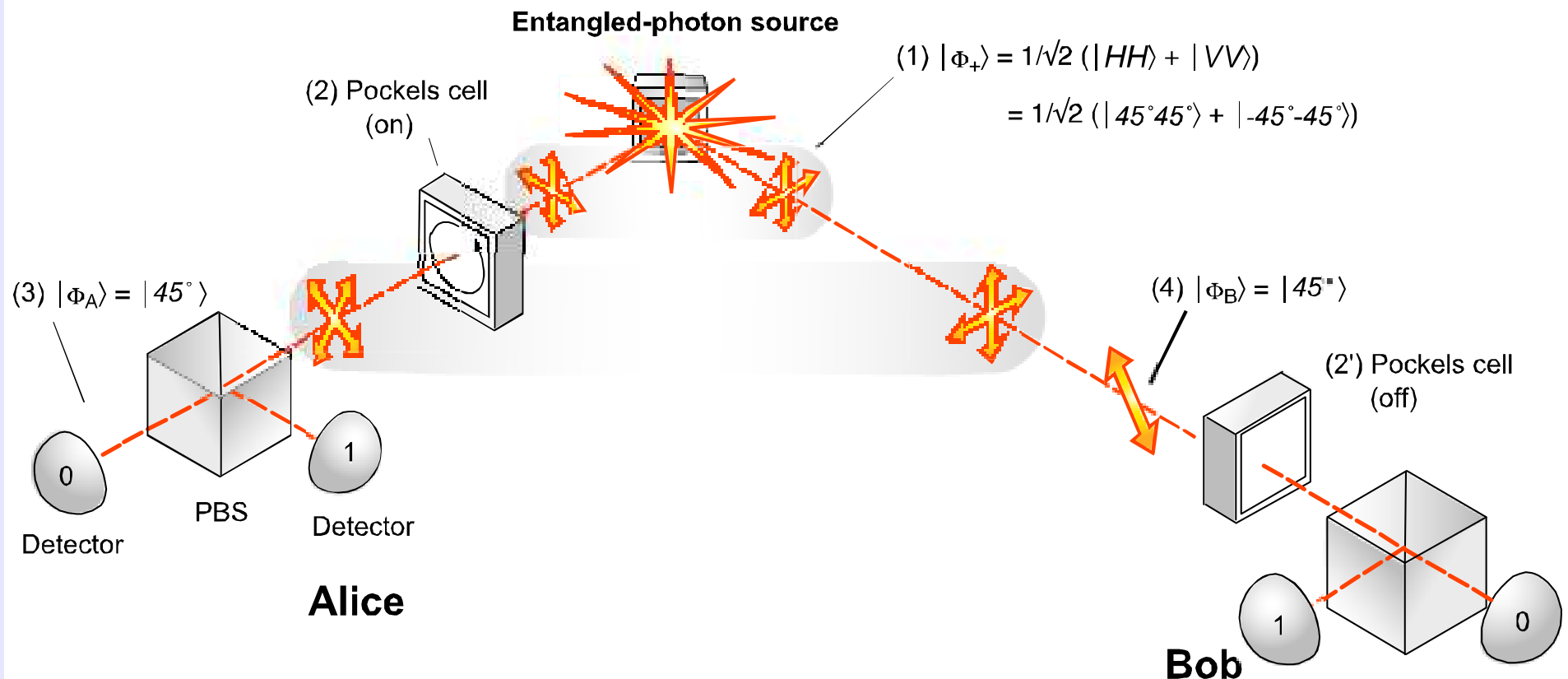
All states had $> 98\%$ state fidelity with the target state

Summary of average fidelities

Direction set	Separable (Th/Exp)
all	73.6/73.2

- Because photon spin is not “pointing” in real space, Alice and Bob require a shared reference frame.

Entangled-Photon Quantum Cryptography



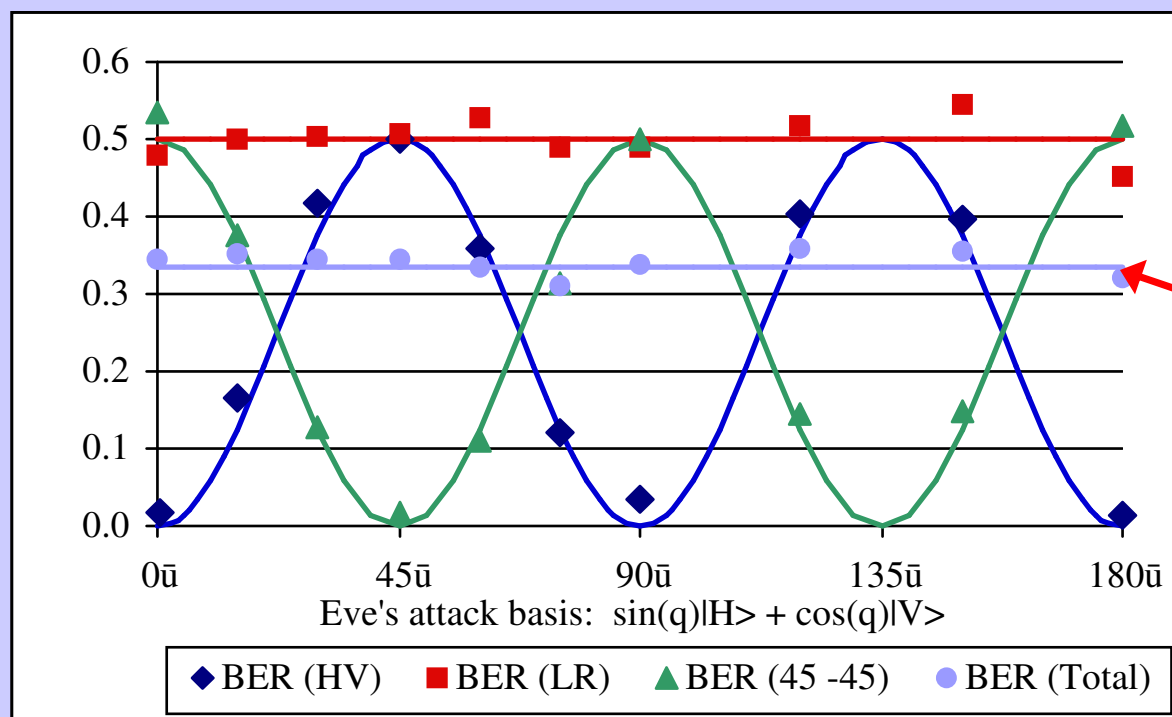
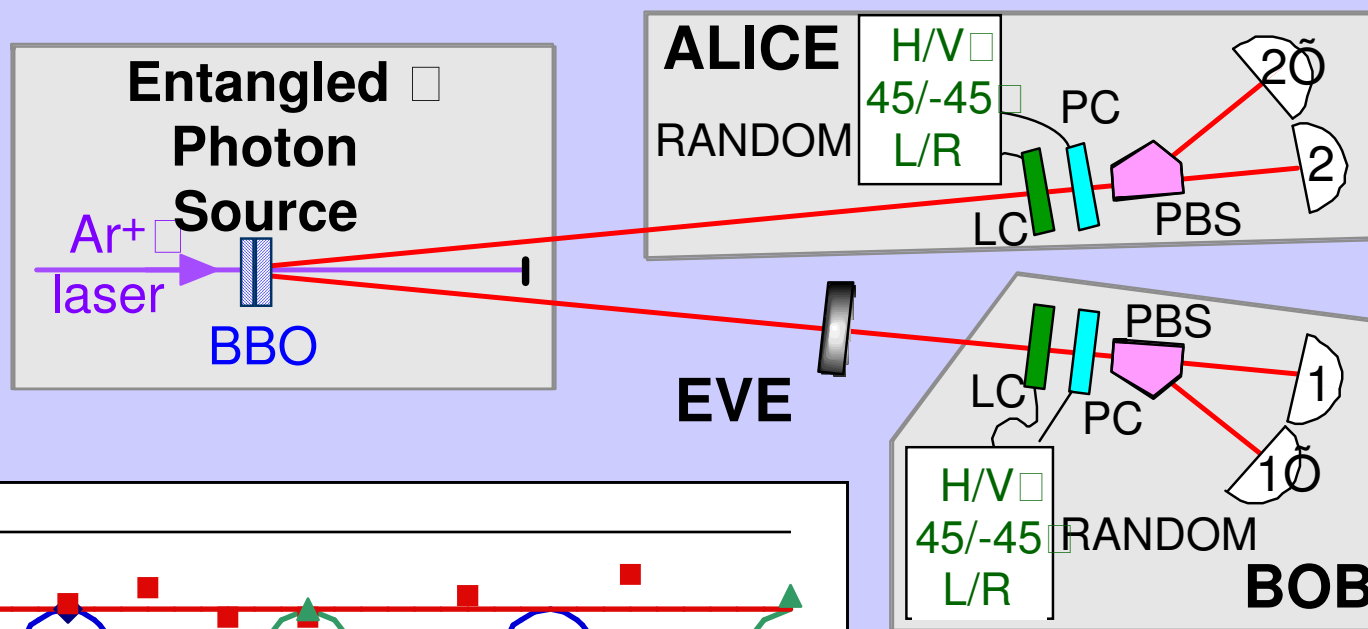
- Alice & Bob randomly measure polarization in the (HV) or the ($45^\circ 45^\circ$) basis.
- Discuss via a “public channel” which bases they used, *but not the results*.
- Discard cases (50%) where they used different bases → uncorrelated results.
- Keep cases where they used the same basis → *perfectly correlated results!*
- Define $H \equiv “0” \equiv 45^\circ$, $V \equiv “1” \equiv -45^\circ$. **They now share a secret key.**

