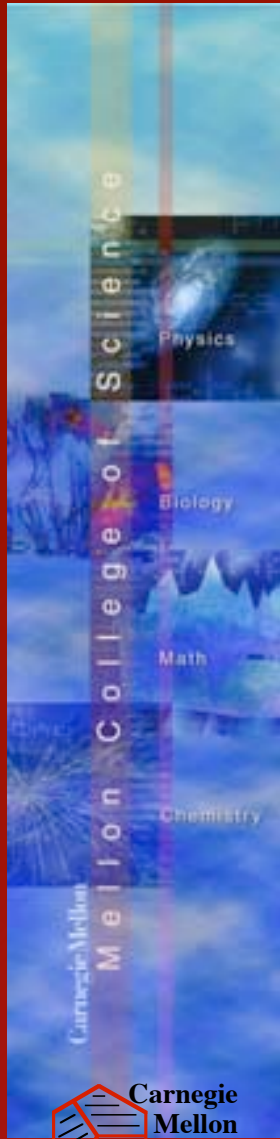


the mesoscale view of interfaces



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joint work with

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Anthony Rollett

Shlomo Ta'asan

Peng Yu (starting Penn State)

supported by the NSF under the MRSEC
program the DMS, and the DoE



Acknowledgments

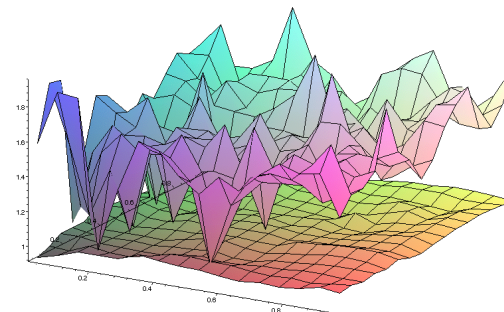
Chun Liu
Darren Mason
W.W. Mullins (†)
K. Barmak
G. Rohrer

References and related work

K, Livshits, Manolache, Rollett & Ta'asan, MRS 652, 2001
Adams, et al., Int. Sci. 1999
Anderson, Srolovitz, et al. Acta Met., 1984
Demirel & Rollett, MRS 652, 2001
Haslam, et al., Mat Sci. Eng. A.
K & Liu, Mat. Met. Mod. Appl. Sci., 2001
K, Livshits, Mason & Ta'asan, Int. Sci. 2002
K, Livshits & Ta'asan, 2003
Kuprat, SIAM Sci. Comp., 2000
Mullins, Acta Met. 1998
Ta'asan, Yu, Livshits, K & Lee

C N A Summer School
May 27 - June 5, 2004
Pittsburgh

Luigi Ambrosio, Scuola Normale Superiore, Pisa (ITALY)
Maria-Carme Calderer, U. of Minnesota, Minneapolis
Irene Fonseca, CMU
Wilfrid Gangbo, Georgia Tech
Nassif Ghoussoub, U. British Columbia (CANADA)
Stefan Müller, Max Planck Institute for Mathematics in the
Sciences, Leipzig (GERMANY)



Motivating Physical Phenomena

The performance of a polycrystalline material is influenced by the types of grain boundaries in the material and the way that they are connected.

Examples:

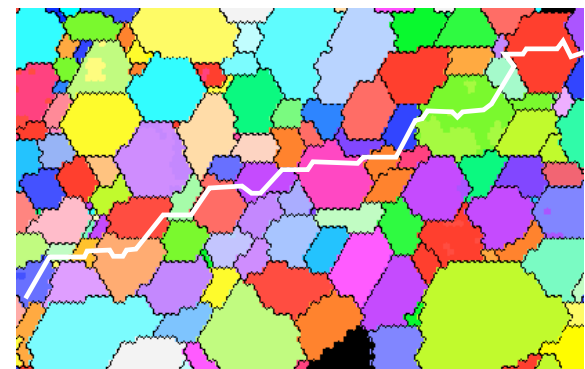
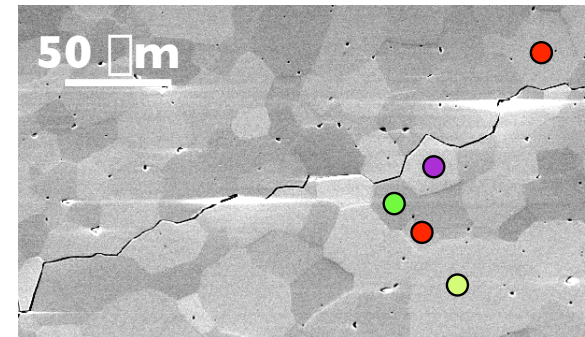
Superconducting Critical Current Density

Electromigration Damage Resistance

Stress Corrosion Cracking

Electrical activity

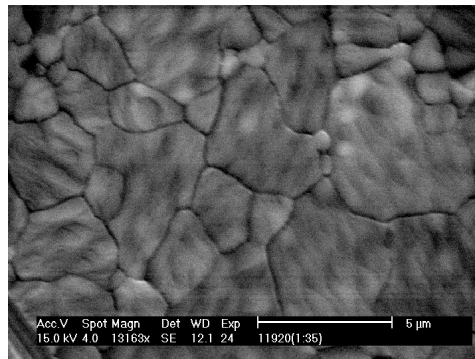
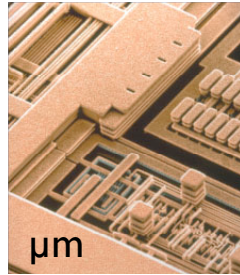
Creep Behavior



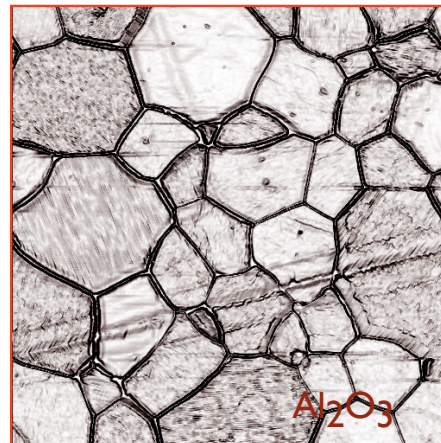
fracture follows
grain boundaries

Viewpoint: Multiscale

- use and occurrence



Al

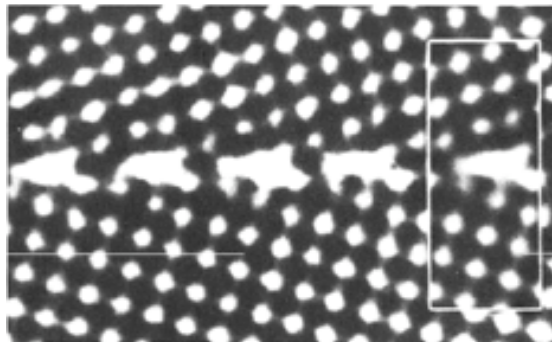


$100\ \mu\text{m}$



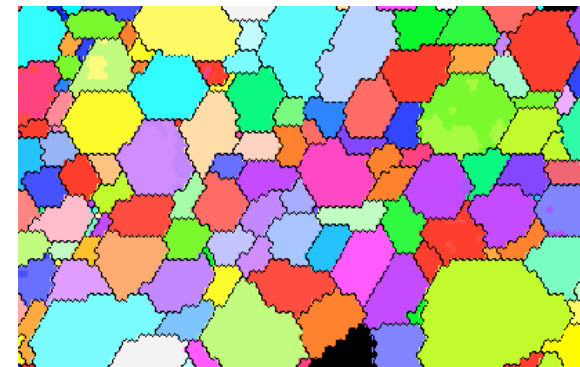
cm

- interrogation



TEM

or orientation imaging
microscopy (shown on
previous slide)



- theory and simulation

atomic level

embedded atom methods

first principle computations

mesoscale

thermodynamics of surfaces (Mullins, Herring, et seq.)

simulations via random methods (Monte Carlo or Potts,
Anderson, Srolovitz, Grest & Sahni, et seq.)

simulations via resolution of the time dependent equations of motion

ensembles of grains

theories of the statistics of grain growth (very large literature)

- rising role of automated data acquisition in materials science
information arrives at a scale at which it can be acquired:
typically a mesoscale
here the scale of geometric and crystallographic information
involves ingredients in relations rather than the desired quantities
- an integral and essential role for simulation and modeling — today's focus

Poster child for grain boundary engineering: Pb for batteries from Ontario Hydro

Changing the Grain Boundary Network Improves Performance



Both electrodes polarized at 200 mV
for 12 days (70°C) in H_2SO_4 .

