

Free fermions and BPS geometries

H. Lin, O. Lunin and J.M., hep-th/0409174

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Emergent space

- Matrix quantum mechanics \rightarrow free fermions (Brezin, Itzykson, Parisi, Zuber '78)



- Low energy $1+0 \rightarrow 1+1$ dimensions: ripples on the fermi sea
- Geometry of fermi sea \sim space-time geometry
- Large $N \sim$ classical limit

Outline

- Example \subset AdS/CFT as 1/2 BPS configurations
- Other theories with $R \times SO(4) \times SO(4)$ symmetry
- Solutions with $\widetilde{SU}(2|4)$ supergroup and associated theories

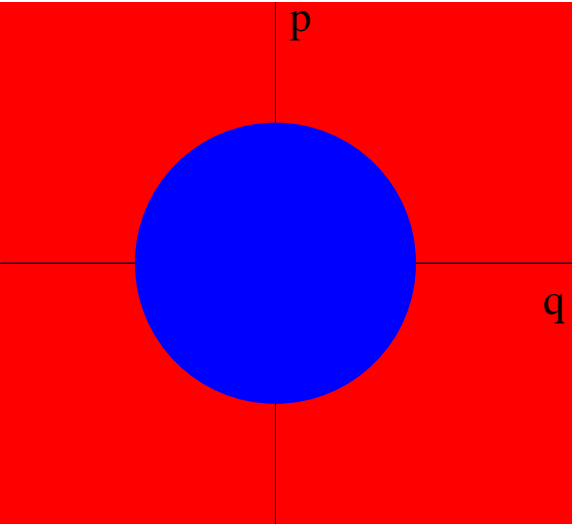
Free fermions from SYM

- Pick J in $SO(6)$, rotates

$$Z = \phi_1 + i\phi_2$$

- Select states with $E = J$ or $\Delta = J$.
- Only the zero mode of the field Z survives. Theory on $R \times S^3 \rightarrow$ harmonic oscillator for Z .

Corley, Jevicki, Ramgoolam; D. Berenstein



Phase space

Gravity Solutions

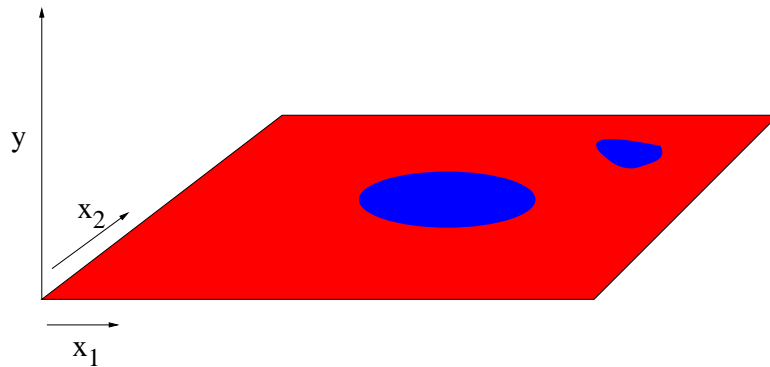
Symmetries:

$$(E - J) + SO(4)_{S^3} + SO(4)_{SO(6)} + 16 \text{ SUSY}$$

$$ds^2 = - * (dt + V)^2 + *(dy^2 + dx_1^2 + dx_2^2) + *d\Omega_3^2 + *d\Omega_3^2$$

The solution is specified by z

$$\partial_1^2 z + \partial_2^2 z + y \partial_y \left(\frac{\partial_y z}{y} \right) = 0$$

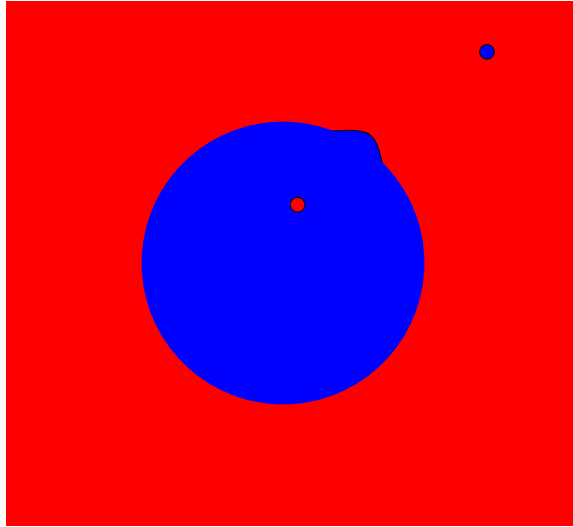


- $y \geq 0$
- 2 types of boundary conditions at $y = 0$

$$z = \frac{1}{2}, \quad \textcolor{red}{S^3} \rightarrow 0$$

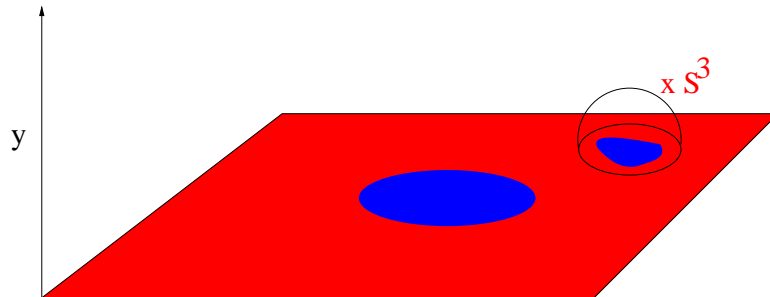
$$z = -\frac{1}{2}, \quad \textcolor{blue}{S^3} \rightarrow 0$$

