

A Heterotic Standard Model

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Burt's talk next week

hep-th/0410055: Elliptic Calabi-Yau Threefolds with $\mathbb{Z}_3 \times \mathbb{Z}_3$ Wilson Lines

hep-th/0501070: A Heterotic Standard Model

hep-th/0502155: A Standard Model from the $E_8 \times E_8$ Heterotic Superstring

hep-th/0504xxx: Standard-Model Vacua in Heterotic String Theory

Table of Contents

- Introduction
- Group Theory
- Calabi-Yau Manifolds
- Vector Bundles
- Low Energy Spectrum
- Conclusions

Introduction

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- $d = 4, \mathcal{N} = 1 \Rightarrow$ stable background.
- ~~$SU(3)_C \times SU(2)_L \times U(1)_Y$~~ .
- $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_{B-L} \Rightarrow$ proton decay suppressed.

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- $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_{B-L} \Rightarrow$ proton decay suppressed.
- No exotic matter.
- All of the ordinary matter fields (including right-handed Neutrino).

Table of Contents

- Introduction
- Group Theory
 - An Organizational Principle
 - Wilson Line Breaking
 - More Group Theory
 - Wish List
- Calabi-Yau Manifolds
- Vector Bundles
- Low Energy Spectrum
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An Organizational Principle

Ancient Lore: $\text{Spin}(10)$ GUT with $\mathbb{Z}_3 \times \mathbb{Z}_3$ Wilson lines “works”:

$\underline{16}$ of $\text{Spin}(10)$: Breaks into one family of quarks and leptons including a right-handed Neutrino.

$\overline{16}$ of $\text{Spin}(10)$: Anti-family.

$\underline{10} = \overline{\underline{10}}$ of $\text{Spin}(10)$: Higgs and color triplets.

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$\underline{10} = \overline{10}$ of $\text{Spin}(10)$: Higgs and color triplets.

However, we do not care about GUTs:

Compactification scale \sim GUT scale

... but nice way to package representations.

Wilson Line Breaking

$$\text{Spin}(10) \supset SU(3) \times SU(2) \times U(1) \times U(1) \times \mathbb{Z}_3 \times \mathbb{Z}_3$$

$$\left\{ \begin{array}{l} \text{Standard Model} \\ \text{gauge group} \end{array} \right\} \times U(1)_{B-L} \times \{\text{Wilson lines}\}$$

$\mathbb{Z}_3 \times \mathbb{Z}_3$ is smallest Wilson line possible.

Wilson Line Breaking

$$\text{Spin}(10) \supset SU(3) \times SU(2) \times U(1) \times U(1) \times \mathbb{Z}_3 \times \mathbb{Z}_3$$

$$\begin{aligned} \underline{16} &= \chi_1^2 \chi_2(\underline{3}, \underline{2}, 1, 1) \oplus \chi_1^2(\underline{1}, \underline{1}, 6, 3) \oplus \\ &\quad \oplus \chi_1^2 \chi_2^2(\overline{\underline{3}}, \underline{1}, -4, -1) \oplus \chi_2^2(\overline{\underline{3}}, \underline{1}, 2, -1) \oplus \\ &\quad \oplus (\underline{1}, \overline{\underline{2}}, -3, -3) \oplus \chi_1(\underline{1}, \underline{1}, 0, 3) \\ \underline{10} &= \chi_1(\underline{1}, \underline{2}, 3, 0) \oplus \chi_1 \chi_2(\underline{3}, \underline{1}, -2, -2) \oplus \\ &\quad \oplus \chi_1^2(\underline{1}, \overline{\underline{2}}, -3, 0) \oplus \chi_1^2 \chi_2^2(\overline{\underline{3}}, \underline{1}, 2, 2) \end{aligned}$$

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Right-handed Neutrino

More Group Theory

$$G = \mathbb{Z}_3 \times \mathbb{Z}_3 = G_1 \times G_2$$

Fix generators g_1 and g_2 .

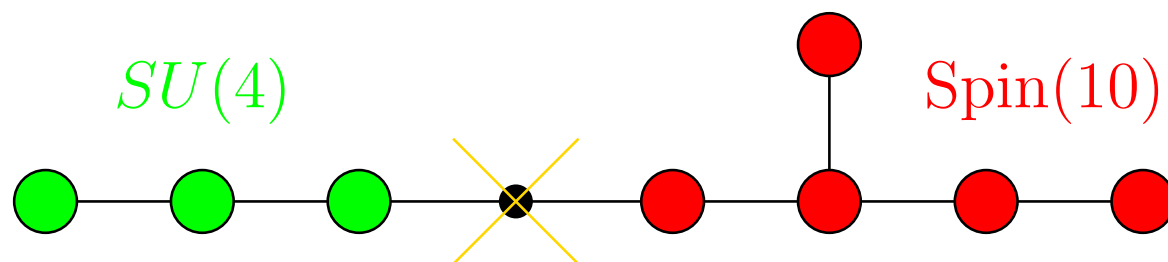
Characters (=1-d representations): Denote generators by χ_1 and χ_2 , where ($\omega = e^{\frac{2\pi i}{3}}$)

$$\begin{aligned}\chi_1(g_1) &= \omega & \chi_1(g_2) &= 1 \\ \chi_2(g_1) &= 1 & \chi_2(g_2) &= \omega .\end{aligned}$$

All other characters are products of χ_1 and χ_2 .