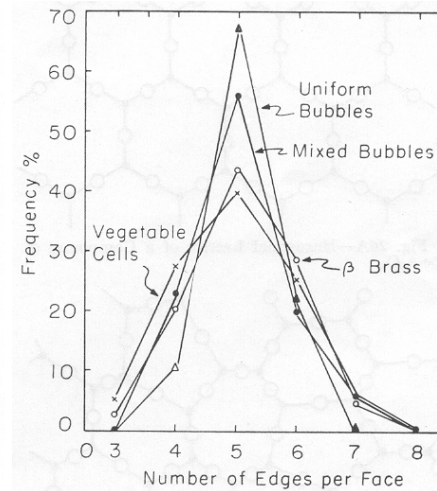
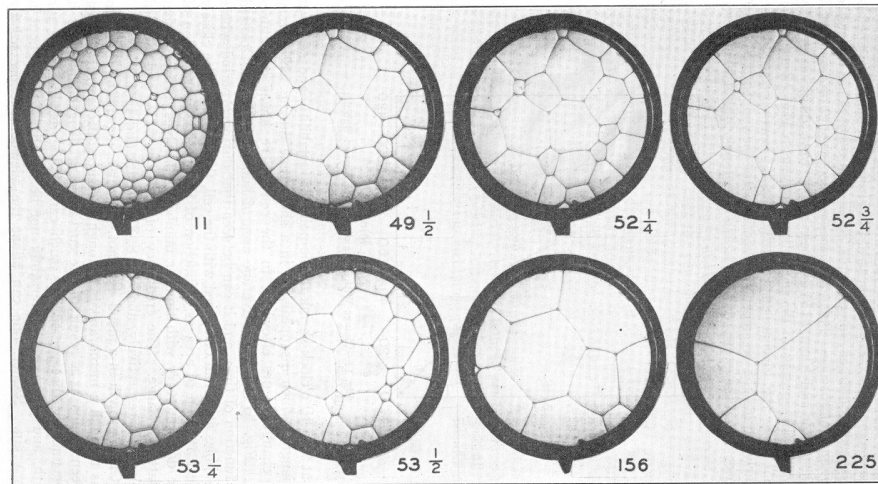
The as cast electrode disintegrated
during the test.

The GBETM electrode maintained its
integrity.

As cast = 10% special boundaries
GBETM = 71% special boundaries

Lehockey, Palumbo, Lin, and Brennenstuhl, Met. Trans. 29A (1998)
387.

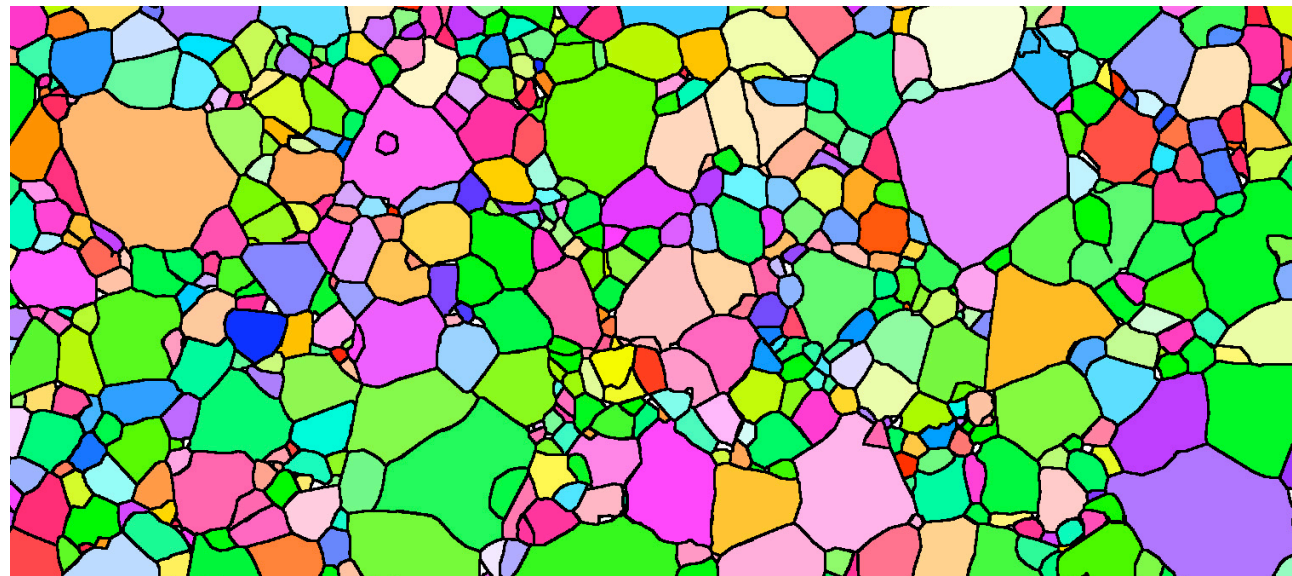
Are microstructures like soap froth?



C.S. Smith,
1951
more alike
than unlike

late stage MgO

very unlike
yet appearances
can deceive



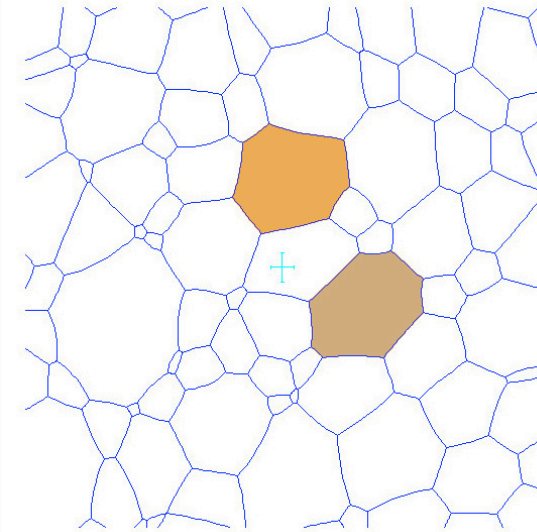
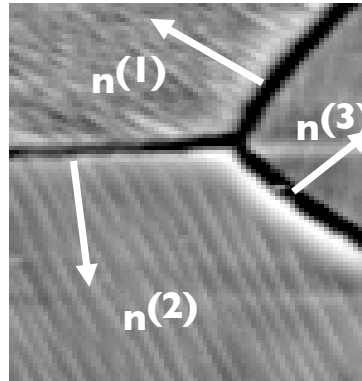
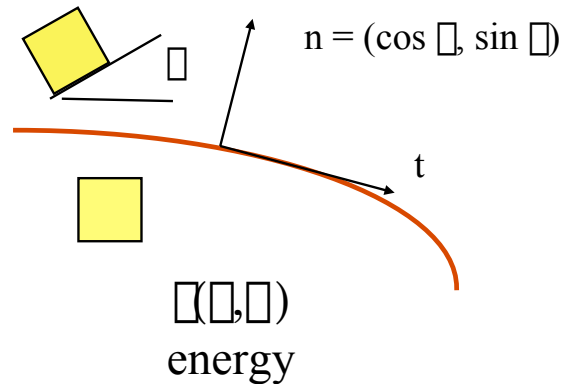
Frank and Ernest