Advantages of Entanglement

- In principle perfect correlations between Alice and Bob \Rightarrow well, not really perfect...
- Automatic randomness of key
- Longer distances accessible (since Bob can know *when* to look for a photon) [But decoy states...]
- Established methods to verify security of key
- Source can be automatically verified (even if “sold” by Evesdropper!)

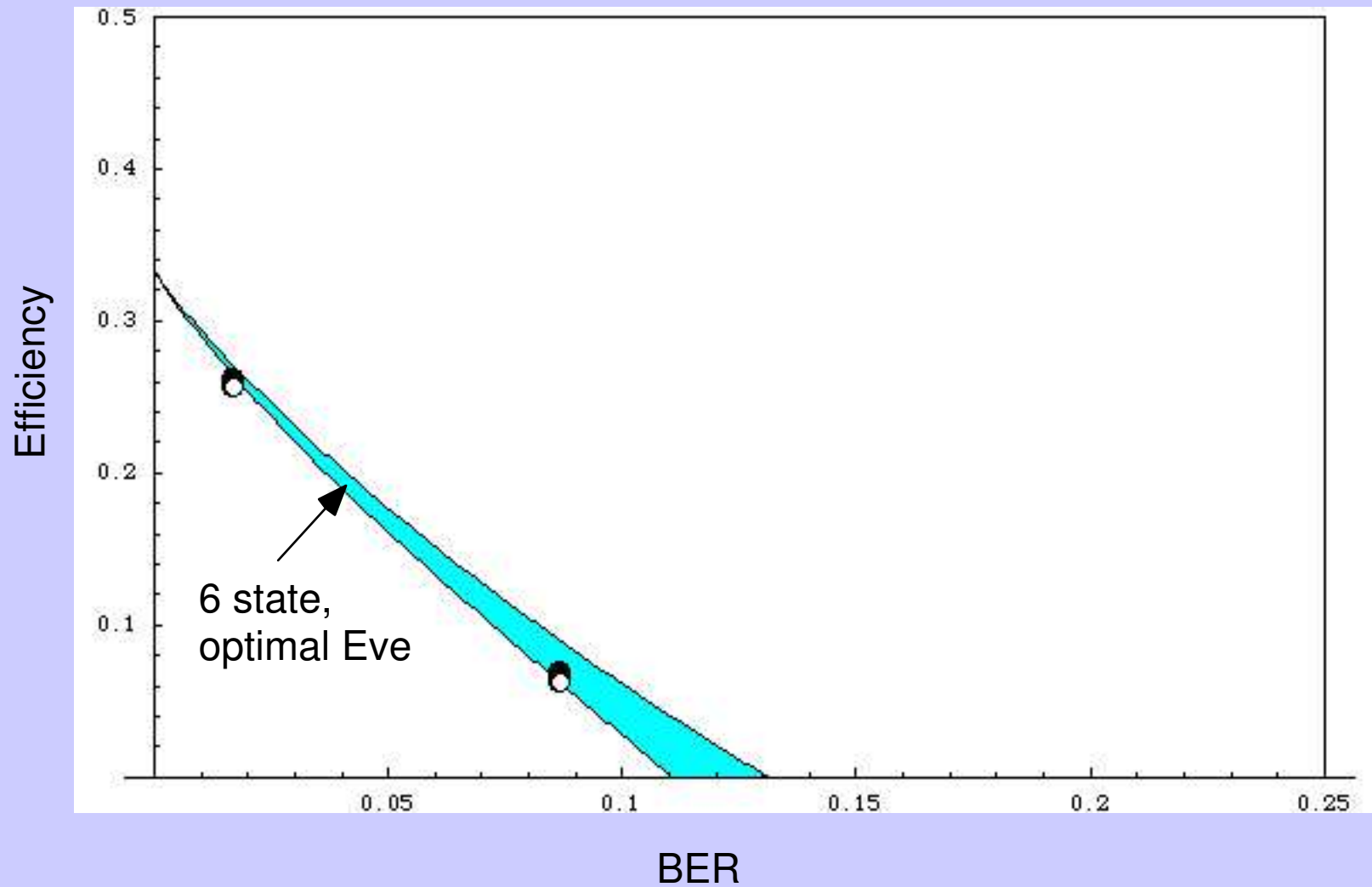
Experimental Realization of Six-State QKD Protocol