Yet More Group Theory

Maximal regular subgroup $SU(4) \times \text{Spin}(10) \subset E_8$:



The adjoint of E_8 (fermions in the $E_8 \times E_8$ heterotic string) decomposes as

$$\underline{248} = (\underline{1}, \underline{45}) \oplus (\underline{15}, \underline{1}) \oplus (\underline{4}, \underline{16}) \oplus (\overline{\underline{4}}, \overline{\underline{16}}) \oplus (\underline{6}, \underline{10})$$

Wish List

To make use of this group theory, we would like

- A Calabi-Yau threefold X with $\mathbb{Z}_3 \times \mathbb{Z}_3$ fundamental group.

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- A $SU(4) \subset E_8$ instanton leaves $\text{Spin}(10)$ unbroken, so we want a rank 4 stable holomorphic vector bundle \mathcal{V} on X .

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- A Calabi-Yau threefold X with $\mathbb{Z}_3 \times \mathbb{Z}_3$ fundamental group.
- The Calabi-Yau should be torus fibered.
- A $SU(4) \subset E_8$ instanton leaves $\text{Spin}(10)$ unbroken, so we want a rank 4 stable holomorphic vector bundle \mathcal{V} on X .
- With the “right” cohomology groups (low energy spectrum).



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- Introduction
- Group Theory
- Calabi-Yau Manifolds
 - Construction
 - Properties
 - Group Actions on del Pezzo Surfaces
 - Hodge Numbers
 - Divisors
- Vector Bundles
- Low Energy Spectrum
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CY Introduction

Calabi-Yau threefold
 X with
 $\pi_1(X) = \mathbb{Z}_3 \times \mathbb{Z}_3$

CY Introduction

Work with

Simply connected
Calabi-Yau threefold
 \tilde{X} with free $\mathbb{Z}_3 \times \mathbb{Z}_3$
action

=

Have in mind

Calabi-Yau threefold
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elliptically fibered

(torus fibered
with section)

Have in mind

Calabi-Yau threefold
 X with
 $\pi_1(X) = \mathbb{Z}_3 \times \mathbb{Z}_3$

torus fibered

(assuming $\mathbb{Z}_3 \times \mathbb{Z}_3$
preserves fibration)

=

Calabi-Yau Construction

Start with two dP_9 surfaces B_1 and B_2 .

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Note: dP_9 are elliptically fibered; Fibers over a generic point $x \in \mathbb{P}^1$ are

$$\beta_1^{-1}(x) \simeq T^2 \subset B_1, \quad \beta_2^{-1}(x) \simeq T^2 \subset B_2.$$

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Calabi-Yau Properties

- $\tilde{X} \stackrel{\text{def}}{=} B_1 \times_{\mathbb{P}^1} B_2$ is a simply connected Calabi-Yau threefold, $c_1(\tilde{X}) = 0$.

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- $h^{1,1}(\tilde{X}) = 19 = h^{2,1}(\tilde{X})$

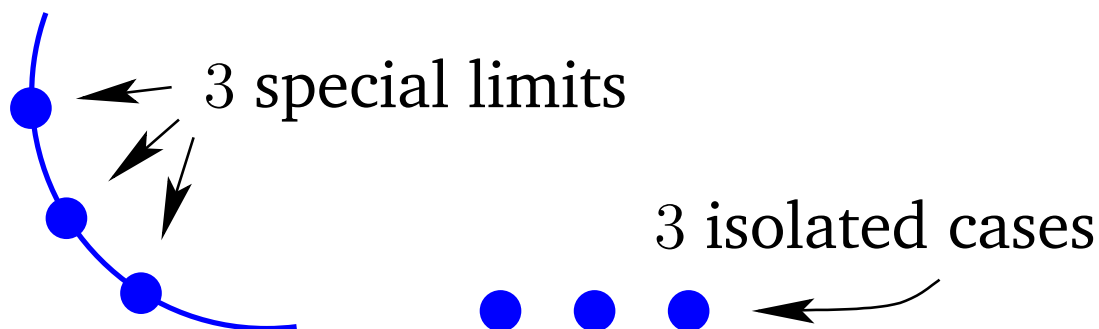
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- $h^{1,1}(\tilde{X}) = 19 = h^{2,1}(\tilde{X})$
- Group actions on B_1, B_2 lift to \tilde{X} if their action on the common base \mathbb{P}^1 is identical.