By Bob Thaves



Grain Growth in 2D everything in 2D today

equilibrium theory



E total energy
 γ gb energy
 n normal
 t tangent
 ϕ angle of normal
 $\Delta\phi$ misorientation
 κ curvature of ϕ

network of curves $\{ \phi^{(i)} \}$

local equilibrium of network

ϕ

$$\left(\frac{d^2 \phi}{d\Gamma^2} + \kappa \phi \right) = 0 \quad \text{on } \phi^{(i)}$$

$$\sum_{TJ} \left(\frac{\partial \phi}{\partial \Gamma} n + \Gamma t \right) = 0 \quad \text{at TJ's}$$

$$E = \sum_i \gamma_i \int |t| ds$$

$$\delta E = 0$$

Herring Condition

dynamics of grain growth Mullins

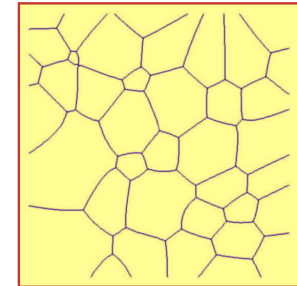
curvature driven growth dynamical problem:

must impose boundary conditions

use Herring Relation, the natural boundary condition for equilibrium

$$\left| \begin{array}{l} v_n = \mu \left(\frac{\partial^2 \gamma}{\partial \theta^2} + \gamma \right) \kappa \quad \text{on } \Gamma^{(i)} \\ \sum_{TJ} \left(\frac{\partial \gamma}{\partial \theta} n + \gamma t \right) = 0 \quad \text{at } TJ's \end{array} \right.$$

$\mu > 0$ mobility
different from tension



$$E = \sum \int_{\Gamma^{(i)}} \gamma |t| ds$$

$$\frac{d}{dt} E = - \sum \int_{\Gamma^{(i)}} \frac{1}{\mu} (v_n)^2 |t| ds + \sum v \cdot \sum_{TJ} \left(\frac{\partial \gamma}{\partial \theta} n + \gamma t \right)$$

$$= - \sum \int_{\Gamma^{(i)}} \frac{1}{\mu} (v_n)^2 |t| ds < 0$$

vanishes when
Herring holds

dynamical system is
dissipative

curvature driven growth

Herring Relation at TJ's

tends to equilibrium but still 'metastable'

dissipative system

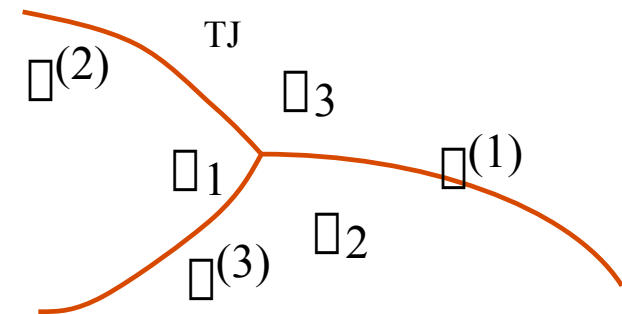
trend to

equilibrium

- When $\gamma = \text{constant}$, Herring condition means segments meet at $2\pi/3$.

When $\gamma = \gamma(\alpha)$, Herring condition is 'Young's Law'

$$\frac{\gamma(\alpha_1)}{\sin(\psi_1)} = \frac{\gamma(\alpha_2)}{\sin(\psi_2)} = \frac{\gamma(\alpha_3)}{\sin(\psi_3)}$$

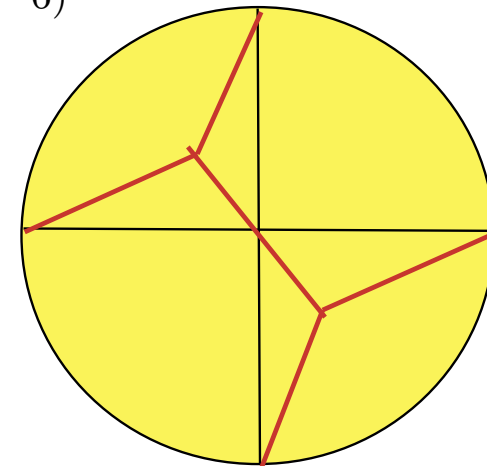


- Mullins – Von Neumann $n - 6$ rule

$\gamma = \text{constant}$, Herring condition

large grains grow and small grains shrink

$$\frac{dA}{dt} = \alpha(n - 6)$$



- Typically triple junctions are stable
- Much analysis related to curve shortening, approximations of multiphase boundaries (Cahn Hilliard, Allen Cahn, ...), and Wulff-type problems (Taylor, Fonseca, ...)
- Direct precursors:
 - Bronsard & Reitich short time local existence: system satisfies complementing conditions, very important!
 - K & Liu long time existence close starting near a stationary solution

excursion



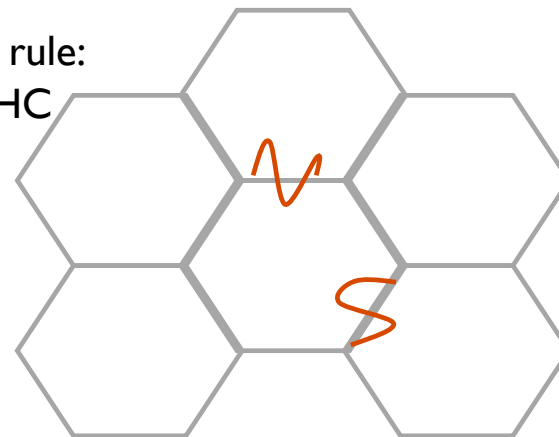
THE GIANT'S CAUSEWAY

Giant's Causeway (Ireland) probably cleavage/fracture

suggests interesting example
based on Mullins-Von Neumann n-6 rule:
deform curves conserving area and HC
 $\langle A \rangle = \text{constant}$
although system in motion

$E = \min \rightarrow$ no interfaces

illustration of metastability
of grain growth system



Simulation and interpretation: strategy

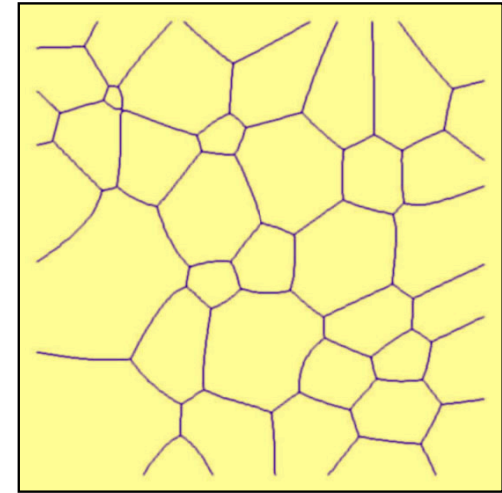
- must fulfill two requirements
 - accurate: fidelity to Mullins-Herring
 - statistical significance: very large scale
- interpret solution through statistics - only possibility
- derive coarse grained descriptions of the simulation useful for prediction
 - equation satisfied by distribution function (histograms of relative area)
 - system of equations satisfied by partial distributions functions (by number of grain facets), which we term a master equation model
- however other issues also intervene (will return to this): we will have some dynamic statistics but we need to ascertain their 'information' content
- identify the material - reconstruct the energy - from the simulation?
 - tantalizing new features

