

Toronto, November 11, 2004

Frobenius Manifolds and Integrable Hierarchies

Boris DUBROVIN

Lecture 3

Systems of Integrable PDEs

Bihamiltonian systems of n PDEs:

$$u_t^i = A_j^i(u)u_x^j + \epsilon \left(B_j^i(u)u_{xx}^j + \frac{1}{2}C_{jk}^i(u)u_x^j u_x^k \right) + O(\epsilon^2)$$

$$= \{u^i(x), H_1\}_1 = \{u^i(x), H_2\}_2$$

$i = 1, \dots, n.$

The rhs are differential ϵ -polynomials of degree 1,

$$\deg u^{(k)} = k, \quad k \geq 1, \quad \deg \epsilon = -1.$$

Hamiltonians are **local functionals**

$$H_k[u] = \int h_k(u; u_x, u_{xx}, \dots; \epsilon) dx, \quad k = 1, 2.$$

Compatible local Poisson brackets, i.e.

$$a_1\{ , \}_1 + a_2\{ , \}_2$$

is a Poisson bracket for any a_1, a_2

Locality

$$\{u^i(x), u^j(y)\}_{1,2} = \eta_{1,2}^{ij}(u(x))\delta'(x-y) + \Gamma_{k,1,2}^{ij}(u)u_x^k\delta(x-y) + O(\epsilon)$$

$$\det \eta_{1,2}^{ij}(u) \neq 0$$

Semisimplicity: characteristic roots of $(A_j^i(u))$ distinct for generic u

$u = (u^1, \dots, u^n) \in M$ local coordinates (later:
 M =Frobenius manifold)

Classify wrt the group of Miura-type transformations

$$u^i \mapsto \tilde{u}^i = f_0^i(u) + \epsilon f_1^i(u; u_x) + \epsilon^2 f_2^i(u; u_x, u_{xx}) + O(\epsilon^3)$$

$$\deg f_m^i(u; u_x, \dots, u^{(m)}) = m, \quad \det \left(\frac{\partial f_0^i(u)}{\partial u^j} \right) \neq 0$$

Remark One Poisson bracket

$$\{u^i(x), u^j(y)\} = \eta^{ij}(u(x))\delta'(x-y) + \Gamma_k^{ij}(u)u_x^k\delta(x-y) + O(\epsilon)$$

$$\det \eta^{ij}(u) \neq 0$$

is equivalent to

$$\{v^\alpha(x), v^\beta(y)\} = \eta^{\alpha\beta}\delta'(x - y), \quad \eta^{\alpha\beta} = \text{const}$$

([B.D., S.Novikov, 1983](#): $\eta^{ij}(u)$ is a **flat metric**;
[E.Getzler, 2001](#): triviality of Poisson cohomology)

Bihamiltonian \Rightarrow commuting Hamiltonians:
 find reducing transformation for the **pencil**

$$\{u^i(x), u^j(y)\}_2 - \lambda \{u^i(x), u^j(y)\}_1$$

$$u^i \mapsto v^\alpha = f_{-1}^\alpha(v; v_x, \dots; \epsilon) + \frac{f_0^\alpha(v; v_x, \dots; \epsilon)}{\lambda} + \frac{f_1^\alpha(v; v_x, \dots; \epsilon)}{\lambda^2} + \dots$$

$$\{v^\alpha(x), v^\beta(y)\} = -\lambda \eta^{\alpha\beta} \delta'(x - y)$$

Theorem The Hamiltonians

$$F_p^\alpha = \int f_p^\alpha(v; v_x, \dots; \epsilon), \quad \alpha = 1, \dots, n, \quad p \geq -1$$

commute

$$\{F_p^\alpha, F_q^\beta\}_{1,2} = 0$$

Proof The functionals

$$F^\alpha(\lambda) := \int v^\alpha dx = F_{-1}^\alpha + \frac{F_0^\alpha}{\lambda} + \frac{F_1^\alpha}{\lambda^2} + \dots$$

are Casimirs of the Poisson pencil

$$\{v^\alpha(x), F^\beta(\lambda)\}_2 - \{v^\alpha(x), F^\beta(\lambda)\}_1$$

$$= -\lambda \eta^{\alpha\beta} \int \delta'(x - u) dy = 0$$

Hierarchy = commuting bihamiltonian flows on the space of vector functions $u(x) = (u^1(x), \dots, u^n(x))$ of the **spatial** variable x organized in n infinite chains

$$t = (t^{\alpha,p}), \quad \alpha = 1, \dots, n, \quad p = 0, 1, 2, \dots$$

$$\frac{\partial u}{\partial t^{\alpha,p}} = \{u(x), H_{\alpha,p}\}_1$$

$$H_{\alpha,p} = \int h_{\alpha,p}(u; u_x, \dots; \epsilon) dx$$

$$\left[\frac{\partial}{\partial t^{\alpha,p}}, \frac{\partial}{\partial t^{\beta,q}} \right] = 0$$