- $y = 0$ plane \leftrightarrow Fermion phase space



- ripples \rightarrow gravitons
- single fermions \rightarrow giant gravitons $S^3 \subset AdS_5$
- single holes \rightarrow giant gravitons $S^3 \subset S^5$

D. Berenstein



- Geometric transitions S^3 that branes wrap
→ contractible
- Fermion/hole number → F_5 flux on 5-sphere
- Topology of solution → topology of droplets

Quantum Foam



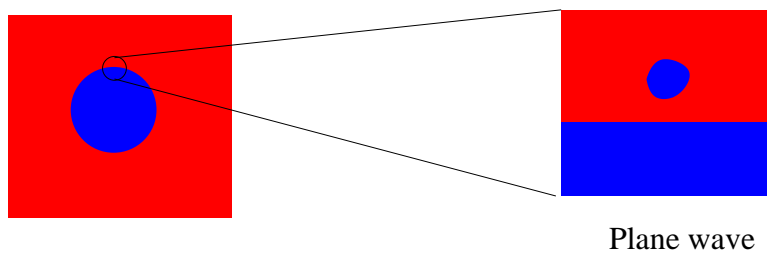
Small topological fluctuations \rightarrow already included as gravitons

$(\text{Boson} = \text{fermion}) \leftrightarrow (\text{Graviton} = \text{Foam})$

If density $0 < \rho < 1 \rightarrow$ null singularity

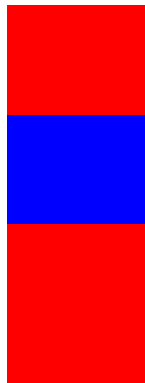
If density $\rho < 0$ or $1 < \rho \rightarrow$ closed time-like curves

Milanesi, O'Laughlin

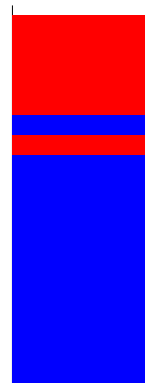


- $1/2$ BPS excitations on the IIB plane wave \rightarrow relativistic fermion.

Compact x_1



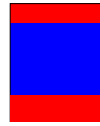
(a)



(b)

- (a) D4 brane on $R \times S^1 \times S^3$. SUSY vacua = states of 2d QCD on a cylinder.
- (b) M2 brane theory with a mass deformation. Pope, Warner ; Bena, Warner

Compact x_1, x_2



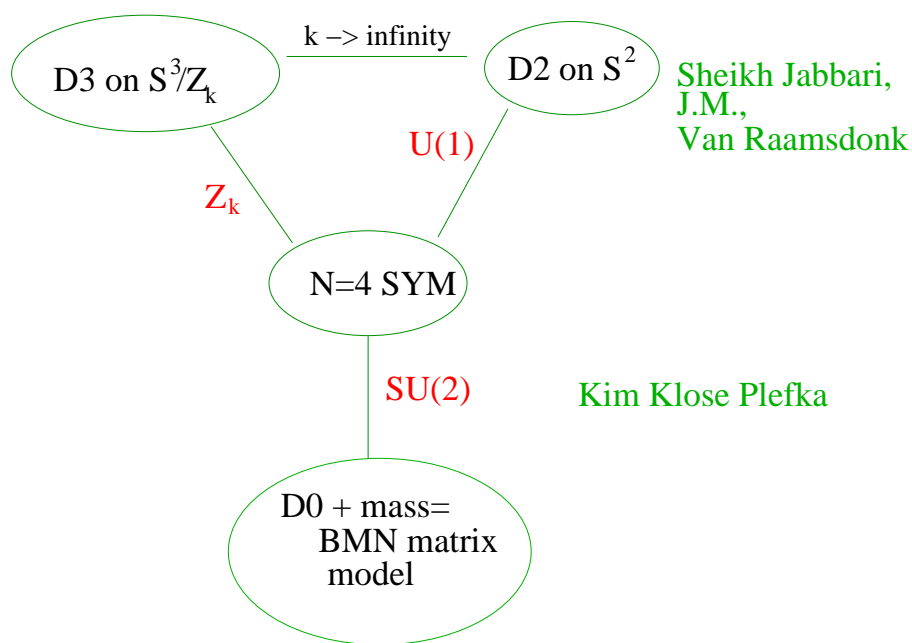
- Little string theory with $R \times T^2 \times S^3_K \times S^3_N$ boundary.
Itzhaki, Kutasov, Seiberg
- Low energies $SU(N)_K$ or $SU(K)_N$
Chern Simons theory on $R \times T^2$.
- Droplet configurations \rightarrow different vacua

Poincare susy algebra in $1+1$ or $2+1$

$$\{Q, Q\} = p_\mu + J_{SU(2)} + J_{SU(2)}$$

non-central charges. Possible in $d \leq 3$.