simulation

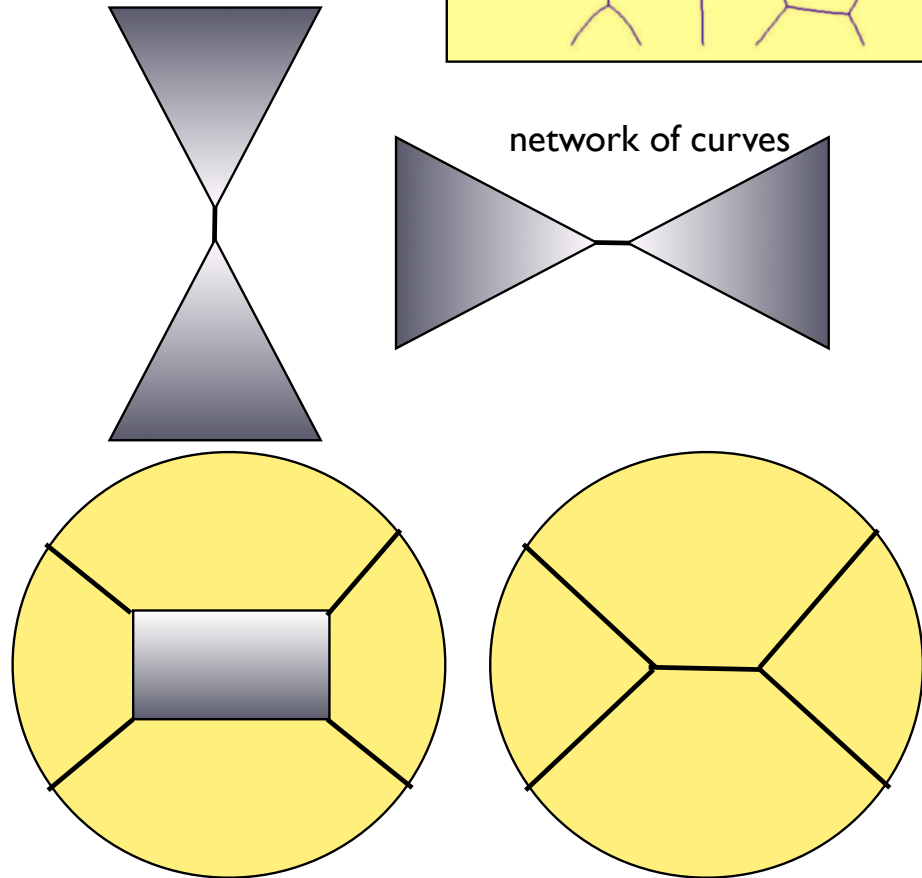
calibration $\gamma = 1$ $\gamma = 1$

evolve as network of curves:
data structure of inferior dimension
present data structure up to 50,000 grains
evolve to 1/4 starting number (typically)
and 1600 time steps (typically)

- strategy:
 - discretize energy with evolution
 - designed to maximize dissipation
 - maintain Herring BC
 - agrees with PDE to 2nd order in space
- critical events:
 - loss of facet (facet flipping)
 - loss of small grain
 - rules based on results of Monte Carlo MD simulations and geometry

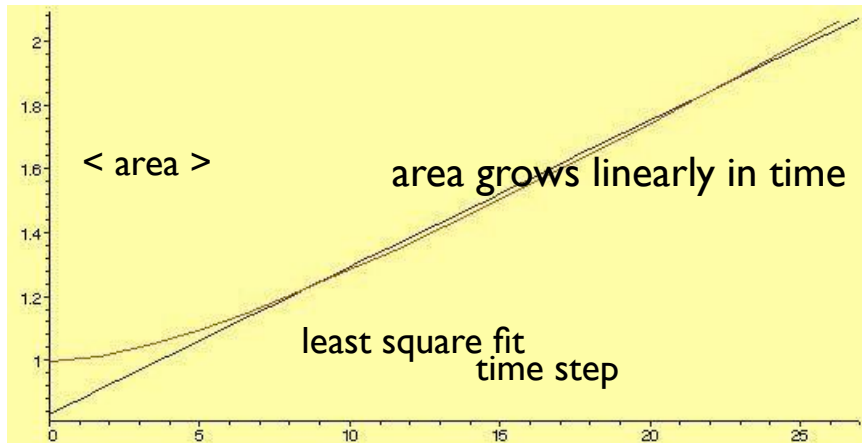


network of curves



First look at results

general features and diagnostics



diagnostics

area: grows linearly in time

n-6 rule: satisfied for individual grains

not subject to critical events

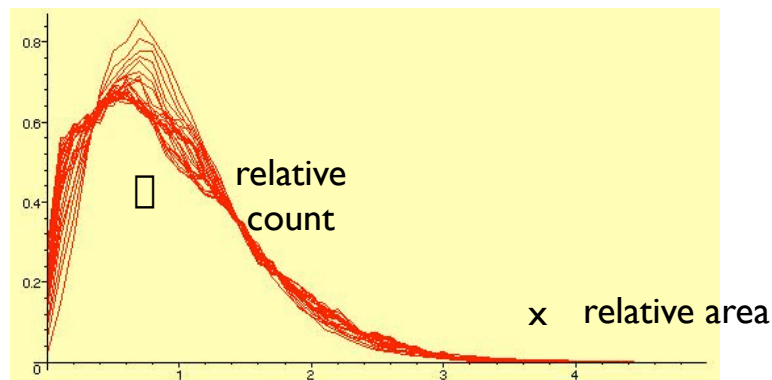
curvature: second order accurate with respect to spatial discretization

(boundary condition is very important)

25,000 grains

histograms at time steps

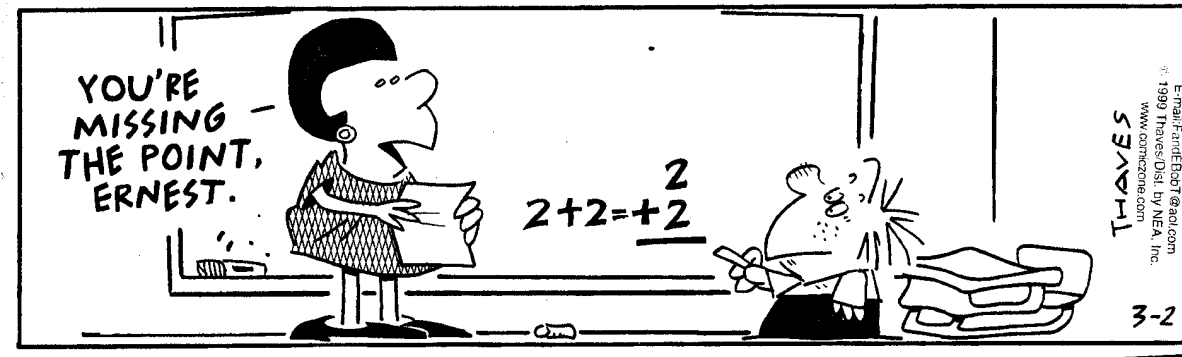
1 - 16 x 1000



*surprisingly high
degree of self-similarity
robust across sample size
and system size*

Frank and Ernest

By Bob Thave



I hope that I shall not miss the point

relative area histograms and their interpretation

interpretation of simulation
simulating metastable systems

interpret simulation: find an equation histograms satisfy
simulated histogram is self similar over a long range but ultimately is dynamic quantity

$$f(\text{Area}, t) = g(\text{Area}/t) \text{ self similar form}$$

try to find by inverse methods

many theories about equations for ρ , generally have form:

$$\frac{\partial \rho}{\partial t} = \frac{\partial}{\partial x} (\psi' \rho)$$

possibly nonlinear transport

Hillert

Mullins

Ryum & Hunderi

$$\frac{\partial \rho}{\partial t} = \sigma \frac{\partial^2 \rho}{\partial x^2} + \frac{\partial}{\partial x} (\psi' \rho)$$

diffusion or Fokker-Planck

Louat

Atkinson no known physical reason for σ

Mullins σ cannot have origins at molecular time scales

but information loss/disorder can give rise to entropy
entropy in an equation is manifested by diffusion
...know from many directions