$$\frac{\partial}{\partial t^{1,0}} = \frac{\partial}{\partial x}$$

satisfying certain **additional properties**

- **triangular** bihamiltonian recursion

$$\{ \cdot , H_{\alpha,p-1} \}_2 = \sum_{q \leq p} R_{\alpha,p}^{\beta,q} \{ \cdot , H_{\beta,q} \}_1$$

starting from the **Casimirs** of $\{ \cdot , \cdot \}_1$

$$H_{\alpha,-1} = \int u_\alpha dx, \quad \{ \cdot , H_{\alpha,-1} \}_1 \equiv 0$$

(recursion for F_p^α : $\{ \cdot , F_{p-1}^\alpha \}_2 = \{ \cdot , F_p^\alpha \}_1$)

- existence of **tau-function**.

Hamilton densities

$$h_{\alpha,p}(u; u_x, \dots; \epsilon) = \epsilon^2 \frac{\partial^2 \log \tau}{\partial x \partial t^{\alpha,p+1}}$$

\Rightarrow **strange symmetry**

$$\frac{\partial h_{\alpha,p-1}}{\partial t^{\beta,q}} = \frac{\partial h_{\beta,q-1}}{\partial t^{\alpha,p}} = \partial_x \Omega_{\alpha,p;\beta,q}(u; u_x, \dots; \epsilon)$$

Tau-structure: choice of a symmetric matrix
 $\Omega_{\alpha,p;\beta,q}(u; u_x, \dots; \epsilon) = \Omega_{\beta,q;\alpha,p}(u; u_x, \dots; \epsilon)$
satisfying the above properties

In particular

$$u_\alpha = \epsilon^2 \frac{\partial^2 \log \tau}{\partial x \partial t^{\alpha,0}}, \quad \alpha = 1, \dots, n.$$

(recall: $x = t^{1,0}$, $h_{\alpha,-1} = u_\alpha$)

Tau-function $\tau(t; \epsilon)$ depends on the choice of a solution

$$u(x, t; \epsilon) = u_0(t) + \epsilon u_1(x, t) + \epsilon^2 u_2(x, t) + \dots$$

$$\epsilon^2 \frac{\partial^2 \log \tau}{\partial t^{\alpha,p} \partial t^{\beta,q}} = \Omega_{\alpha,p;\beta,q}(u; u_x, \dots; \epsilon)$$

Results (B.D., Y.Zhang)

Theorem 1 At $\epsilon = 0$

hierarchies \leftrightarrow semisimple Frobenius manifolds M^n (or their degenerations)

Must also allow the operation of **changing the spatial direction**

$$\frac{\partial}{\partial x} = \frac{\partial}{\partial t^{1,0}} \mapsto \frac{\partial}{\partial \tilde{x}} = \sum b^i \frac{\partial}{\partial t^{i,0}}$$

$\Rightarrow \frac{n(n-1)}{2}$ parameters* of dispersionless ($=$ no ϵ) integrable hierarchies with n dependent variables.

moduli = holonomy of the deformed flat connection $\tilde{\nabla}$ on $M^n \times \mathbb{C}^$

Construction of the **Principal Hierarchy** associated with M (B.D. 1992)

$$\frac{\partial v}{\partial t^{\alpha,p}} = \nabla \theta_{\alpha,p}(v) \cdot v_x = \{v(x), H_{\alpha,p}\}_1$$

$$H_{\alpha,p} = \int \theta_{\alpha,p+1}(v) dx$$

Here the functions

$$f_\alpha(v, z) = \sum_{p \geq 0} \theta_{\alpha,p}(v) z^{p+\mu_\alpha}, \quad \alpha = 1, \dots, n$$

are such that

$$\tilde{\nabla} df_\alpha = 0$$

$\mu_1, \mu_2, \dots, \mu_n$ = eigenvalues of $\frac{d-2}{2} - \nabla E$ (here assuming semisimplicity and **nonresonancy** of ∇E , $\mu_\alpha - \mu_\beta \notin \mathbb{Z}_{>0}$).