Theories with $\widetilde{SU}(2|4)$ symmetry



$$Z_k \subset U(1) \subset SU(2)_L \subset SO(4) \subset SO(2,4)$$

All have $SO(6)$ symmetry

Gravity solutions

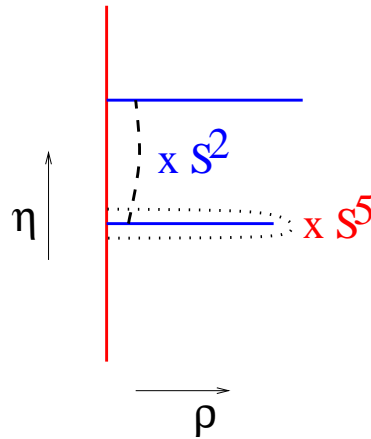
View the system in IIA.

$$ds^2 = - * (dt + \omega)^2 + *(d\rho^2 + d\eta^2) + * d\Omega_5^2 + * d\Omega_2^2 \quad (1)$$

Single function

$$\frac{1}{\rho} \partial_\rho (\rho \partial_\rho V) + \partial_\eta^2 V = 0$$

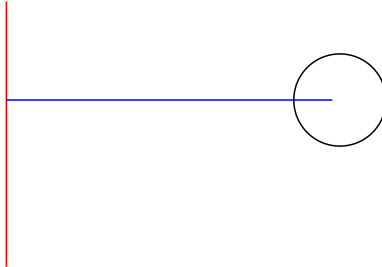
Electrostatic problem



- Charge on disks: **D2 charge**
- Separation: **NS-5 charge**

- Different configurations of disks \rightarrow different vacua
- Study 1/2 BPS states on these geometries with $E = J$, $J \subset SO(6)$
- We can study near BPS states (pp-wave limit)
- These states live near the tip of the disks

Near BPS states for D2 on S^2



Single disk \rightarrow IIA pp wave

$$ds^2 = -2dx^+dx^- - (\vec{r}^2 + 4\vec{y}^2)(dx^+)^2 + d\vec{r}^2 + d\vec{y}^2$$

Hyun, Shin

Bosons: masses 1 and 2. Universal result.

Near BPS:

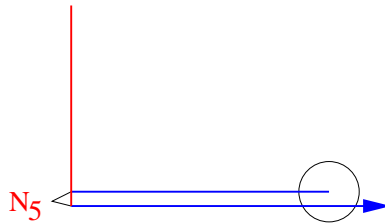
$$(E - J)_n = \sqrt{1 + (g_{YM}^2 N)^{\frac{2}{3}} \frac{n^2}{J^2}}$$

Coefficient is not universal.

Weak coupling:

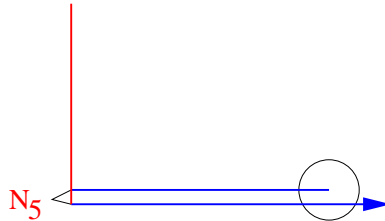
$$(E - J)_n = 1 + (g_{YM}^2 N) \frac{n^2}{J^2} + \dots$$

Solutions dual to the BMN matrix model



- Conducting plane at $\eta = 0$
- Vacua of BMN matrix model \rightarrow
dimension N $SU(2)$ representations
 \rightarrow
configurations of disks with fixed total dipole moment.

Near BPS states in the BMN matrix model



- IIA pp wave + N_5 NS-5-branes.
- Lighcone worldsheet \rightarrow (4,4) super-symmetric theory.
- (linear dilaton) \times $SU(2)_{N_5}$ WZW + potential

Near BPS, large N_5

$$(E - J)_n = \sqrt{1 + (g_{YM}^2 N_2)^{\frac{1}{2}} \frac{n^2}{J^2}}$$

$$N = N_2 N_5$$

Weak coupling $N_5 = 1$:

$$(E - J)_n = 1 + f(x) \frac{n^2}{J^2} + \dots$$

$$f(x) = x + c_2 x^2 + c_3 x^3 + c_4 x^4 + \dots$$

$$x = g_{YM}^2 N$$

Klose, Plefka ; Fischbacher, Klose, Plefka

We cannot do $N_5 = 1$

Conclusions

- Precise matching between field theories with large number of vacua, $e^{\sqrt{N}}$ and supergravity
- Unified gravity description for all theories with $\widetilde{SU}(2|4)$ supergroup.
- Interesting strong coupling results for many spin chains. Close relation to $\mathcal{N} = 4$ SYM.