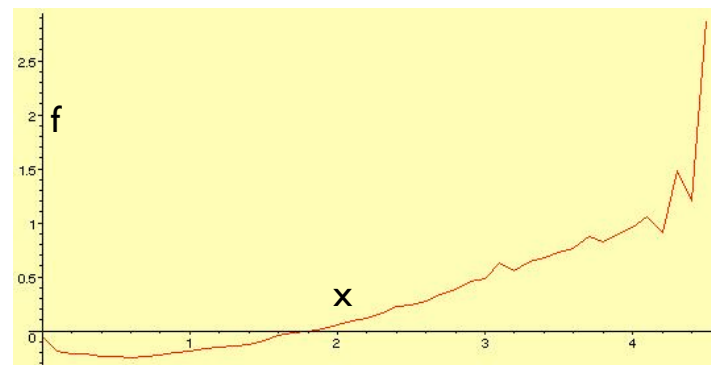
will look for F-P equation

inverse method to discover Fokker-Planck

a. identify an equilibrium solution



equilibrium solution \square



shape of potential f

$$\sigma \frac{\partial^2}{\partial x^2} \phi + \frac{\partial}{\partial x} (b\phi) = 0 \quad x > 0$$
$$\Rightarrow$$
$$\phi(x) = \frac{1}{Z} \exp(-\psi(x)/\sigma) \quad \text{with } b = \psi'$$

can determine a pattern or shape

$$f(x) = -\log \phi(x) \quad \text{and} \quad \psi(x) = \sigma f(x)$$

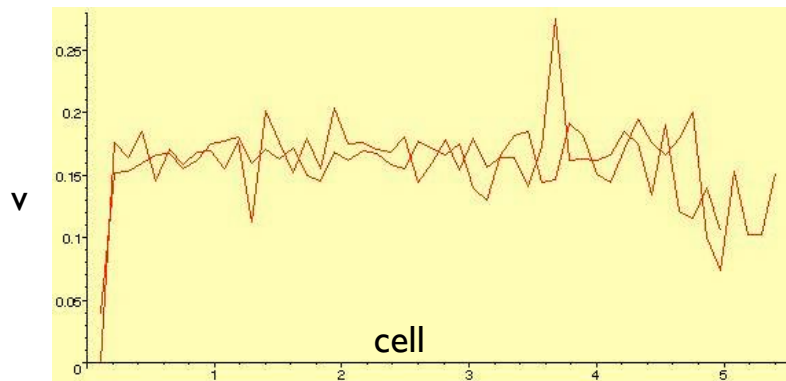
b. find diffusion coefficient
recall how $\varphi(x,t)$ determined

can regard the information that gives the
relative histograms as data for determining
the equation

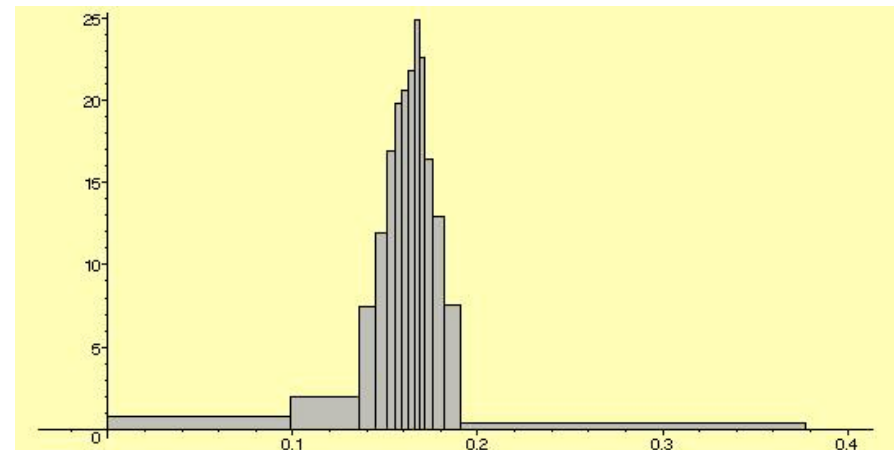


$$\text{diffusion coefficient} = \sigma = \frac{1}{2} \frac{v}{\Delta t} = 0.17$$

and $\Delta t = h^2$ by scaling



variances in cells at several times

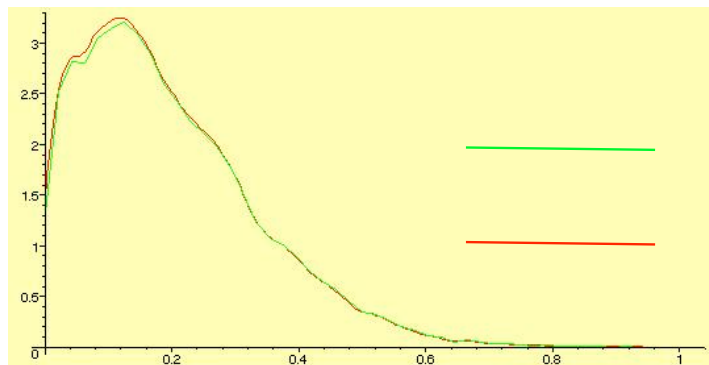
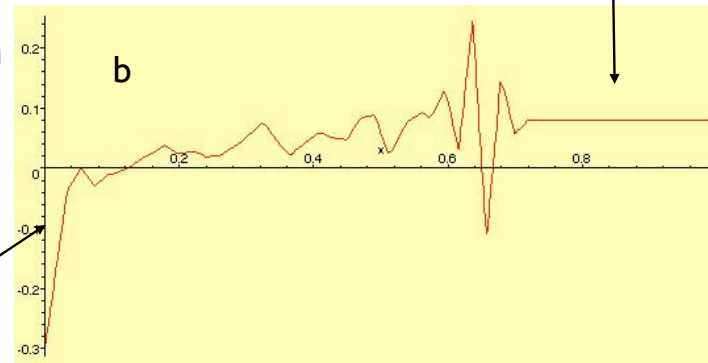


histogram of empirical variances (diff coeffs)
over simulation: **nearly constant**

final step:

for derived ρ and ϕ solve the F-P equation
for time = transient time

(unable to significantly vary initial configurations
but ...)