(First) Hamiltonian structure

$$\{v^\alpha(x), v^\beta(y)\}_1 = \eta^{\alpha\beta} \delta'(x - y)$$

i.e.

$$\frac{\partial v^\beta}{\partial t^{\alpha,p}} = \partial_x \eta^{\beta\gamma} \frac{\delta H_{\alpha,p}}{\delta v^\gamma(x)} = \partial_x \eta^{\beta\gamma} \frac{\partial \theta_{\alpha,p+1}}{\partial v^\gamma}$$

Proof uses the *horizontality equations*

$$\tilde{\nabla} df_\alpha = 0, \text{ i.e.}$$

$$\partial_\lambda \partial_\mu \theta_{\alpha,p+1} = c_{\lambda\mu}^\nu(v) \partial_\nu \theta_{\alpha,p}$$

Second Hamiltonian structure

$$\{u^\alpha(x), u^\beta(y)\}_2 = g^{\alpha\beta}(v(x)) \delta'(x-y) + \Gamma_\gamma^{\alpha\beta}(v) v_x^\gamma \delta(x-y)$$

with

$$g^{\alpha\beta}(v) := E^\gamma(v) c_\gamma^{\alpha\beta}(v)$$

another flat (outside *discriminant* $\det g = 0$) **metric** on M

Example Dispersionless KdV

$$u_{t_p} = \frac{u^p}{p!} u_x$$

Here $n = 1$

$$\Omega_{pq} = \frac{u^{p+q+1}}{p! q! (p+q+1)}$$

$$\frac{\partial \Omega_{pq}}{\partial t^r} = \frac{u^p}{p!} \frac{u^q}{q!} \frac{u^r}{r!} u_x$$

$$\{u(x), u(y)\}_1 = \delta'(x - y)$$

$$\{u(x), u(y)\}_2 = u(x) \delta'(x - y) + \frac{1}{2} u_x \delta(x - y)$$

$$H_p = \int \frac{u^{p+2}}{(p+2)!} dx$$

$$\{ \cdot, H_{p-1} \}_2 = \left(p + \frac{1}{2} \right) \{ \cdot, H_p \}_1$$

Solutions $u = u(\mathbf{t})$ given in implicit form

$$\sum_{p \geq 0} (t_p - c_p) \frac{u^p}{p!} = 0$$

arbitrary constants c_p

Tau-function

$$\log \tau = \frac{1}{2} \sum \Omega_{pq}(u(\mathbf{t}))(t_p - c_p)(t_q - c_q)$$

Theorem 1'. The dispersionless limit of any tau-structure \leftrightarrow trivialization of the **deformed flat connection** on $M \times \mathbb{C}^*$

$$\tilde{\nabla}_a b = \nabla_a b + z a \cdot b, \quad a, b \in T_x M$$

$$\begin{aligned}\tilde{\nabla}_{\frac{d}{dz}} b &= \partial_z b + E \cdot b + \frac{\mu}{z} b \\ \mu &= \frac{d-2}{2} - \nabla E\end{aligned}$$

(= *calibrated* Frobenius manifold in the Givental's terminology)

Conversely: from tau-structures to Frobenius manifolds.

$$\Omega_{\alpha,0;\beta,0}|_{\epsilon=0} = \frac{\partial^2 F(u)}{\partial u^\alpha \partial u^\beta}$$

The function $F(u)$ satisfies $WDVV \Rightarrow$ a **structure of semisimple Frobenius manifold** on the space of fields (u^1, \dots, u^n)

(must also allow the operation of change of the spatial direction

$$\frac{\partial}{\partial t^{1,0}} = \frac{\partial}{\partial x} \mapsto \sum b^\alpha \frac{\partial}{\partial t^{\alpha,0}} = \frac{\partial}{\partial x}$$

for constant b^α)

E.g., for the (extended) Toda the change

$$\frac{\partial}{\partial s_0} \leftrightarrow \frac{\partial}{\partial t_0}$$

corresponds to

$$\text{Toda} \leftrightarrow \text{NLS}$$

Next step: classification of tau-structures with the **given** dispersionless limit (i.e., associated with a given calibrated semisimple Frobenius manifold)

Deformation theory of bihamiltonian structures:

two differentials ∂_1 and ∂_2 on multivectors (Schouten - Nijenhuis brackets with $\{ , \}_1$ and $\{ , \}_2$ resp.)

$$\partial_1^2 = \partial_2^2 = \partial_1 \partial_2 + \partial_2 \partial_1 = 0$$

In our case both ∂_1 and ∂_2 are acyclic
 \Rightarrow the space of infinitesimal deformations $H^2(\partial_1, \partial_2)$ equals

$$H^2(\partial_1, \partial_2) = \text{Ker } \partial_1 \partial_2 / \text{Im } \partial_1 + \text{Im } \partial_2$$

(Here $\partial_1 \partial_2$ acts on vector fields).