Group Actions on the del Pezzo I

We classified all $\mathbb{Z}_3 \times \mathbb{Z}_3$ actions on dP_9 surfaces.
The moduli space looks like this:

A one parameter family



Group Actions on the del Pezzo II

All such dP_9 surfaces with $G = \mathbb{Z}_3 \times \mathbb{Z}_3$ action give rise to a G action on \tilde{X} .

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Group Actions on the del Pezzo II

All such dP_9 surfaces with $G = \mathbb{Z}_3 \times \mathbb{Z}_3$ action give rise to a G action on \tilde{X} .

- The 3 isolated cases never yield a free $\mathbb{Z}_3 \times \mathbb{Z}_3$ action.
- The one-parameter family and its limits can give a free $\mathbb{Z}_3 \times \mathbb{Z}_3$ action on \tilde{X} .

We only consider this one-parameter family in the following.

Group Action on Cohomology

$G = \mathbb{Z}_3 \times \mathbb{Z}_3$ acts on the cohomology groups of B_1 , B_2 , and \tilde{X} .

On $H^2(B_i, \mathbb{Z}) = \mathbb{Z}^{10}$ the group generators g_1, g_2 act as

$$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 3 & -1 & -1 & -1 \\ 0 & 1 & 0 & 1 & 0 & 0 & 3 & -1 & -1 & -1 \\ 0 & 0 & 0 & 0 & 1 & 0 & -2 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -2 & 0 & 2 & 1 \\ 0 & 0 & 0 & -1 & 0 & 1 & -2 & 1 & 1 & 1 \\ 0 & 0 & 1 & 0 & 0 & 0 & -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & -1 & 1 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3 & 0 & 3 & 1 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 3 & 0 & 0 & 0 & 0 & -1 & 2 & -1 \\ 0 & 1 & 3 & 0 & 0 & 1 & 1 & -1 & 2 & -1 \\ 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 & -1 & 1 \\ 0 & 0 & -2 & 0 & 0 & -1 & 0 & 1 & -2 & 1 \\ 0 & 0 & -2 & 0 & 0 & 0 & -1 & 1 & -1 & 1 \\ 0 & 0 & -1 & 1 & 0 & -1 & 0 & 0 & -1 & 0 \\ 0 & 0 & -1 & 0 & 1 & 0 & -1 & 1 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 1 & 0 & -3 & 0 & 0 & 0 & 0 & 1 & -2 & 2 \end{pmatrix}$$

Similarly, 19×19 matrices for g_1, g_2 on $H^{1,1}(\tilde{X}) = \mathbb{C}^{19}$.

Invariant Cohomology

$$G = \mathbb{Z}_3 \times \mathbb{Z}_3 \text{ action free} \Rightarrow H^{p,q}(X) = H^{p,q}(\tilde{X})^G$$

$$\Rightarrow \text{Hodge diamond } h^{p,q}(X) = \begin{array}{ccccc} & & & 1 & \\ & & & 0 & 0 \\ & & 0 & 3 & 0 \\ & 1 & 3 & 3 & 1 \\ & 0 & 3 & 0 & \\ & & 0 & 0 & \\ & & & 1 & \end{array}$$

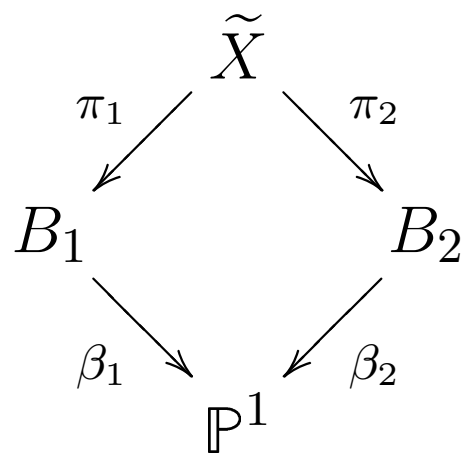
$h^{1,1}(X) = 3$ dimensional space of divisor classes.