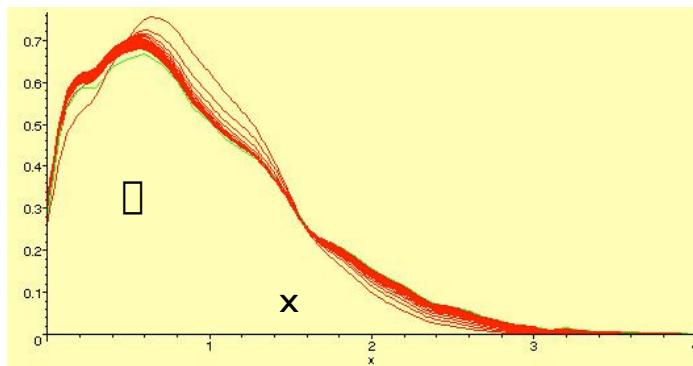
distribution

ρ
from simulation

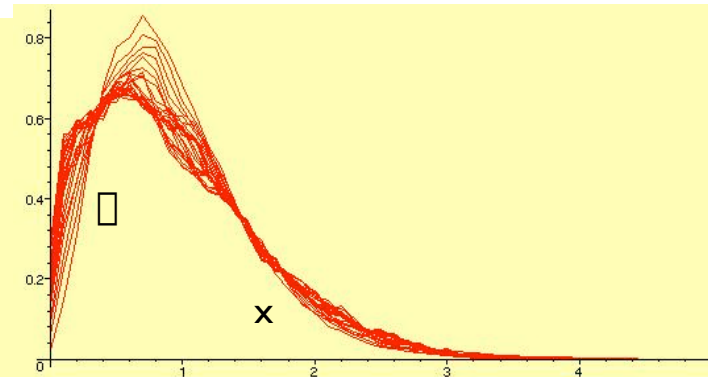
ρ
from solving F-P

approximate b
to solve F-P
arrows point to deviations
from the differenced ρ

this checks that we obtain
the correct equilibrium state

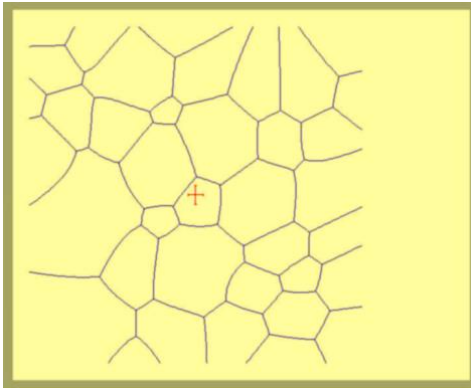


solution of Fokker-Planck
simulation



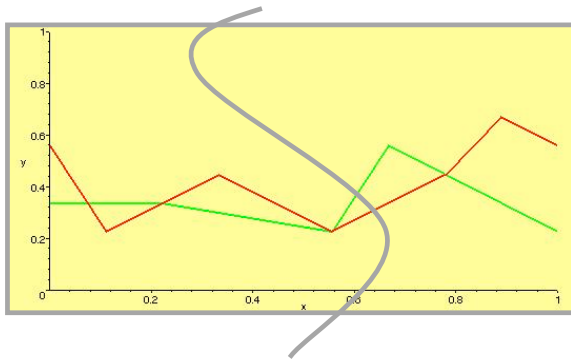
grain boundary
simulation histograms

have managed to approximately capture the transient in fact, there are other considerations



large scale simulation

computation at 'microscale'
need to interpret



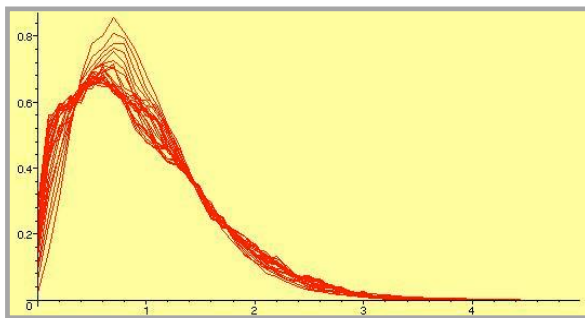
master equation description

'mesoscale'

grain trajectories: Mullins von Neumann $n - 6$ rule

$$\frac{dA}{dt} = \alpha(n - 6)$$

interrupted by random edge loss/gain or grain disappearance



relative area histograms

one point statistics

suggests Fokker-Planck

Interpretation of system: at intermediate scale by master equation model

$F_n(x,t)$ = density of n -facet grains with area x at time t

$$\frac{\partial}{\partial t} F_n(x,t) + c(n-6) \frac{\partial}{\partial x} F_n(x,t) = I_n(x,t)$$

$$F_n(N_a, t) = 0, \quad n = 3, 4, 5 \quad \text{and} \quad F_n(0, t) = 0 \quad n = 6, \dots$$

$$I_n = -p_n F_n + p_{n+1} F_{n+1} - (q_n + r_n) F_n + q_{n+1} F_{n+1} - r_{n-1} F_{n-1}$$

$p_n(x,t)$ = facet loss from grain disappearance (depends on $F_n(0,t)$)

$q_n(x)$ = facet loss from grain boundary flipping

$r_n(x)$ = facet gain from grain boundary flipping

determine from the the first
few time steps of the simulation
tend to vary with $1/\langle \text{area} \rangle$

earlier work by

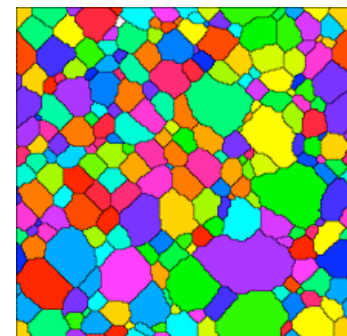
Fradkov

Flyvberg

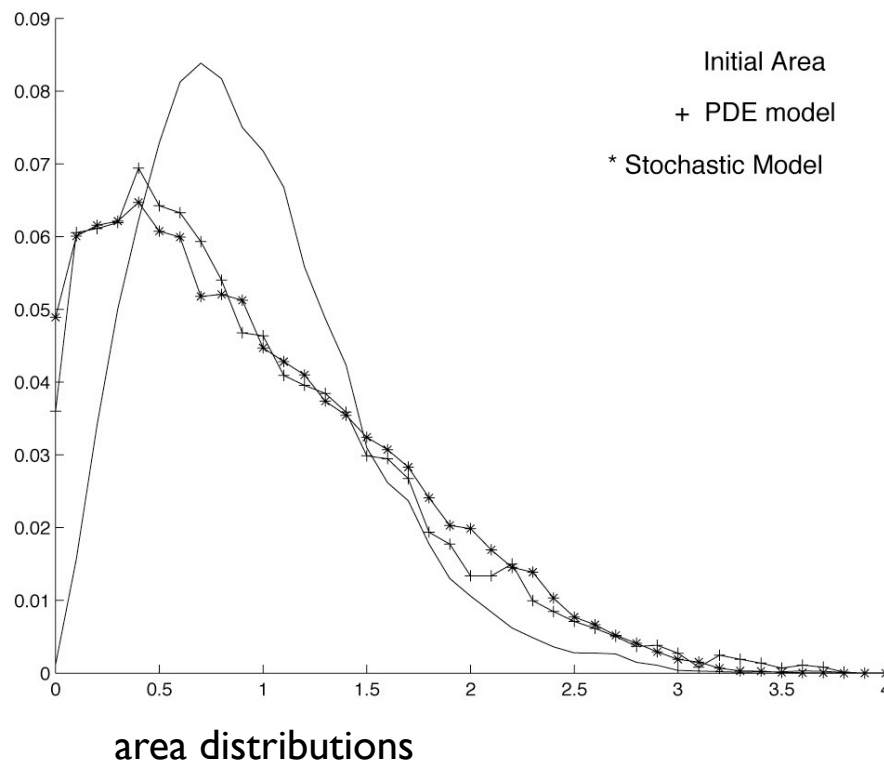
other descriptions:

Monte Carlo/Potts models

vertex models (also Henseler, Niethammer, Otto)