{D. Enzer et al., New Journal Physics 4, 45.1 (2002)}

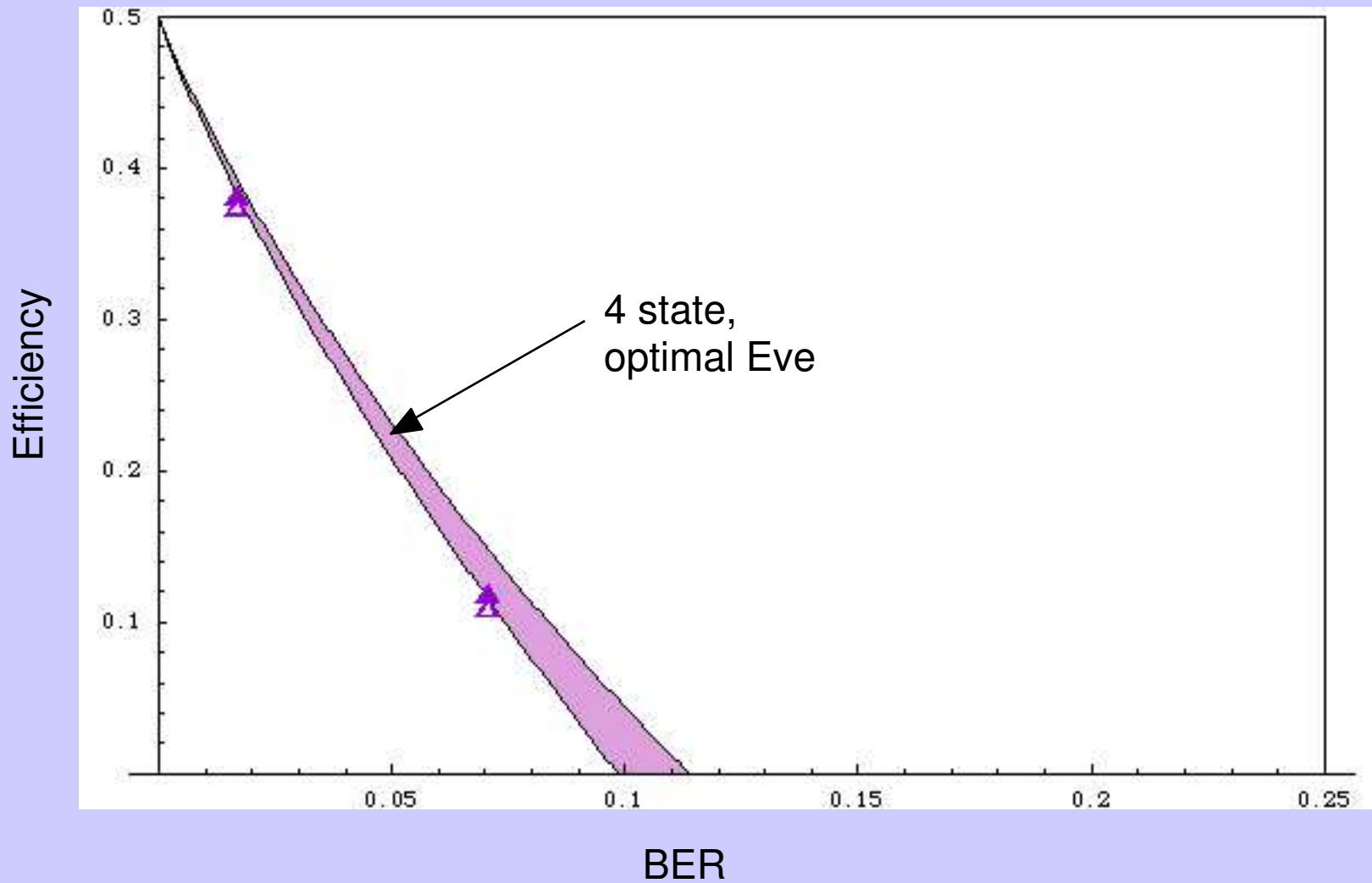


Total BER is 33%, independent of attack strategy
 (c.f. to 25% BER in BB84 4-state protocol)

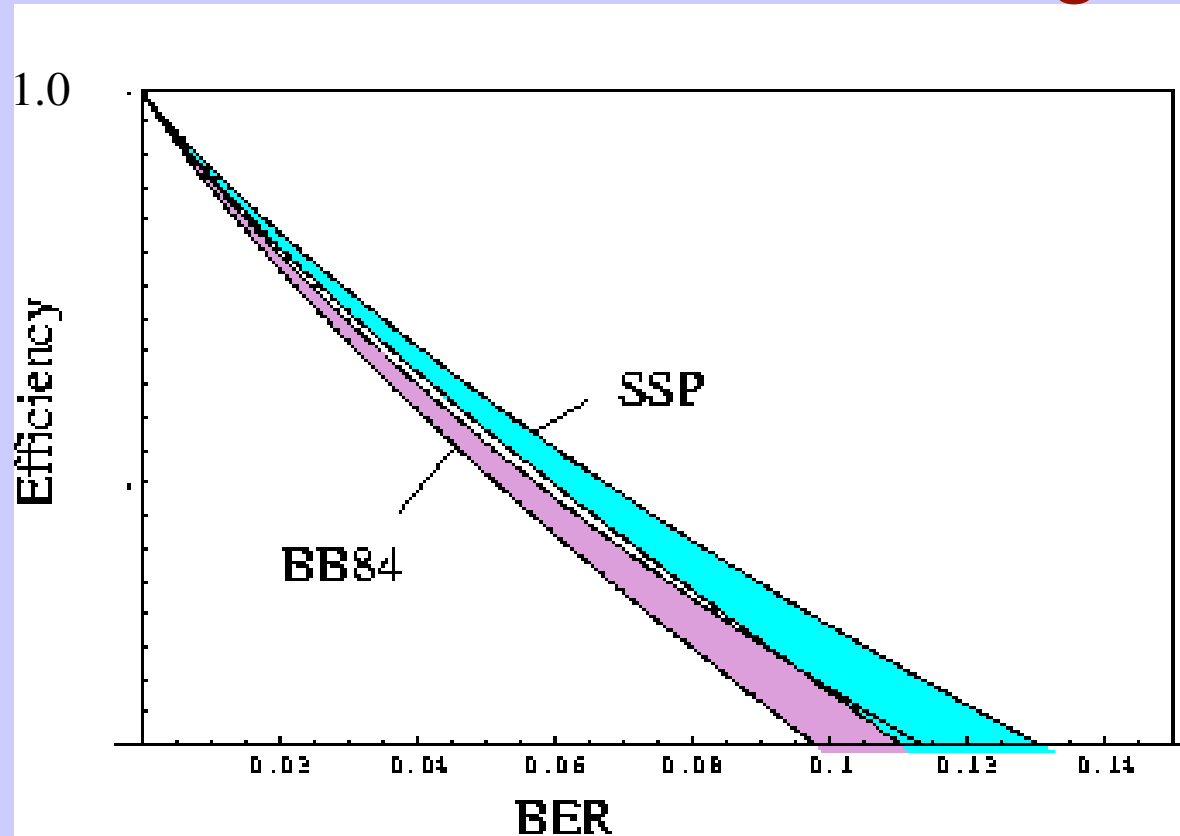
Bit yield, after Error Detection & Privacy Amplification



