

Holography and hydrodynamics

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Collaboration and some Refs.

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Prologue

- ∇ Our goal is to understand thermal gauge theories, e.g. thermal QCD
- ∇ Of particular interest is the regime described by fluid dynamics, e.g. quark-gluon plasma
- ∇ This near-equilibrium regime is completely characterized by values of transport coefficients, e.g. shear and bulk viscosity
- ∇ Transport coefficients are hard to compute from “first principles”, even in perturbation theory. For example, no perturbative calculation of bulk viscosity in gauge theory is available.

Prologue (continued)

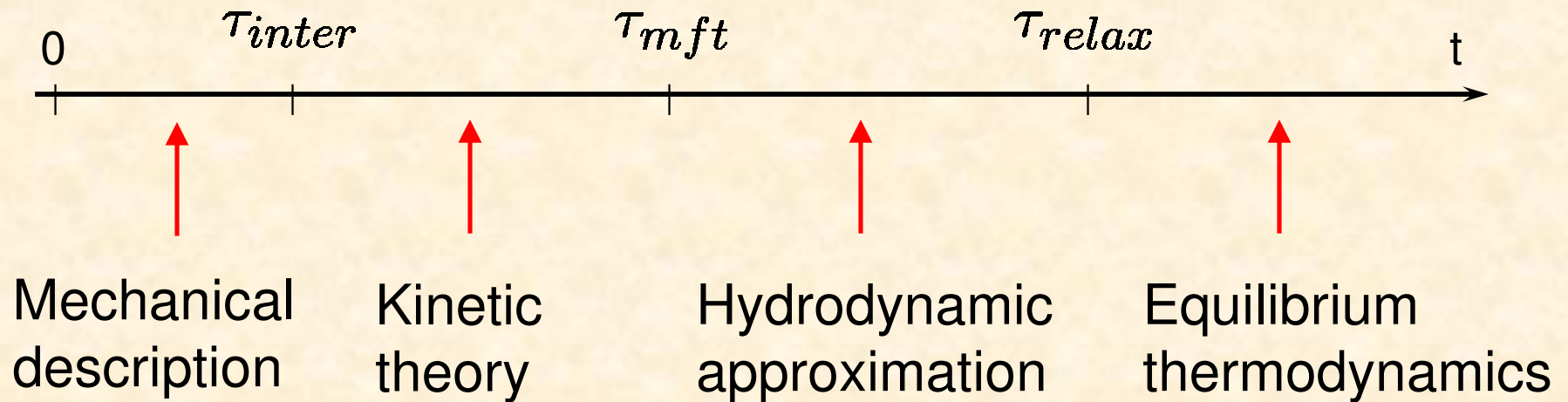
- ∇ Transport coefficients of some gauge theories can be computed in the regime described by string (gravity) duals – usually at large N and large 't Hooft coupling
- ∇ Corrections can in principle be computed
- ∇ Shear viscosity result is *universal*. Model-independent results may be of relevance for RHIC physics
- ∇ Certain results are *predicted* by hydrodynamics. Finding them in gravity provides a check of the AdS/CFT conjecture

Outline

- Ø Shear and bulk viscosities
- Ø Computing viscosity in AdS/CFT
- Ø Universality of shear viscosity in the regime described by gravity duals
- Ø Bulk viscosity and the speed of sound from gravity
- Ø Relation to RHIC (if any)
- Ø Conclusions

What is hydrodynamics?

Hierarchy of times (example)



Hierarchy of scales

$$l_{mfp} \ll l \ll L$$

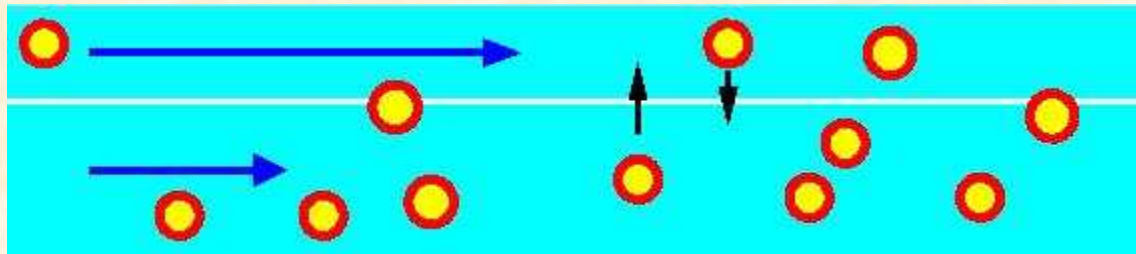
(L is a macroscopic size of a system)

What is viscosity?