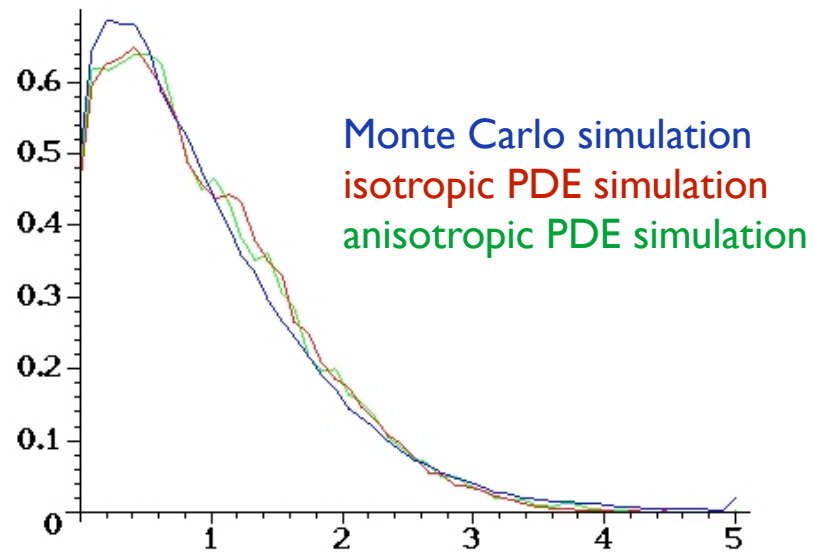
preliminary results



master equations give good
description of stationary
distribution

a few statistics from the initial
time period of the evolution,
determines stationary
distribution:
so in some fashion
can predict the stationary
distribution but do not understand
how to unravel this

only one difficulty



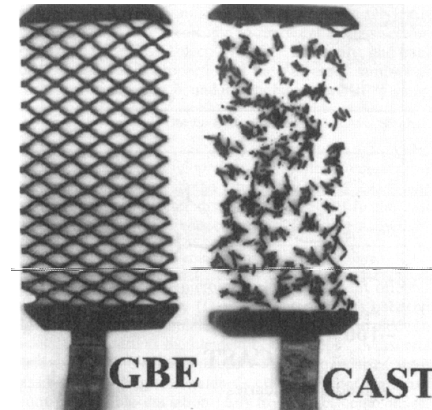
- stationary relative area histogram is extremely robust
- relative area histograms do not distinguish the energy (or mobility)

Crucial role of anisotropy

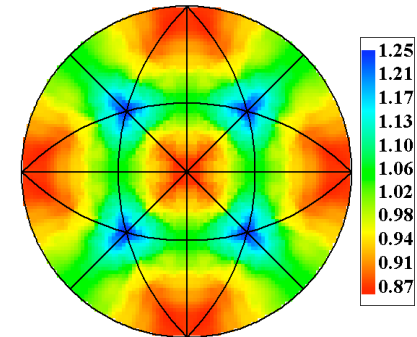
is anisotropy important?



fracture toughness optimized by specific arrangements of grains



corrosion resistance depends on high fraction of low energy boundaries



occurrence

Interrupted grain growth experiment on Al specimen to monitor change in $I(\Delta g, n)$ (relative normal averaged over misorientations) during growth.

what is the ideal distribution of boundaries in a well annealed sample?
all boundaries minimum energy?

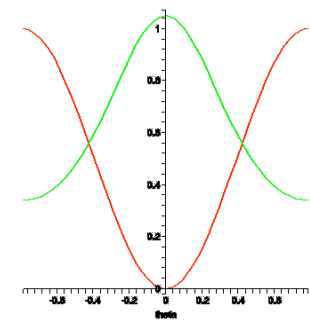
cannot satisfy Herring condition

independent trials with respect to energy?

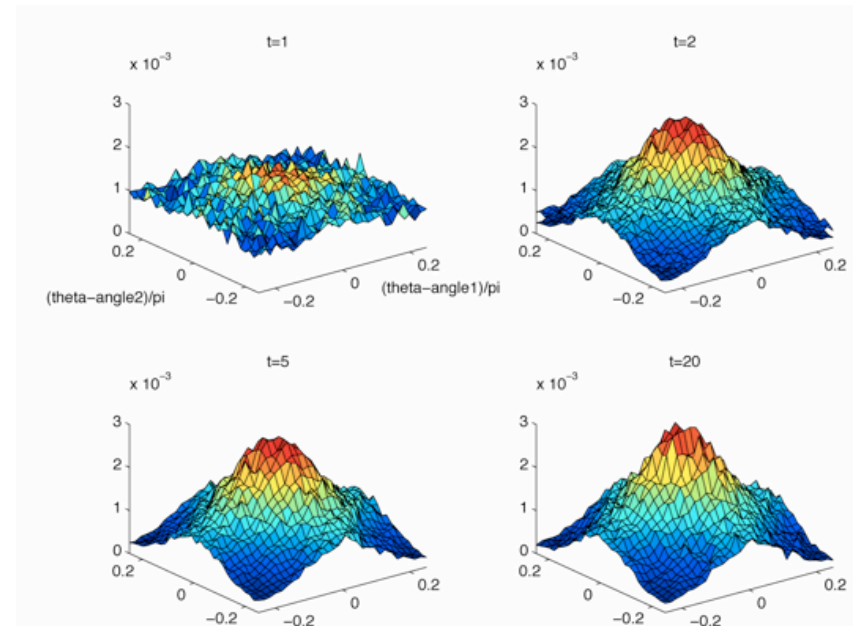
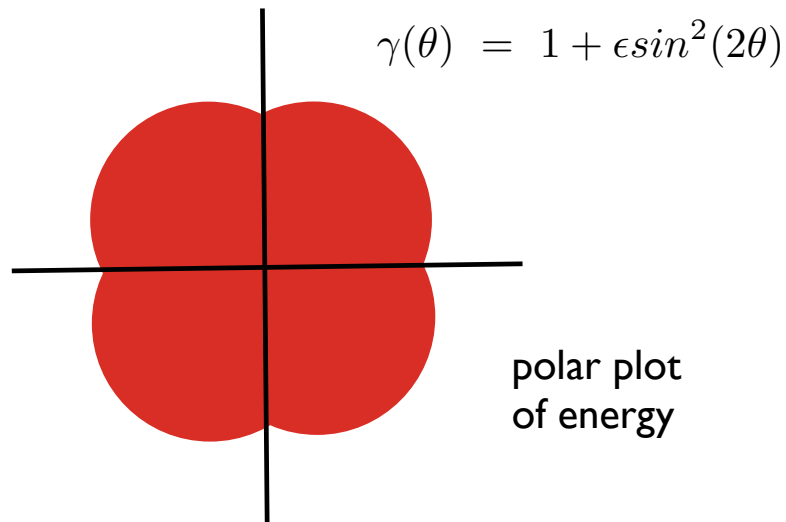
for $\gamma = \gamma(\alpha)$, $\alpha = \text{misorientation}$

$$\rho(\alpha) = e^{k\gamma(\alpha)} / Z$$

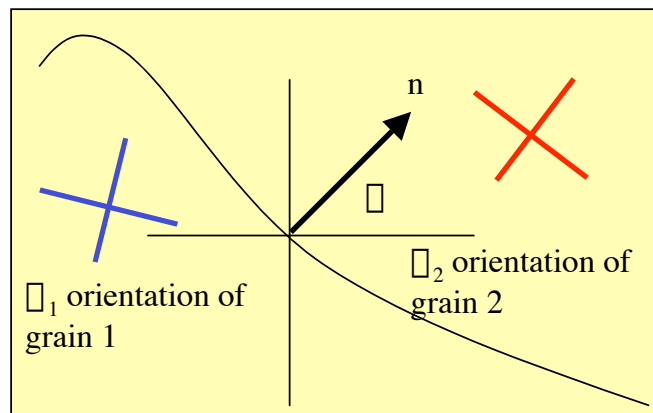
can verify this and confirm work of Holm et al., Upmanyu et al.



energy depends also on normal
 our's is the first (perhaps the only) with ability
 to execute large scale simulations with $\gamma = \gamma(\theta, \alpha)$