Bit yield, after Error Detection & Privacy Amplification



The Trouble with Sifting

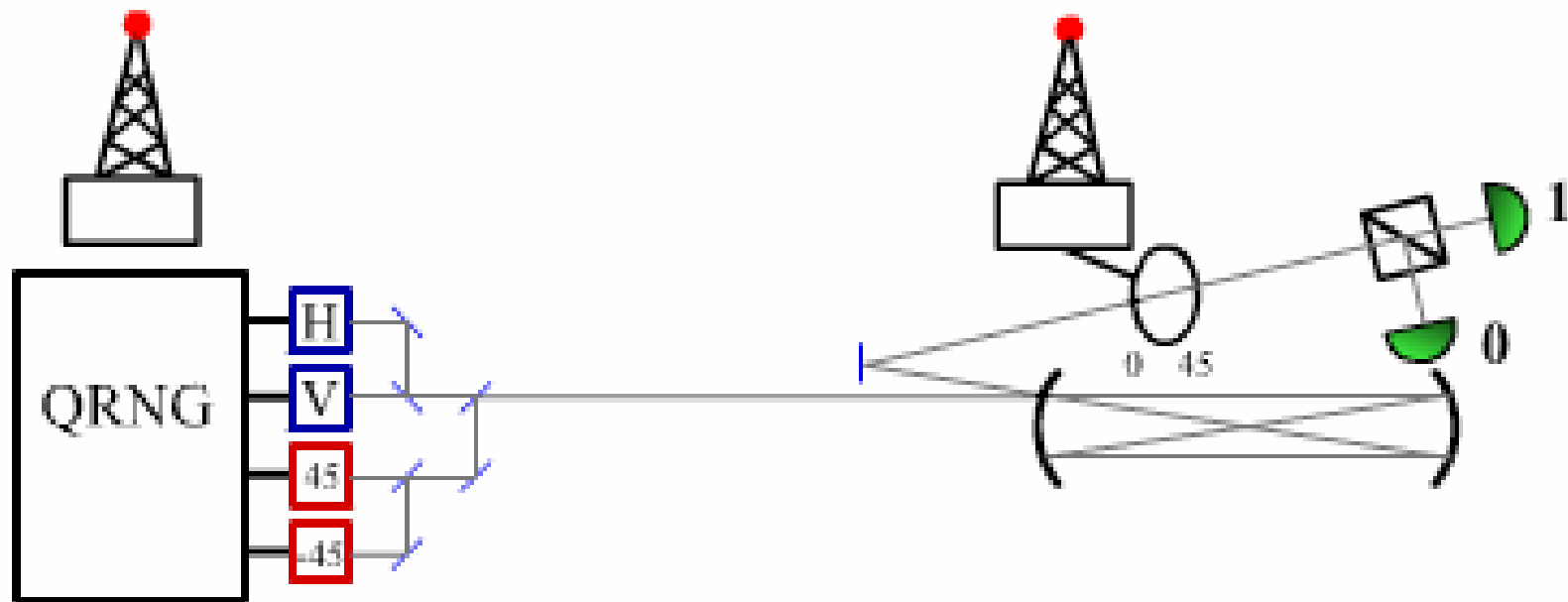


Eliminating the sifting \Rightarrow double efficiency of BB84
 \Rightarrow triple efficiency of SSP

- In principle, every photon contributes to key!
- SSP is always advantageous

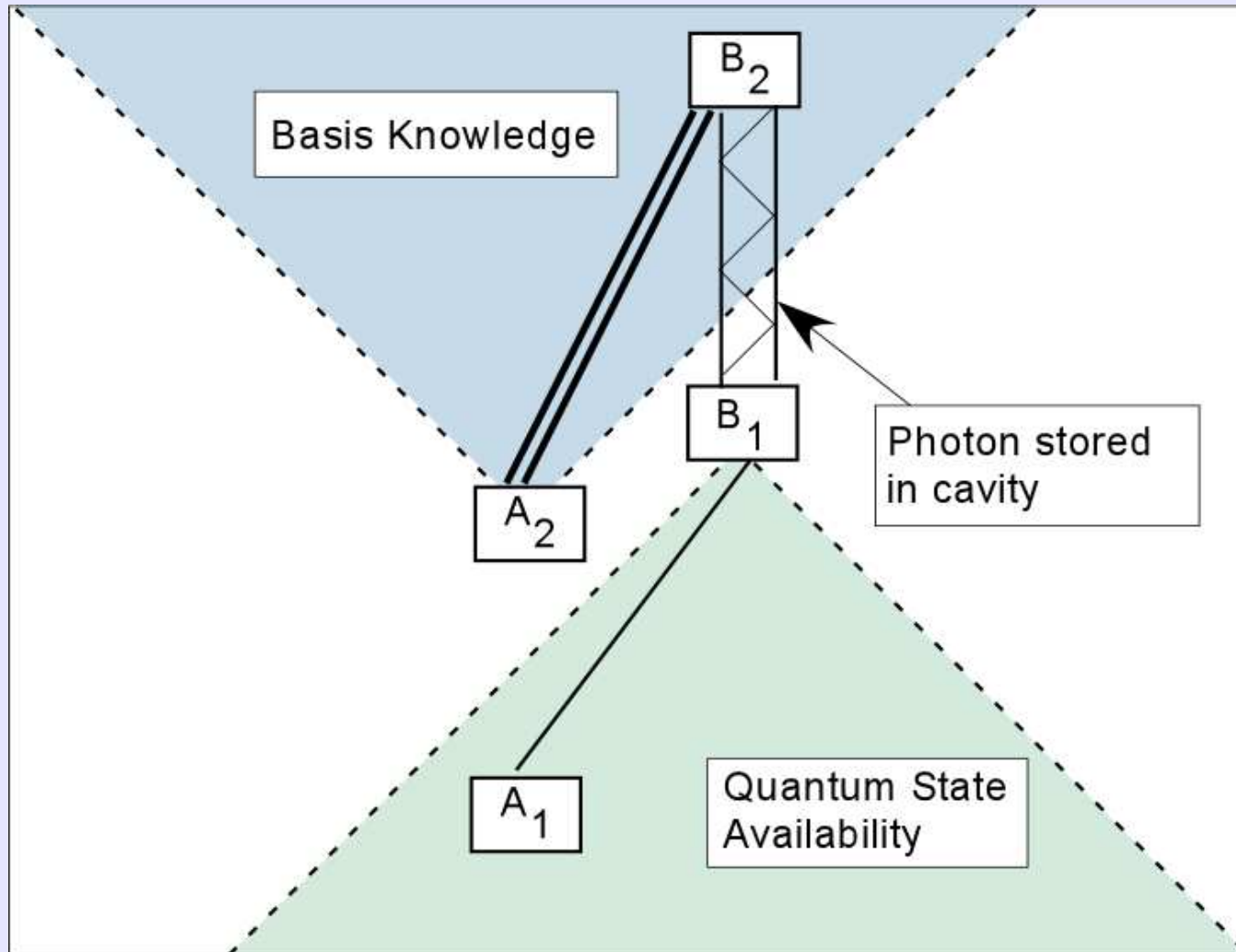
How can we eliminate sifting...