Friction in Newton's equation: $\frac{d(mv_i)}{dt} + \gamma v_i = F_i$

Friction in Euler's equations

$$\frac{\partial(\rho v_i)}{\partial t} = -\frac{\partial}{\partial x^k} (P\delta_{ik} + \rho v_i v_k) + \frac{\partial}{\partial x^k} \sigma_{ik}^{fric}$$



$$\sigma_{ik}^{fric} \sim \partial v_i / \partial x^k$$

$$\sigma_{ik}^{fric} \sim \partial v_i / \partial x^k + \partial v_k / \partial x^i$$

$$\sigma_{ik}^{fric} = \eta \left(\frac{\partial v_i}{\partial x^k} + \frac{\partial v_k}{\partial x^i} - \frac{2}{d} \delta_{ik} \frac{\partial v_l}{\partial x^l} \right) + \zeta \delta_{ik} \frac{\partial v_l}{\partial x^l} + \dots$$

Viscosity of gases and liquids

Gases (Maxwell, 1867): $\eta \sim \rho \bar{v} l_{mfp} \sim \frac{m_o \bar{v}}{\sigma} \sim \frac{m_o^{1/2}}{\sigma} \sqrt{T}$

Viscosity of a gas is

- § independent of pressure

- § scales as square of temperature

- § inversely proportional to cross-section

Liquids (Frenkel, 1926): $\eta \sim A(P, T) \exp \frac{W}{T}$

- § W is the “activation energy”

- § In practice, A and W are chosen to fit data

Computing transport coefficients from “first principles”

Fluctuation-dissipation theory
(Callen, Welton, Green, Kubo)

Kubo formulae allows one to calculate transport coefficients from microscopic models

$$\eta = \lim_{\omega \rightarrow 0} \frac{1}{2\omega} \int dt d^3x e^{i\omega t} \langle [T_{xy}(t, x), T_{xy}(0, 0)] \rangle$$

In the regime described by a gravity dual
the correlator can be computed using
AdS/CFT

Universality of shear viscosity in the regime described by gravity duals

$$ds^2 = f(w) (dx^2 + dy^2) + g_{\mu\nu}(w) dw^\mu dw^\nu$$

$$\left. \begin{aligned} \eta &= \lim_{\omega \rightarrow 0} \frac{1}{2\omega} \int dt dx e^{i\omega t} \langle [T_{xy}(t, x), T_{xy}(0, 0)] \rangle \\ \sigma_{abs} &= -\frac{16\pi G}{\omega} \text{Im } G^R(\omega) \\ &= \frac{8\pi G}{\omega} \int dt dx e^{i\omega t} \langle [T_{xy}(t, x), T_{xy}(0, 0)] \rangle \end{aligned} \right\} \eta = \frac{\sigma_{abs}(0)}{16\pi G}$$

Graviton's component h_y^x obeys equation for a minimally coupled massless scalar. But then $\sigma_{abs}(0) = A_H$.