Divisors on the Calabi-Yau I

$\dim_{\mathbb{C}} = 3 :$

$\dim_{\mathbb{C}} = 2 :$

$\dim_{\mathbb{C}} = 1 :$

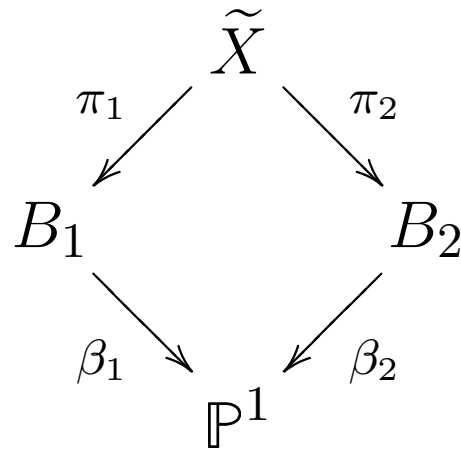


Divisors on the Calabi-Yau I

$\dim_{\mathbb{C}} = 3 :$

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$\dim_{\mathbb{C}} = 1 :$



$$H^{1,1}(B_1)^G = \mathbb{C}f_1 \oplus \mathbb{C}t_1$$

$$H^{1,1}(B_2)^G = \mathbb{C}f_2 \oplus \mathbb{C}t_2$$

Divisors on the Calabi-Yau I

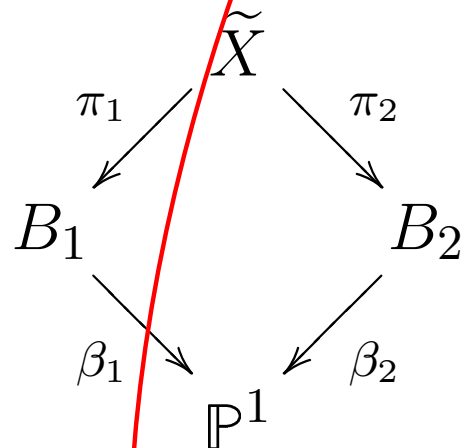


$$H^{1,1}(\tilde{X})^G = \mathbb{C}\pi_1^{-1}(f_1) + \mathbb{C}\pi_1^{-1}(t_1) + \mathbb{C}\pi_2^{-1}(f_2) + \mathbb{C}\pi_2^{-1}(t_2)$$

$\dim_{\mathbb{C}} = 3 :$

$\dim_{\mathbb{C}} = 2 :$

$\dim_{\mathbb{C}} = 1 :$



$$H^{1,1}(B_1)^G = \mathbb{C}f_1 \oplus \mathbb{C}t_1$$

$$H^{1,1}(B_2)^G = \mathbb{C}f_2 \oplus \mathbb{C}t_2$$



Divisors on the Calabi-Yau II



$$\begin{aligned} H^{1,1}(\tilde{X})^G &= \mathbb{C}\pi_1^{-1}(f_1) + \mathbb{C}\pi_1^{-1}(t_1) + \mathbb{C}\pi_2^{-1}(f_2) + \mathbb{C}\pi_2^{-1}(t_2) \\ &= \mathbb{C}\phi \oplus \mathbb{C}\tau_1 \oplus \mathbb{C}\tau_2 \end{aligned}$$

There is one relation:

$$\begin{aligned} \pi_1^{-1}(f_1) &= \left\{ T^4 \text{ fiber of } \begin{array}{c} \tilde{X} \\ \downarrow \\ \mathbb{P}^1 \end{array} \right\} = \pi_2^{-1}(f_2) \stackrel{\text{def}}{=} \phi \\ \pi_1^{-1}(t_1) &\stackrel{\text{def}}{=} \tau_1 \qquad \pi_2^{-1}(t_2) \stackrel{\text{def}}{=} \tau_2 \end{aligned}$$



Table of Contents



- Introduction
- Group Theory
- Calabi-Yau Manifolds
- Vector Bundles
 - Line Bundles
 - Equivariant Line Bundles
 - The Serre Construction
 - Equivariant Vector Bundles
 - Constructing Vector Bundles
 - Our Solution
- Low Energy Spectrum
- Conclusions



Line Bundles



On any variety Y , we have

$$\left\{ \text{Divisors } D \right\} / \sim = \left\{ \text{Line bundles } \mathcal{O}_Y(D) \right\}$$



Line Bundles

On any variety Y , we have

$$\{\text{Divisors } D\} / \sim = \{\text{Line bundles } \mathcal{O}_Y(D)\}$$

Linear equivalence

For B_1, B_2, \tilde{X} that is just cohomology class of the divisor in $H^{1,1}$.

Every line bundle is of the form

- $\mathcal{O}_{\tilde{X}}(x_1\tau_1 + x_2\tau_2 + x_3\phi)$, $x_1, x_2, x_3 \in \mathbb{Z}$.
- $\mathcal{O}_{B_i}(y_1t_i + y_2f_i)$, $y_1, y_2 \in \mathbb{Z}$.

Equivariant Line Bundles I

Line bundles on
 $X = \tilde{X}/G$

Equivariant Line Bundles I

Work with

G -equivariant line
bundles on \tilde{X}

=

Have in mind

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An equivariant line bundle is a line bundle \mathcal{L} together with a group action $\gamma : G \times \mathcal{L} \rightarrow \mathcal{L}$:

$$\begin{array}{ccc} \mathcal{L} & \xrightarrow{\gamma_g} & \mathcal{L} \\ \downarrow & & \downarrow \\ \tilde{X} & \xrightarrow{g} & \tilde{X} \end{array}$$

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- Most line bundles on \tilde{X} cannot be made equivariant.