“Relativistic” Quantum Cryptography



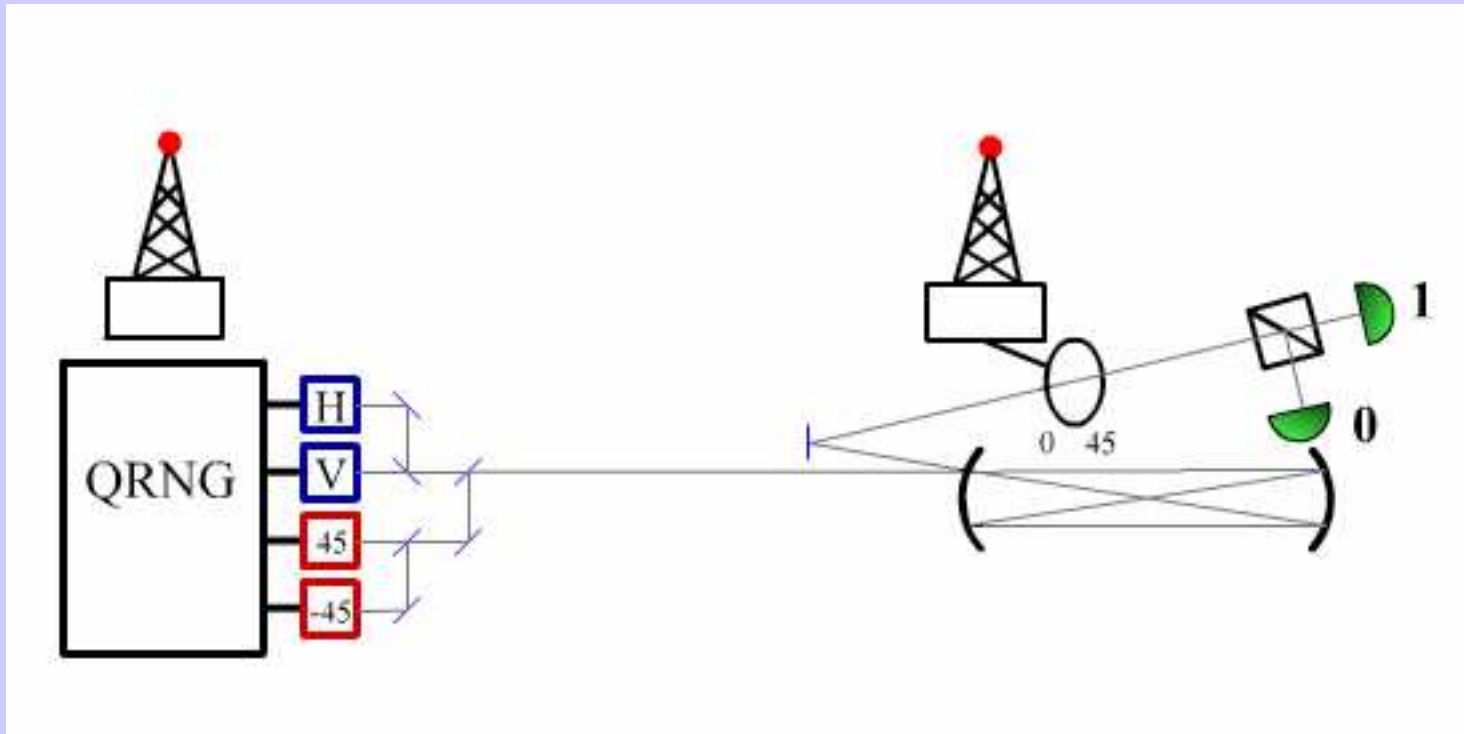
Bob stores each photon until Alice tells him which basis to use
à net efficiency is increased to 100% (in principle)
à same security as BB84 (Eve's p cannot depend on Bob)

QKD and Special Relativity



- These two light cones **must not** overlap
- A_2 may be before B_1 in some reference frames
- *Alice and Bob must know their space-time coordinates*

“Relativistic” Quantum Cryptography

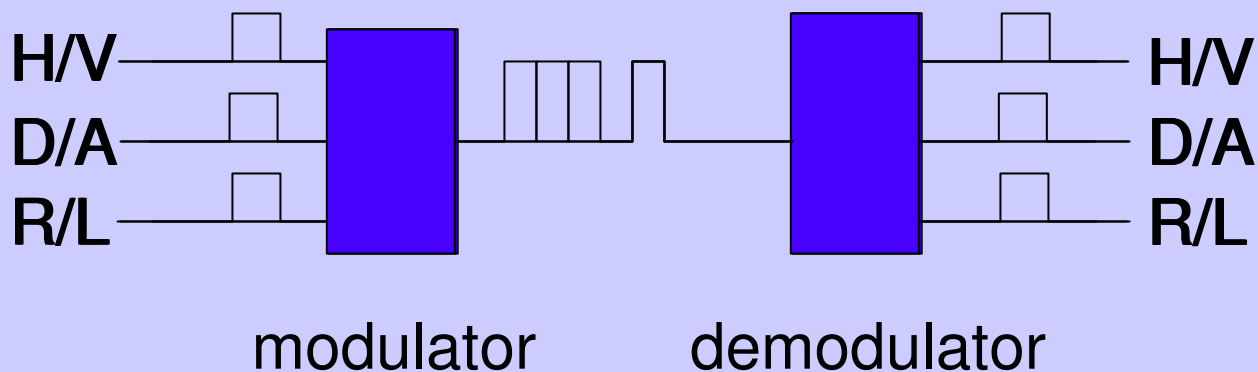


Special components
à fast classical modulation system
à quantum memory
à simpler 6-state analysis system

Modulation System

The classical basis information must be sent over a low-latency communication channel (since Bob can only store the photon for $\sim 1 \mu\text{s}$).