Since the entropy (density) is $s = A_H/4G$ we get

$$\frac{\eta}{s} = \frac{1}{4\pi}$$

Three roads to universality of η/s

∅ The absorption argument

D. Son, P. Kovtun, A.S., hep-th/0405231

∅ Direct computation of the correlator in Kubo formula from AdS/CFT

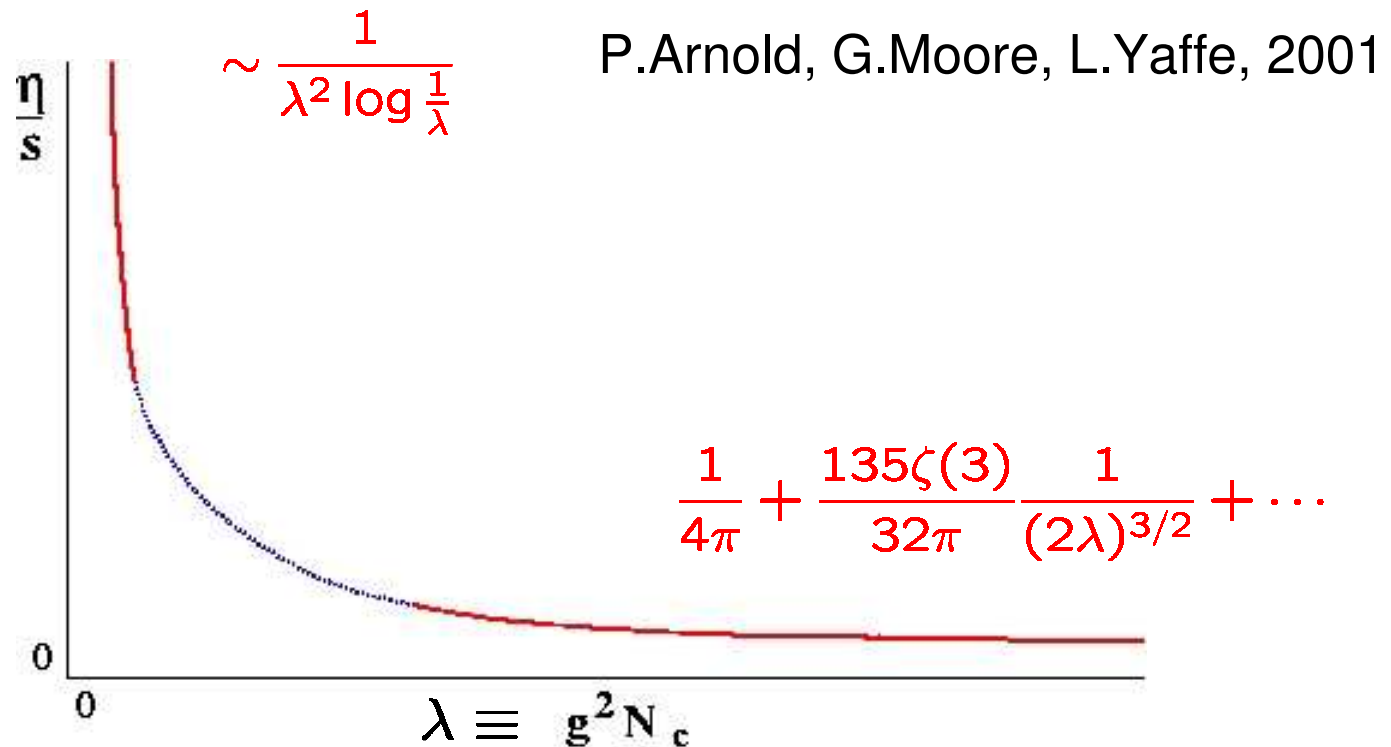
A.Buchel, hep-th/0408095

∅ “Membrane paradigm” general formula for diffusion coefficient + interpretation as lowest quasinormal frequency = pole of the shear mode correlator + Buchel-Liu theorem

P. Kovtun, D.Son, A.S., hep-th/0309213, A.S., to appear,

P.Kovtun, A.S., hep-th/0506184, A.Buchel, J.Liu, hep-th/0311175

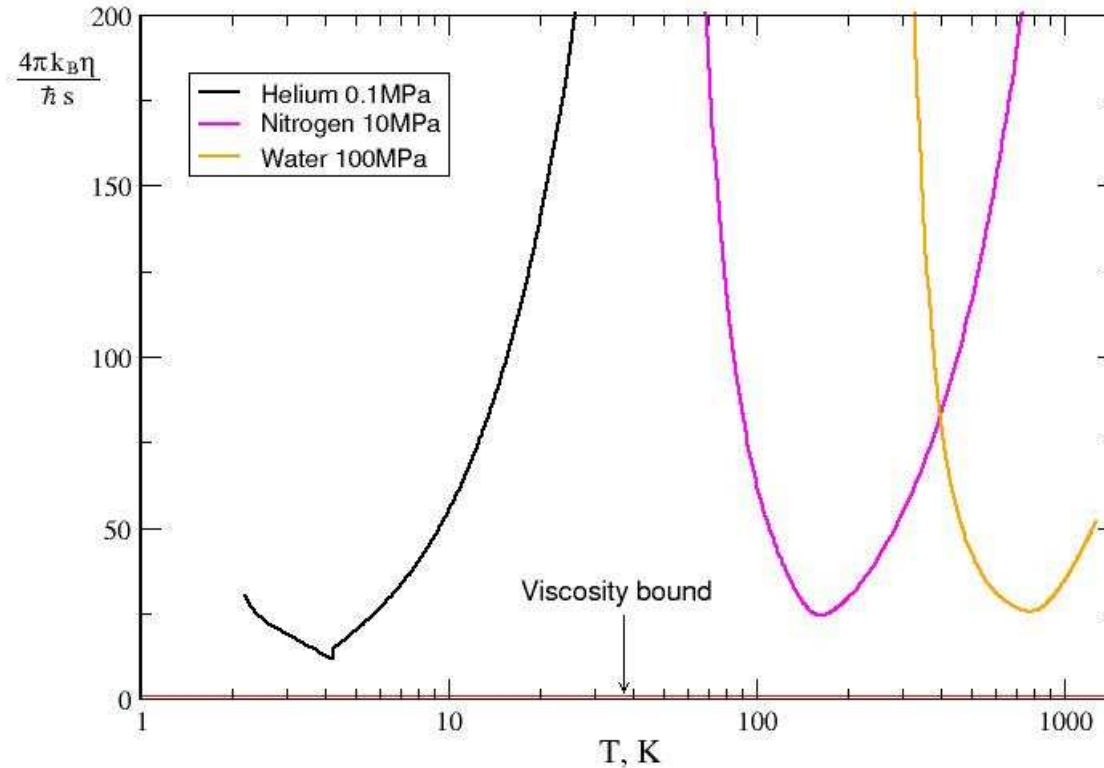
Shear viscosity in $\mathcal{N} = 4$ SYM