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 $x_1, x_2, x_3 \in \mathbb{Z}$ with $x_1 + x_2 \equiv 0 \pmod{3}$ allow for a
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 $x_1, x_2, x_3 \in \mathbb{Z}$ with $x_1 + x_2 \equiv 0 \pmod{3}$ allow for a
 $G = \mathbb{Z}_3 \times \mathbb{Z}_3$ action.
- In these cases, there is always more than one G
action
 \Rightarrow Different equivariant line bundles!

Equivariant Structures

Consider the trivial line bundle $\mathcal{O}_{\tilde{X}} = \tilde{X} \times \mathbb{C}$.

- Obvious equivariant action

$$\gamma_g : \tilde{X} \times \mathbb{C} \rightarrow \tilde{X} \times \mathbb{C}, \quad (p, v) \mapsto (g(p), v)$$

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- Different equivariant action by multiplying with a character

$$\chi\gamma_g : \tilde{X} \times \mathbb{C} \rightarrow \tilde{X} \times \mathbb{C}, (p, v) \mapsto (g(p), \chi(g)v)$$

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- We write $\chi\mathcal{O}_{\tilde{X}}$ for this different equivariant line bundle.

The Serre Construction

A way to construct many stable rank 2 vector bundles on a surface (here: B_1 and B_2).

- Take two line bundles $\mathcal{L}_1, \mathcal{L}_2$.
- An ideal sheaf I (sheaf of functions vanishing at some fixed points).
- Define \mathcal{S} as an extension

$$0 \longrightarrow \mathcal{L}_1 \longrightarrow \mathcal{S} \longrightarrow \mathcal{L}_2 \otimes I \longrightarrow 0$$

- Cayley-Bacharach property \Rightarrow generic extension is a rank 2 vector bundle.

Serre Construction Example

$$0 \longrightarrow \mathcal{O}_{B_2}(-2f_2) \longrightarrow \mathcal{W} \longrightarrow \mathcal{O}_{B_2}(2f_2) \otimes I_9 \longrightarrow 0$$

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- I_9 is the ideal sheaf of (certain) nine points.
- Has the Cayley-Bacharach property.
- $\text{Ext}^1 \left(\mathcal{O}_{B_2}(2f_2) \otimes I_9, \mathcal{O}_{B_2}(-2f_2) \right) = \mathbb{C}^9$,
so there exist extensions.

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Vector bundles on
 $X = \tilde{X}/G$

Problem: Even if \mathcal{E}, \mathcal{F} are equivariant,

$$0 \longrightarrow \mathcal{E} \longrightarrow \mathcal{V} \longrightarrow \mathcal{F} \longrightarrow 0$$

Extension is not necessarily equivariant!

Only extensions in $\mathrm{Ext}^1(\mathcal{F}, \mathcal{E})^G$ are equivariant.

Equivariant Example

$$0 \longrightarrow \mathcal{O}_{B_2}(-2f_2) \longrightarrow \mathcal{W} \longrightarrow \chi_2 \mathcal{O}_{B_2}(2f_2) \otimes I_9 \longrightarrow 0$$

- $\mathcal{O}_{B_2}(-2f_2), \chi_2 \mathcal{O}_{B_2}(2f_2)$ are equivariant.

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 $1 \oplus \chi_1 \oplus \chi_1^2 \oplus \chi_2 \oplus \chi_1 \chi_2 \oplus \chi_1^2 \chi_2 \oplus \chi_2^2 \oplus \chi_1 \chi_2^2 \oplus \chi_1^2 \chi_2^2$

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so there exist extensions.

Constructing Vector Bundles



Building blocks:

- Line bundles on \tilde{X} .
- Rank 2 bundles pulled back from B_1, B_2 .



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Operations:

- Tensor product of bundles.
- Sums of bundles.



Constructing Vector Bundles



Building blocks:

- Line bundles on \tilde{X} .
- Rank 2 bundles pulled back from B_1, B_2 .

Operations:

- Tensor product of bundles.
- ~~Sums of bundles~~ **Never (slope-) stable!**
- Extensions of bundles.



Our Solution

Define these two rank 2 vector bundles

$$\begin{aligned}\mathcal{V}_1 &\stackrel{\text{def}}{=} \chi_2 \mathcal{O}_{\tilde{X}}(-\tau_1 + \tau_2) \oplus \chi_2 \mathcal{O}_{\tilde{X}}(-\tau_1 + \tau_2) = \\ &= 2\chi_2 \mathcal{O}_{\tilde{X}}(-\tau_1 + \tau_2)\end{aligned}$$

$$\mathcal{V}_2 \stackrel{\text{def}}{=} \mathcal{O}_{\tilde{X}}(\tau_1 - \tau_2) \otimes \pi_2^*(\mathcal{W})$$

We define the rank 4 bundle \mathcal{V} finally as a generic extension

$$0 \longrightarrow \mathcal{V}_1 \longrightarrow \mathcal{V} \longrightarrow \mathcal{V}_2 \longrightarrow 0$$

Table of Contents

- Introduction
- Group Theory
- Calabi-Yau Manifolds
- Vector Bundles
- Low Energy Spectrum
 - Spectrum
 - Cohomology
 - Heterotic Anomaly
 - Spectral Sequences
- Conclusions

Low Energy Spectrum I

The massless spectrum
= zero modes of \mathcal{D}_{E_8}
= H^1 cohomology of the adjoint bundle $\mathcal{E}_8^{\mathcal{V}/G}$.