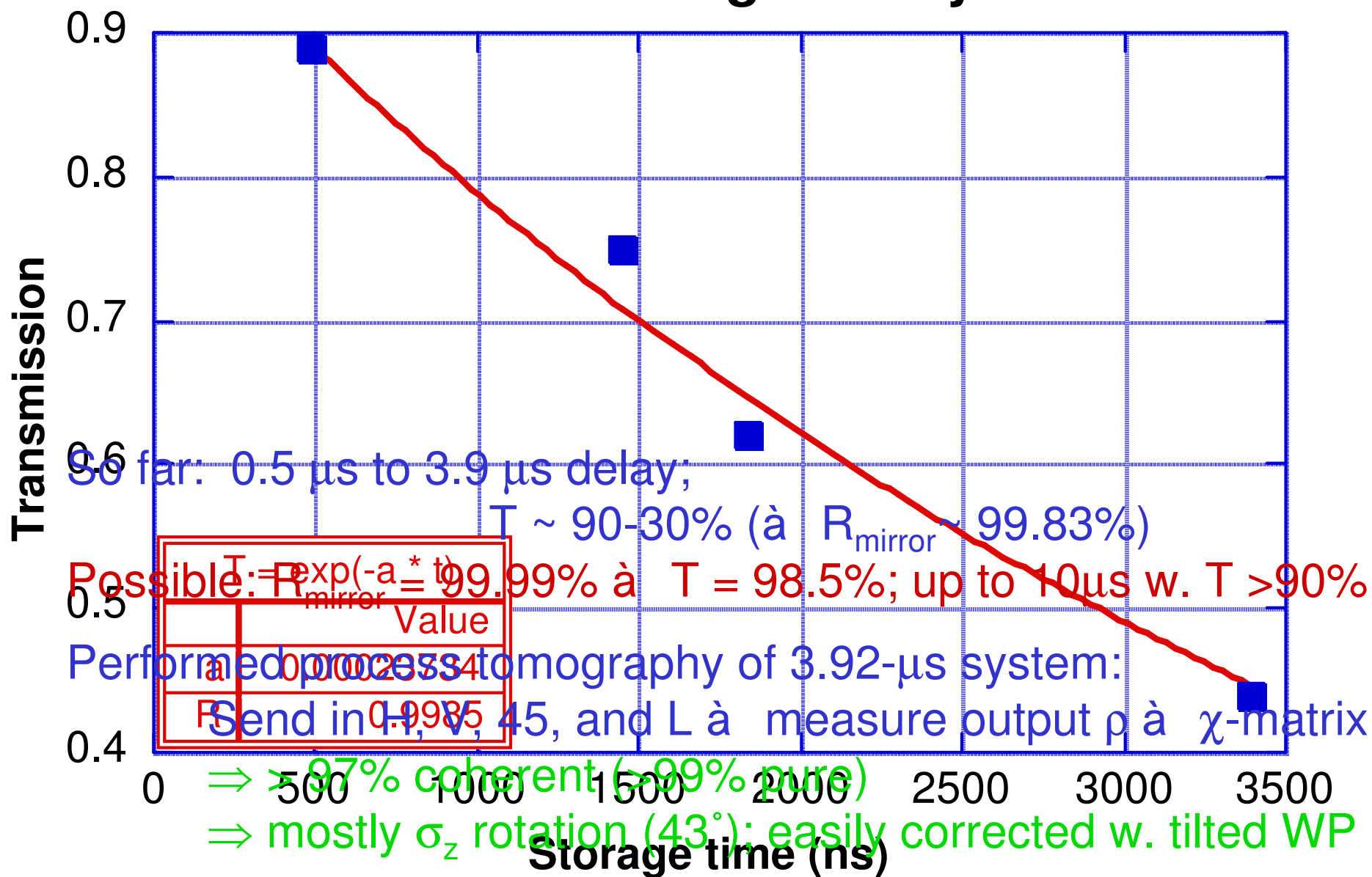
We implement a finite state machine using fast programmable logic (CPLD) to drive a diode laser:

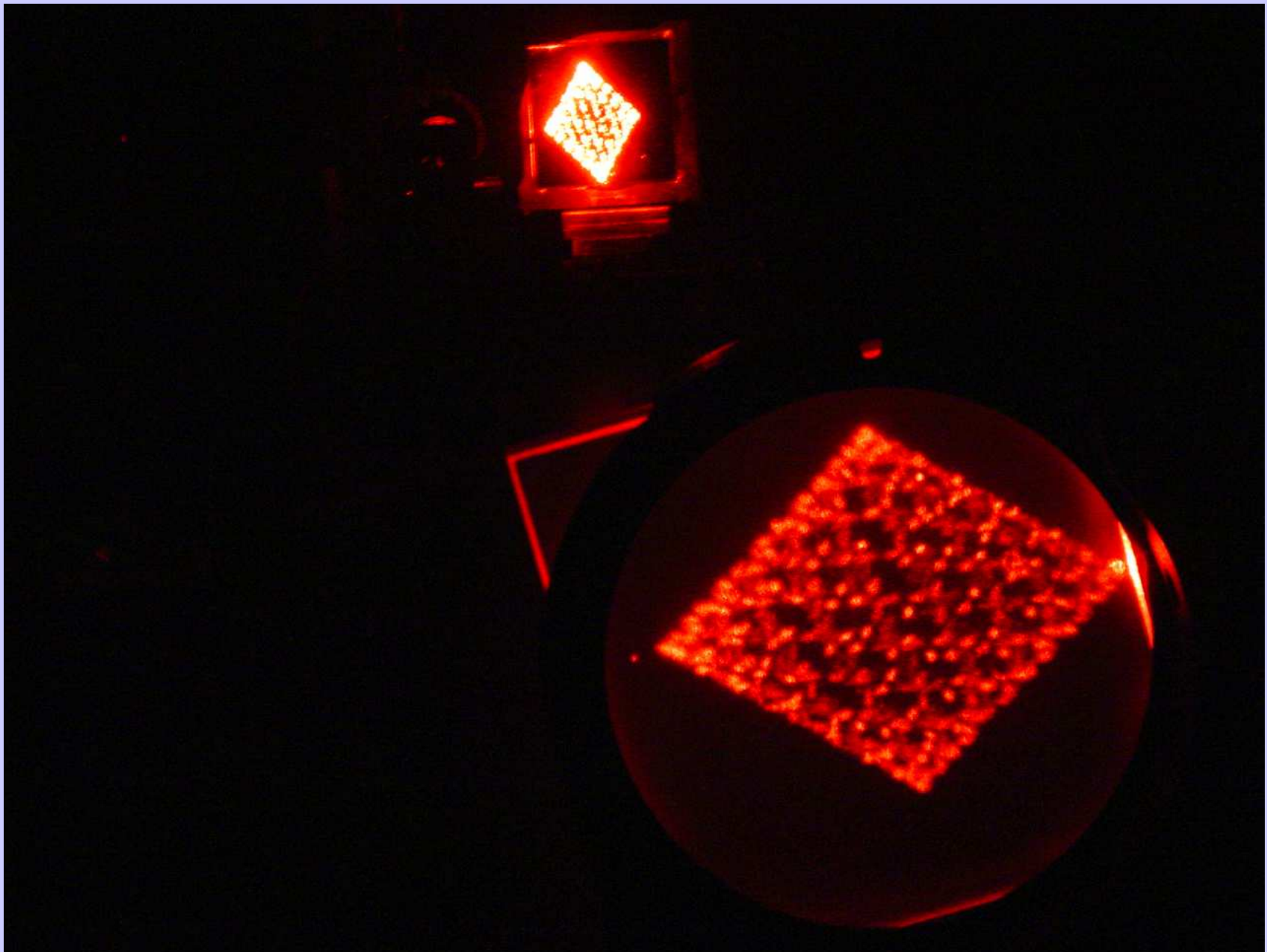


Clock rate = 80 MHz

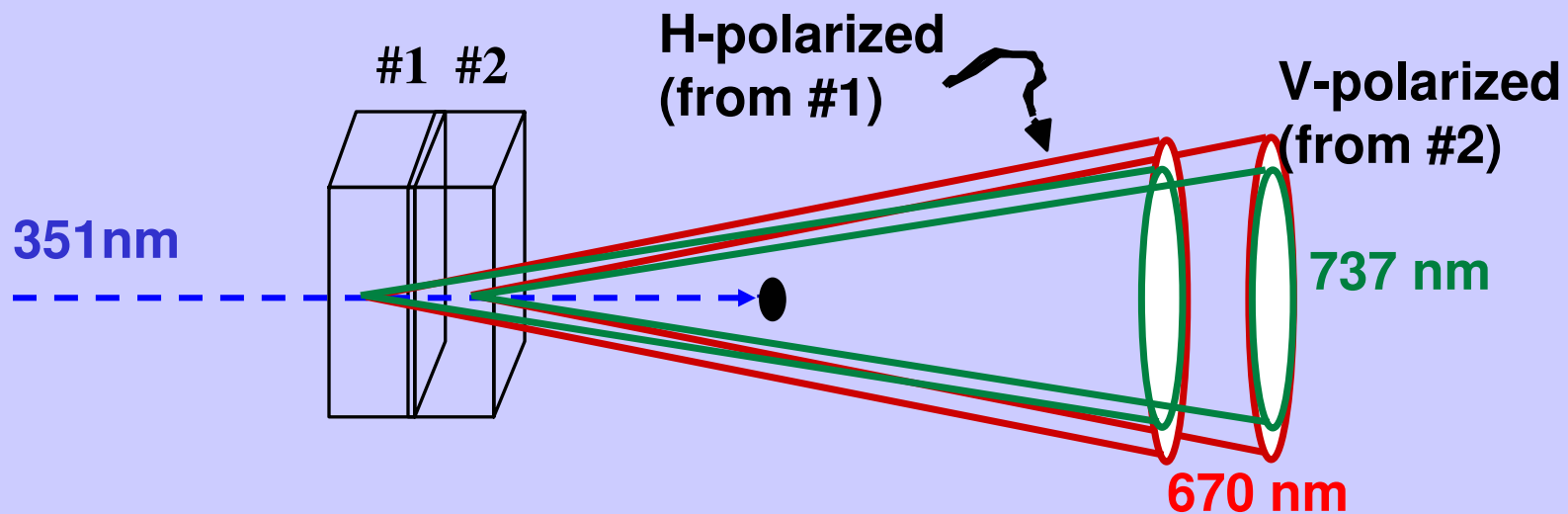
Total latency < 120 ns

Photon Storage Cavity

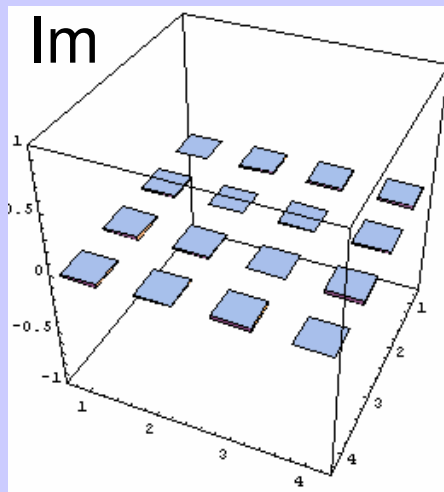
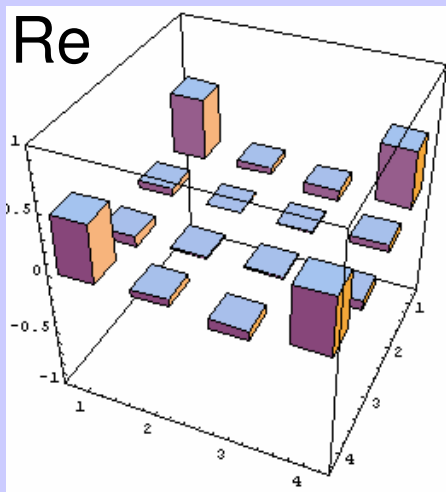




Non-degenerate Polarization-entanglement



$$|\psi\rangle = \frac{1}{\sqrt{2}} (|H\rangle |H\rangle + e^{i\varphi} |V\rangle |V\rangle)$$



$$F_{\text{with HH+VV}} = 0.968$$

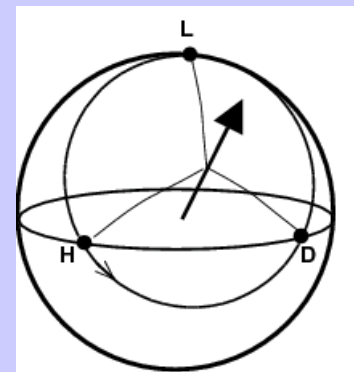
$$\text{Linear entropy} = 0.027$$

$$\text{Tangle} = 0.957$$

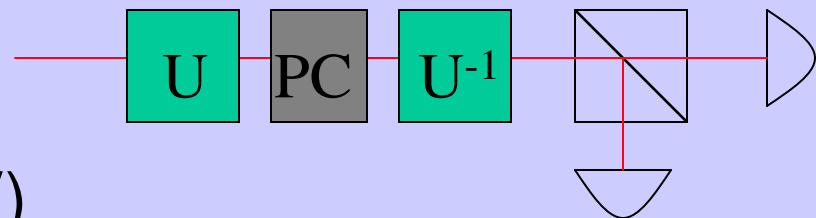
Six-State Polarization Analysis

- Measuring in three bases typically requires two electro-optic devices
- We desire a “Minimum parts-count analyzer”

Try rotating about the $\{1,1,1\}$ axis on Poincaré sphere



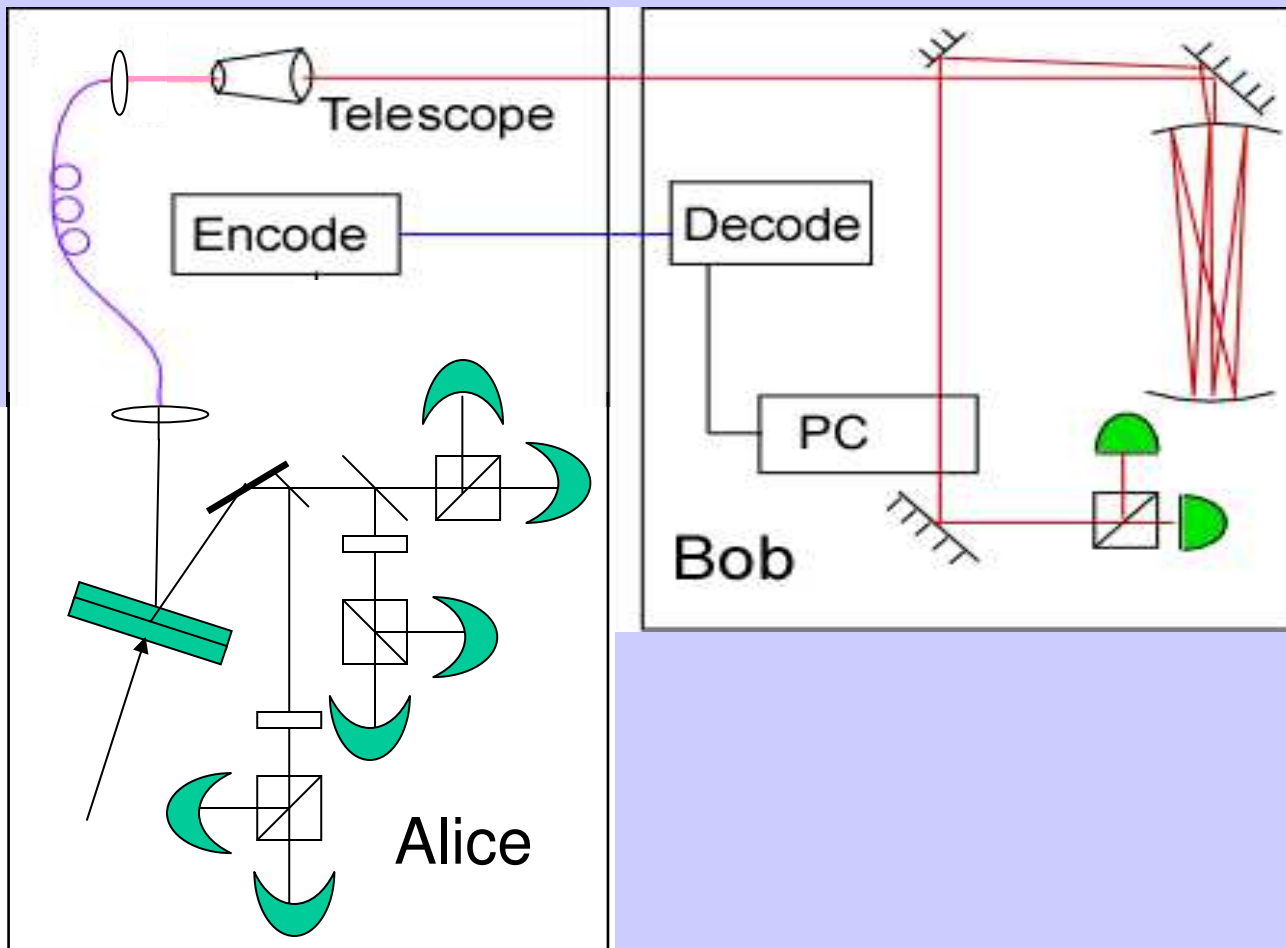
Use fixed waveplates plus **one** Pockels cell (with *three* voltages: 0, $\pm V$)