Theorem 2 The deformation space of a given Principal Hierarchy with n dependent variables is at most n -dimensional

Construction of the parameters c_1, \dots, c_n for a bihamiltonian structure

$$\{w^\alpha(x), w^\beta(y)\}_1 = \eta^{\alpha\beta}\delta'(x-y) + \epsilon^2 A_1^{\alpha\beta}(w(x))\delta'''(x-y) + \dots$$

$$\{w^\alpha(x), w^\beta(y)\}_2 = g^{\alpha\beta}(w(x))\delta'(x-y) + \Gamma_\gamma^{\alpha\beta}(w)w_x^\gamma\delta(x-y)$$

$$+ \epsilon^2 A_2^{\alpha\beta}(w(x))\delta'''(x-y) + \dots$$

- Canonical coordinates u_1, \dots, u_n on M : roots of

$$\det(g^{\alpha\beta}(v) - \lambda \eta^{\alpha\beta}) = 0$$

Main property:

$$\partial/\partial u_i \cdot \partial/\partial u_j = \delta_{ij} \partial/\partial u_i$$

The two metrics are diagonal in the canonical coordinates

$$\eta = \sum_{i=1}^n \eta_{ii}(u) du_i^2, \quad g = \sum_{i=1}^n \eta_{ii}(u) \frac{du_i^2}{u_i}.$$

- Define

$$B_k^{ij}(u) := \frac{\partial u^i}{\partial w^\alpha} \frac{\partial u^j}{\partial w^\beta} A_k^{\alpha\beta}(w), \quad k = 1, 2$$

- Then

$$c_i := \frac{1}{3} \eta_{ii}^2(u) (B_2^{ii}(u) - u_i B_1^{ii}(u)), \quad i = 1, \dots, n.$$

This is the **complete** set of invariants of the bihamiltonian structure with the **given leading term** at $\epsilon = 0$.

Example For $n = 1$ all tau-structures are equivalent to KdV

$$u_t = u u_x + 2c \epsilon^2 u_{xxx}$$

Proof uses results of [S.-Q.Liu, Y.Zhang](#) on deformations of bihamiltonian structures and also **Quasitriviality Theorem**

Quasitiviality Theorem: any tau-structure can be obtained from the dispersionless limit by a transformation

$$v_\alpha \mapsto w_\alpha = v_\alpha + \sum_{k \geq 1} \epsilon^k F_\alpha^{[k]}(v; v_x, v_{xx}, \dots)$$

with some functions $F_\alpha^{[k]}$ **rational** in the derivatives (B.D., S.-Q.Liu, Y.Zhang, math/0410027)

Example 1

Riemann wave \mapsto KdV

$$v_t + v v_x = 0 \quad \mapsto \quad w_t + w w_x + \frac{\epsilon^2}{12} w_{xxx} = 0$$

The substitution

$$\begin{aligned} w &= v + \frac{\epsilon^2}{24} \partial_x^2 (\log v_x) \\ &+ \epsilon^4 \partial_x^2 \left(\frac{v^{IV}}{1152 v_x^2} - \frac{7 v_{xx} v_{xxx}}{1920 v_x^3} + \frac{v_{xx}^3}{360 v_x^4} \right) + O(\epsilon^6). \end{aligned}$$

Baikov, Gazizov, Ibragimov, 1989

Example 2

Riemann wave \mapsto Camassa-Holm

$$v_t = \frac{3}{2} v v_x \mapsto w_t = (1 - \epsilon^2 \partial_x^2)^{-1} \left(\frac{3}{2} w w_x - \epsilon^2 \left[w_x w_{xx} + \frac{1}{2} w w_{xxx} \right] \right)$$

$$\begin{aligned} w = & v + \epsilon^2 \partial_x \left(\frac{v v_{xx}}{3 v_x} - \frac{v_x}{6} \right) \\ & + \epsilon^4 \partial_x \left(\frac{7 v_{xx}^2}{45 v_x} + \frac{45 v v_{xx}^3}{16 v_x^3} - \frac{45 v^2 v_{xx}^4}{32 v_x^5} - \frac{v_{xxx}}{8} - \frac{59 v v_{xx} v_{xxx}}{90 v_x^2} \right. \\ & \quad \left. + \frac{37 v^2 v_{xx}^2 v_{xxx}}{30 v_x^4} - \frac{7 v^2 v_{xxx}^2}{30 v_x^3} + \frac{5 v v^{IV}}{18 v_x} \right. \\ & \quad \left. - \frac{31 v^2 v_{xx} v^{IV}}{90 v_x^3} + \frac{v^2 v^V}{18 v_x^2} \right) + O(\epsilon^6) \end{aligned}$$

Lorenzoni, 2002

From Frobenius Manifolds to Integrable Systems of the Topological Type

Hierarchy of the **topological type**:

$$c_1 = c_2 = \dots = c_n = \frac{1}{24}$$

Main goal: for any n construct **universal** integrable hierarchy of the topological type