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$$\begin{aligned} H^1\left(X, \mathcal{E}_8^{\mathcal{V}/G}\right) &= \\ &= H^1\left(X, \mathcal{E}_8^{\mathcal{V}}/G\right) \end{aligned}$$

Low Energy Spectrum I

The massless spectrum
= zero modes of \mathcal{D}_{E_8}
= H^1 cohomology of the adjoint bundle $\mathcal{E}_8^{\mathcal{V}/G}$.

Work with

$$H^1\left(\tilde{X}, \mathcal{E}_8^{\mathcal{V}}\right)^G$$

=

Have in mind

$$\begin{aligned} H^1\left(X, \mathcal{E}_8^{\mathcal{V}/G}\right) &= \\ &= H^1\left(X, \mathcal{E}_8^{\mathcal{V}}/G\right) \end{aligned}$$

Low Energy Spectrum II



$$\underline{248} = (\underline{1}, \underline{45}) \oplus (\underline{15}, \underline{1}) \oplus (\underline{4}, \underline{16}) \oplus (\overline{\underline{4}}, \overline{\underline{16}}) \oplus (\underline{6}, \underline{10})$$

$$\begin{aligned} \underline{10} = & \chi_1(\underline{1}, \underline{2}, 3, 0) \oplus \chi_1\chi_2(\underline{3}, \underline{1}, -2, -2) \oplus \\ & \oplus \chi_1^2(\underline{1}, \underline{2}, -3, 0) \oplus \chi_1^2\chi_2^2(\underline{3}, \underline{1}, 2, 2) \end{aligned}$$

Correspondingly, the fermions split as...



Low Energy Spectrum II

$$\begin{aligned} \mathcal{E}_8^V = & \left(\mathcal{O}_{\tilde{X}} \otimes \theta(\underline{45}) \right) \oplus \left(\text{ad}(\mathcal{V}) \otimes \theta(\underline{1}) \right) \oplus \\ & \oplus \left(\mathcal{V} \otimes \theta(\underline{16}) \right) \oplus \left(\mathcal{V}^\vee \otimes \theta(\overline{16}) \right) \oplus \left(\wedge^2 \mathcal{V} \otimes \theta(\underline{10}) \right) \end{aligned}$$

where $\theta(\dots)$ is the trivial bundle.

$$\begin{aligned} \theta(\underline{10}) = & \left[\chi_1 \theta(\underline{1}, \underline{2}, 3, 0) \right] \oplus \left[\chi_1 \chi_2 \theta(\underline{3}, \underline{1}, -2, -2) \right] \oplus \\ & \oplus \left[\chi_1^2 \theta(\underline{1}, \underline{2}, -3, 0) \right] \oplus \left[\chi_1^2 \chi_2^2 \theta(\underline{3}, \underline{1}, 2, 2) \right] \end{aligned}$$

Low Energy Spectrum III

For example, focus on the fields in the 10:

$$\begin{aligned} H^1\left(\tilde{X}, \mathcal{E}_8^V\right)^G &= (\text{lots of other fields}) \oplus \\ &\oplus \left[\chi_1 \otimes H^1\left(\tilde{X}, \wedge^2 \mathcal{V}\right)\right]^G \otimes (\underline{\mathbf{1}}, \underline{\mathbf{2}}, 3, 0) \oplus \\ &\oplus \left[\chi_1 \chi_2 \otimes H^1\left(\tilde{X}, \wedge^2 \mathcal{V}\right)\right]^G \otimes (\underline{\mathbf{3}}, \underline{\mathbf{1}}, -2, -2) \oplus \\ &\oplus \left[\chi_1^2 \otimes H^1\left(\tilde{X}, \wedge^2 \mathcal{V}\right)\right]^G \otimes (\underline{\mathbf{1}}, \underline{\mathbf{2}}, -3, 0) \oplus \\ &\oplus \left[\chi_1^2 \chi_2^2 \otimes H^1\left(\tilde{X}, \wedge^2 \mathcal{V}\right)\right]^G \otimes (\overline{\mathbf{3}}, \underline{\mathbf{1}}, 2, 2) . \end{aligned}$$