In fact, because of polarization entanglement, Alice can set the bases by which measurements *she* uses (c.f., Remote State Preparation)

Incorporate *entangled* photon source

Non-degenerate polarization-entangled state
(351 nm \rightarrow 670 nm + 737 nm)

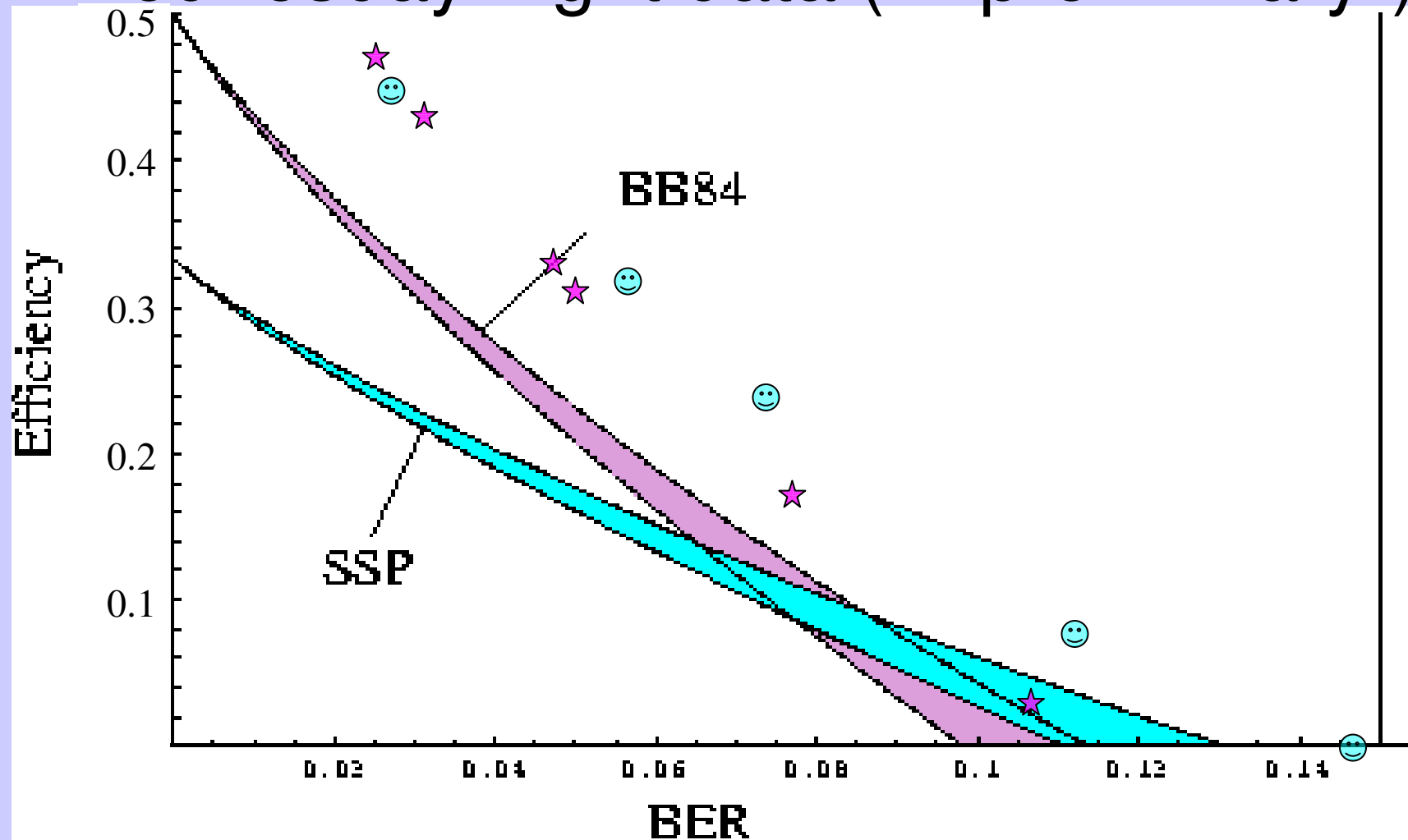


BB84, 30 mW pump power
94 sifted bits/second
2.5% error rate
 \rightarrow 65.5 secret bits/second

BB84, 90 mW pump power
214 bits/second
3.1% error rate
 \rightarrow 136 secret bits/second
 \rightarrow **yield enhancement = 1.3**

SSP, 90 mW pump power
371 bits/second
2.7% error rate
 \rightarrow 251 bits/second
 \rightarrow **yield enhancement = 2.1**

Wednesday night data (= “preliminary”)



- Cavity loss ($\sim 33\%$) prevents yield from exceeding $1/2$.
- Rate is ...non-optimal (~ 50 - $100/s$).
- Six-state protocol now advantageous if $BER > \sim 4\%$.

Relativistic QKD: Summary

- By allowing photon-storage, Bob can use the correct basis for every measurement \Rightarrow enhanced yield by 2 (BB84) or 3 (6-state)
- Security constraint = “why it’s relativistic”
 - non-overlapping light cones of Bob’s receipt of photon and Alice’s classical basis transmission
 - Alice and Bob must know their relative space-time coordinates!
- New technologies
 - Twisted cylindrical-mirror cavity (450ns with $T = 67\%$) [CLEAN mirrors!]
 - Low latency classical modulation sender/receiver (CPLD logic, < 120 ns).
 - 3-basis analysis with single Pockel cell (three voltages: 0, $\pm V$).
- Non-degenerate polarization-entangled source
 - Preliminary yield enhancements of ~ 1.3 (BB84) and ~ 2.1 (SSP)
 - Mirror cleaning may/should improve this (up to 1.6 [BB84] and 2.6 [SSP])
- First demo. of Überquantum advantage, i.e., $QM + SR > QM$
 - Next steps: rates, stability, eavesdropping...