Correction to $1/4\pi$: A. Buchel, J. Liu, A.S., hep-th/0406264

A viscosity bound conjecture

$$\frac{\eta}{s} \geq \frac{\hbar}{4\pi k_B} \approx 6.08 \cdot 10^{-13} \text{ K} \cdot \text{s}$$

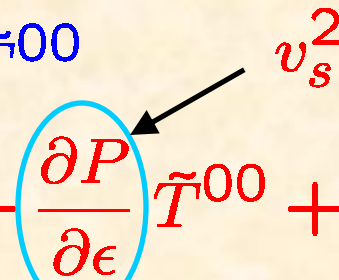


P.Kovtun, D.Son, A.S., hep-th/0309213, hep-th/0405231

Hydrodynamics as an effective theory

Thermodynamic equilibrium: $\langle T^{00} \rangle = \epsilon$, $\langle T^{0i} \rangle = 0$
 $T^{ij} = P(\epsilon) \delta^{ij}$

Near-equilibrium: $T^{00} = \epsilon + \tilde{T}^{00}$
 $T^{ij} = P \delta^{ij} + \frac{\partial P}{\partial \epsilon} \tilde{T}^{00} + \tilde{T}^{ij}$



$$\tilde{T}^{ij} = -\frac{1}{\epsilon + P} \left[\eta \left(\partial_i \tilde{T}^{0j} + \partial_j \tilde{T}^{0i} - \frac{2}{3} \delta^{ij} \partial_k \tilde{T}^{0k} \right) + \zeta \delta^{ij} \partial_k \tilde{T}^{0k} \right] + \dots$$

Eigenmodes of the system of equations $\partial_\mu T^{\mu\nu} = 0$

Shear mode (transverse fluctuations of \tilde{T}^{0i}): $\omega = -\frac{i\eta}{\epsilon + P} q^2$

Sound mode: $\omega = v_s q - \frac{i}{2} \frac{1}{\epsilon + P} \left(\zeta + \frac{4}{3} \eta \right) q^2$

For CFT we have $\zeta = 0$ and $\epsilon = 3P \implies v_s = 1/\sqrt{3}$

Two-point correlation function of stress-energy tensor

Field theory

Zero temperature: $\langle T_{\mu\nu} T_{\alpha\beta} \rangle_{T=0} = \Pi_{\mu\nu,\alpha\beta} F(k^2) + Q_{\mu\nu,\alpha\beta} G(k^2)$

Finite temperature: $\langle T_{\mu\nu} T_{\alpha\beta} \rangle_T = S_{\mu\nu,\alpha\beta}^{(1)} G_1(\omega, q) + S_{\mu\nu,\alpha\beta}^{(2)} G_2(\omega, q) + S_{\mu\nu,\alpha\beta}^{(3)} G_3(\omega, q) + S_{\mu\nu,\alpha\beta}^{(4)} G_4 + S_{\mu\nu,\alpha\beta}^{(5)} G_5$

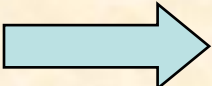
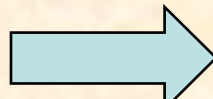
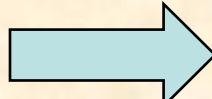
Dual gravity