Cohomology I



The necessary cohomology groups for \mathcal{V} are

$$H^1(\tilde{X}, \mathcal{V}) = 3 \operatorname{Reg}(G)$$

$$H^1(\tilde{X}, \mathcal{V}^\vee) = 0$$

$$\begin{aligned} H^1(\tilde{X}, \wedge^2 \mathcal{V}) &= H^1(\tilde{X}, \mathcal{V}_1 \otimes \mathcal{V}_2) = \\ &= 2 \oplus 2\chi_1 \oplus 2\chi_2 \oplus 2\chi_1^2 \oplus 2\chi_2^2 \oplus \\ &\quad \oplus 2\chi_1\chi_2^2 \oplus 2\chi_1^2\chi_2 \end{aligned}$$



Cohomology II



$$H^1\left(\tilde{X}, \wedge^2 \mathcal{V}\right) = \\ 2 \oplus 2\chi_1 \oplus 2\chi_2 \oplus 2\chi_1^2 \oplus 2\chi_2^2 \oplus 2\chi_1\chi_2^2 \oplus 2\chi_1^2\chi_2$$



Cohomology II



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$$2 = \left[\chi_1 \otimes H^1(\tilde{X}, \wedge^2 \mathcal{V}) \right]^G \quad \text{up Higgs}$$



Cohomology II



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Cohomology II

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Heterotic Anomaly I

So far, we only considered the visible E_8 bundle!

It turns out that we need an instanton \mathcal{H} in the hidden E_8 for anomaly cancellation.

$$c_2(T\tilde{X}) - c_2(\mathcal{V}) - c_2(\mathcal{H}) = PD(C) \in H^4(\tilde{X}, \mathbb{Z})$$

where C is the curve wrapped by five-branes.

$$c_2(T\tilde{X}) = 12 (\tau_1^2 + \tau_2^2)$$

$$c_2(\mathcal{V}) = -2\tau_1^2 + 7\tau_2^2 + 4\tau_1\tau_2$$

Heterotic Anomaly II

We found a rank 2 bundle with

- no hidden matter whatsoever.
- which requires five branes.

$$c_2(T\tilde{X}) - c_2(\mathcal{V}) - c_2(\mathcal{H}_{(1)}) = 6\tau_1^2.$$

Note: the curve τ_1^2 is just an elliptic fiber,

$$\tau_1^2 = \pi_1^{-1}(\{\text{pt.}\})$$

Heterotic Anomaly III

We found another (reducible) rank 4 bundle with

- two hidden $\text{Spin}(12)$ multiplets.
- which requires no five branes.

$$c_2(T\tilde{X}) - c_2(\mathcal{V}) - c_2(\mathcal{H}_{(2)}) = 0.$$

Spectral Sequences I

How did we compute all these cohomology groups?

Leray spectral sequence for any sheaf \mathcal{F} on $\tilde{X} \rightarrow B_2$:

$$E_2^{p,q} = H^p\left(B_2, R^q\pi_{2*}\mathcal{F}\right) \Rightarrow H^{p+q}\left(\tilde{X}, \mathcal{F}\right)$$

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Think of $E_2^{p,q}$ as the “forms with p legs along the base and q legs along the fiber”.

Spectral Sequences II

Example: $H^1\left(\tilde{X}, \wedge^2 \mathcal{V}\right) = H^1\left(\tilde{X}, 2\chi_2 \pi_2^*(\mathcal{W})\right)$

$$\pi_{2*}\left(2\chi_2 \pi_2^*(\mathcal{W})\right) = 2\chi_2 \mathcal{W}$$

$$R^1 \pi_{2*}\left(2\chi_2 \pi_2^*(\mathcal{W})\right) = 2\chi_1 \chi_2 \mathcal{W} \otimes \mathcal{O}_{B_2}(-f_2)$$

Compute $H^p(B_1, \dots)$ by two more Leray SS...

$$\Rightarrow E_2^{p,q} = \begin{array}{c|ccc} & \begin{array}{c} q=1 \\ \uparrow \\ q=0 \end{array} & \begin{array}{c} 0 \\ 0 \end{array} & \begin{array}{c} 2 \oplus 2\chi_1 \oplus 2\chi_2 \oplus 2\chi_1^2 \oplus 2\chi_2^2 \oplus 2\chi_1 \chi_2^2 \oplus 2\chi_1^2 \chi_2 \\ 2 \oplus 2\chi_1 \oplus 2\chi_2 \oplus 2\chi_1^2 \oplus 2\chi_2^2 \oplus 2\chi_1 \chi_2^2 \oplus 2\chi_1^2 \chi_2 \end{array} & \begin{array}{c} 0 \\ 0 \end{array} \\ \begin{array}{c} p=0 \\ \leftarrow \\ \rightarrow \\ p=1 \end{array} & & & & \begin{array}{c} p=2 \\ \rightarrow \end{array} \end{array}$$

Table of Contents

- Introduction
- Group Theory
- Calabi-Yau Manifolds
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- Low Energy Spectrum
- Conclusions
 - Important Lessons
 - Equivariant Structures
 - Summary

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 \Rightarrow no enhanced spectrum.