- Ø Five gauge-invariant combinations Z_1, Z_2, Z_3, Z_4, Z_5 of $h_{\mu\nu}$ and other fields determine G_1, G_2, G_3, G_4, G_5
- Ø Z_1, Z_2, Z_3, Z_4, Z_5 obey a system of coupled ODEs
- Ø Their (quasinormal) spectrum determines singularities of the correlator

Classification of fluctuations and universality

$$ds^2 = \frac{r^2}{R^2} \left(-f(r)dt^2 + dx^2 + dy^2 + dz^2 \right) + \frac{R^2}{r^2 f} dR^2$$

$$\delta g_{\mu\nu} \sim e^{-i\omega t + iqz} h_{\mu\nu}(r) \quad \text{O(2) symmetry in x-y plane}$$

Shear channel:	h_{tx} h_{zx} h_{ty} h_{zy}		Z_1
Sound channel:	h_{tt} h_{tz} h_{zz} $h_{xx} + h_{yy}$		Z_2
Scalar channel:	h_{xy} $h_{xx} - h_{yy}$		Z_3

Other fluctuations (e.g. $\delta\varphi_1, \dots, \delta\varphi_n$) may affect sound channel

But not the shear channel  universality of η/s

Bulk viscosity and the speed of sound in $\mathcal{N} = 2^*$ SYM

$\mathcal{N} = 2^*$ is a “mass-deformed” $\mathcal{N} = 4$ (Pilch-Warner flow)

Ø Finite-temperature version: A.Buchel, J.Liu, hep-th/0305064

Ø The metric is known explicitly for $m/T \ll 1$

Ø Speed of sound and bulk viscosity:

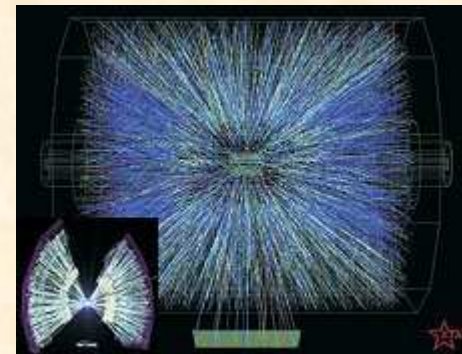
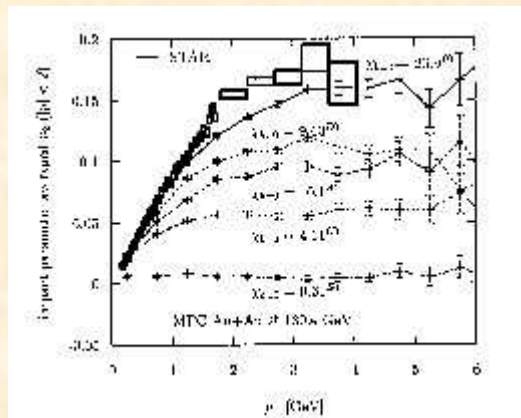
$$v_s = \frac{1}{\sqrt{3}} \left(1 - \frac{[\Gamma(\frac{3}{4})]^4}{3\pi^4} \left(\frac{m_f}{T}\right)^2 - \frac{1}{18\pi^4} \left(\frac{m_b}{T}\right)^4 + \dots \right)$$

$$\frac{\zeta}{\eta} = \beta_f^\Gamma \frac{[\Gamma(\frac{3}{4})]^4}{3\pi^3} \left(\frac{m_f}{T}\right)^2 + \frac{\beta_b^\Gamma}{432\pi^2} \left(\frac{m_b}{T}\right)^4 + \dots$$

$$\frac{\zeta}{\eta} = -\kappa \left(v_s^2 - \frac{1}{3} \right)$$

Relation to RHIC

- Ø IF quark-gluon plasma is indeed formed in heavy ion collisions
- Ø IF a hydrodynamic regime is unambiguously proven to exist
- Ø THEN hydrodynamic MODELS describe experimental results for e.g. elliptic flows well, provided $\eta/s \sim 1/4\pi$



- ▽ Bulk viscosity and speed of sound results are potentially interesting

Epilogue

- 1 AdS/CFT gives insights into physics of thermal gauge theories in the nonperturbative regime
- 1 Generic hydrodynamic predictions can be used to check validity of AdS/CFT
- 1 General algorithm exists to compute transport coefficients and the speed of sound in any gravity dual
- 1 Model-independent statements can presumably be checked experimentally