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- Unique solution?

- Equivariant actions are the key.

Photograph



Equivariant Structures I

Recall our solution: $0 \longrightarrow \mathcal{V}_1 \longrightarrow \mathcal{V} \longrightarrow \mathcal{V}_2 \longrightarrow 0$

$$\mathcal{V}_1 \stackrel{\text{def}}{=} \chi_2 \mathcal{O}_{\tilde{X}}(-\tau_1 + \tau_2) \oplus \chi_2 \mathcal{O}_{\tilde{X}}(-\tau_1 + \tau_2)$$

$$\mathcal{V}_2 \stackrel{\text{def}}{=} \mathcal{O}_{\tilde{X}}(\tau_1 - \tau_2) \otimes \pi_2^*(\mathcal{W})$$

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not a bundle!

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Constraint on characters:

$SU(4)$ bundle, not $U(4)$

$$\Leftrightarrow \wedge^4 \mathcal{V} = \mathcal{O}_{\tilde{X}}.$$

Without this, \mathcal{W} is
not a bundle!

Equivariant Structures II

Different G -action:

$$0 \longrightarrow \mathcal{V}'_1 \longrightarrow \mathcal{V}' \longrightarrow \mathcal{V}'_2 \longrightarrow 0$$

$$\mathcal{V}'_1 \stackrel{\text{def}}{=} \bigcirc \mathcal{O}_{\tilde{X}}(-\tau_1 + \tau_2) \oplus \bigcirc \chi_2^2 \mathcal{O}_{\tilde{X}}(-\tau_1 + \tau_2)$$

$$\mathcal{V}'_2 \stackrel{\text{def}}{=} \mathcal{O}_{\tilde{X}}(\tau_1 - \tau_2) \otimes \pi_2^*(\mathcal{W}')$$

$$0 \longrightarrow \mathcal{O}_{B_2}(-2f_2) \longrightarrow \mathcal{W}' \longrightarrow \bigcirc \chi_2 \mathcal{O}_{B_2}(2f_2) \otimes I_9 \longrightarrow 0$$

This is a different equivariant bundle on \tilde{X}

\Rightarrow a different quotient bundle on X !

Equivariant Structures III

$$\begin{aligned} H^1\left(\tilde{X}, \wedge^2 \mathcal{V}\right) &= 2 \oplus 2\chi_1 \oplus 2\chi_2 \oplus 2\chi_1^2 \oplus 2\chi_2^2 \oplus \\ &\quad \oplus 2\chi_1\chi_2^2 \oplus 2\chi_1^2\chi_2 \\ H^1\left(\tilde{X}, \wedge^2 \mathcal{V}'\right) &= 2 \oplus \chi_1 \oplus \chi_1^2 \oplus 2\chi_2 \oplus 2\chi_1\chi_2 \oplus \\ &\quad \oplus \chi_1^2\chi_2 \oplus 2\chi_2^2 \oplus \chi_1\chi_2^2 \oplus 2\chi_1^2\chi_2^2 \end{aligned}$$

Different spectrum:

Equivariant Structures III

$$H^1(\tilde{X}, \wedge^2 \mathcal{V}) = 2 \oplus 2\chi_1 \oplus 2\chi_2 \oplus 2\chi_1^2 \oplus 2\chi_2^2 \oplus \\ \oplus 2\chi_1\chi_2^2 \oplus 2\chi_1^2\chi_2$$

$$H^1(\tilde{X}, \wedge^2 \mathcal{V}') = 2 \oplus \chi_1 \oplus \chi_1^2 \oplus 2\chi_2 \oplus 2\chi_1\chi_2 \oplus \\ \oplus \chi_1^2\chi_2 \oplus 2\chi_2^2 \oplus \chi_1\chi_2^2 \oplus 2\chi_1^2\chi_2^2$$

Different spectrum:

One up Higgs, one down Higgs,

Equivariant Structures III

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Different spectrum:

One up Higgs, one down Higgs, two $\underline{3}$, and two $\overline{3}$.

Spectrum Summary

- 3 families of quarks and leptons.
- Zero anti-families.
- 4 Higgs (twice MSSM).
- Doublets and triplets are completely split, all triplets are projected out.
- Hidden pure E_7 or $\text{Spin}(12)$ with 2 matter fields.

Within our Ansatz, there is only one solution that is even close to a realistic spectrum.

Future Directions

- Supersymmetry breaking.
- $U(1)_{B-L}$ breaking.
- Yukawa couplings.
- Moduli stabilization.
- Revisit $SU(5)$ with \mathbb{Z}_2 Wilson line: no $U(1)_{B-L}$.