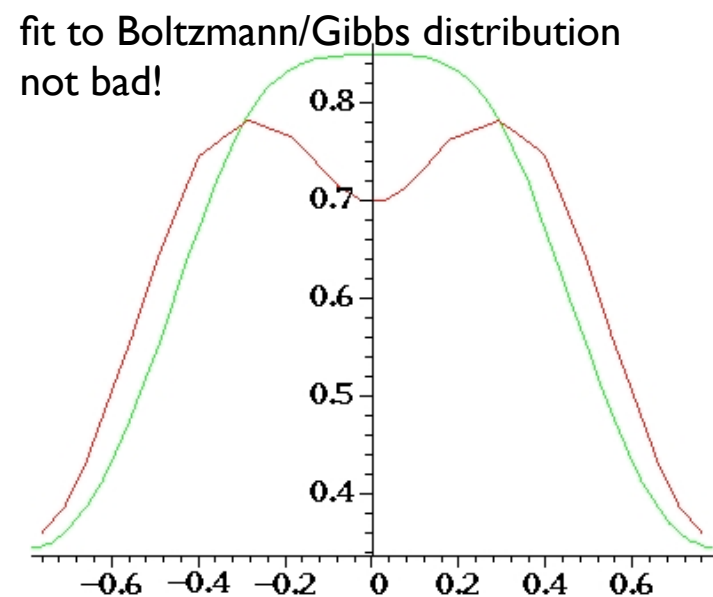
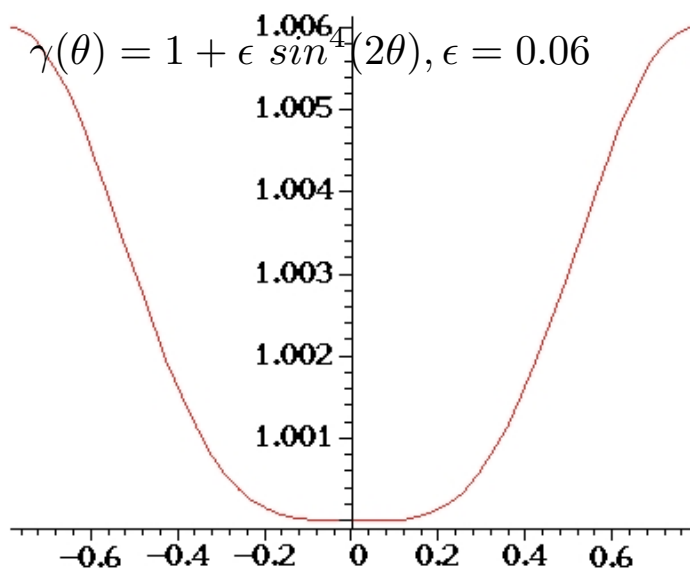
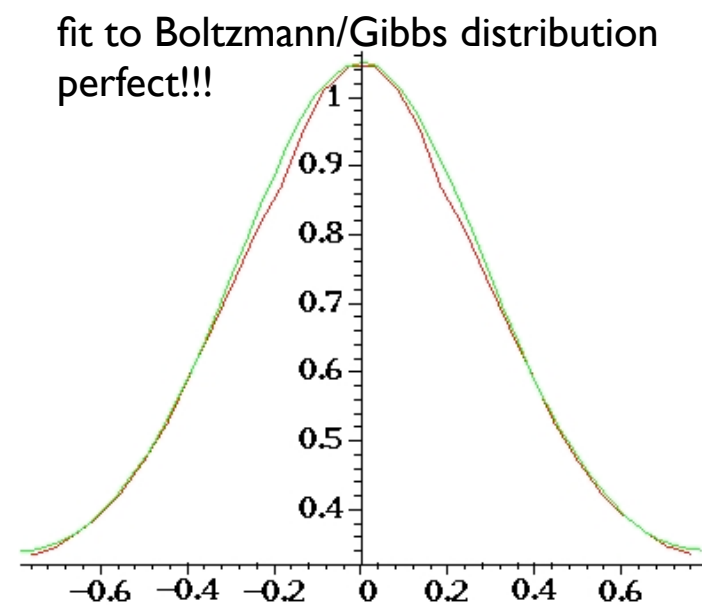
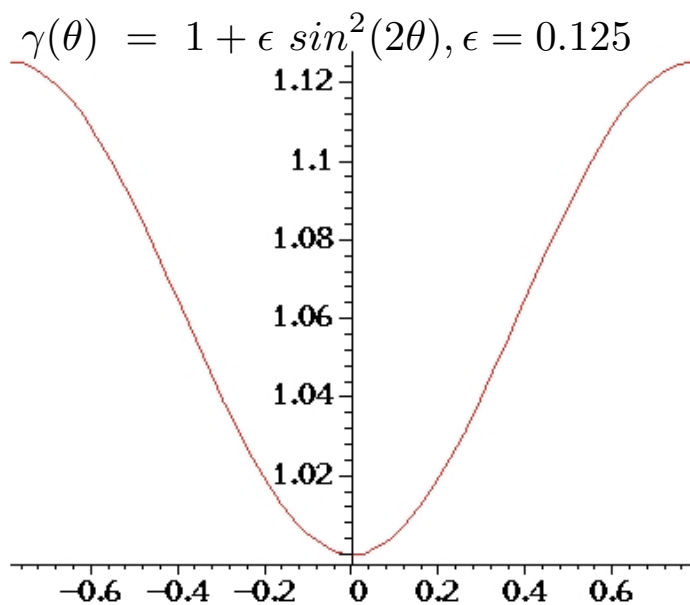


simulation shows development of a structure
 where low energy boundaries predominate

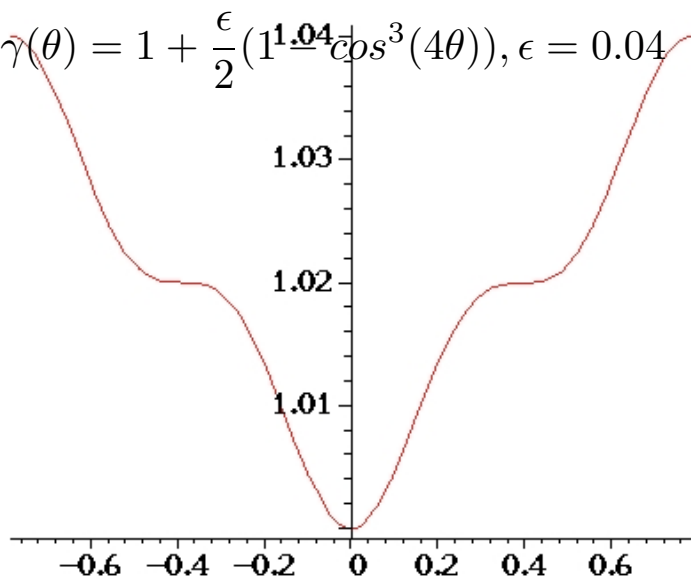


approximate grain boundary energy as average
 of two surface energies as seen in experiments
 on MgO:

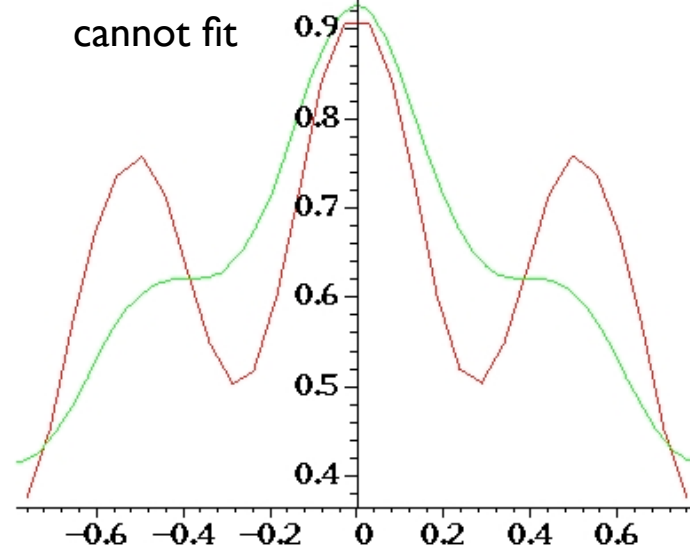
$$\gamma(\theta; \omega_1, \omega_2) = 1 + \epsilon(\sin^2(2(\theta - \omega_1)) + \sin^2(2(\theta - \omega_2)))$$



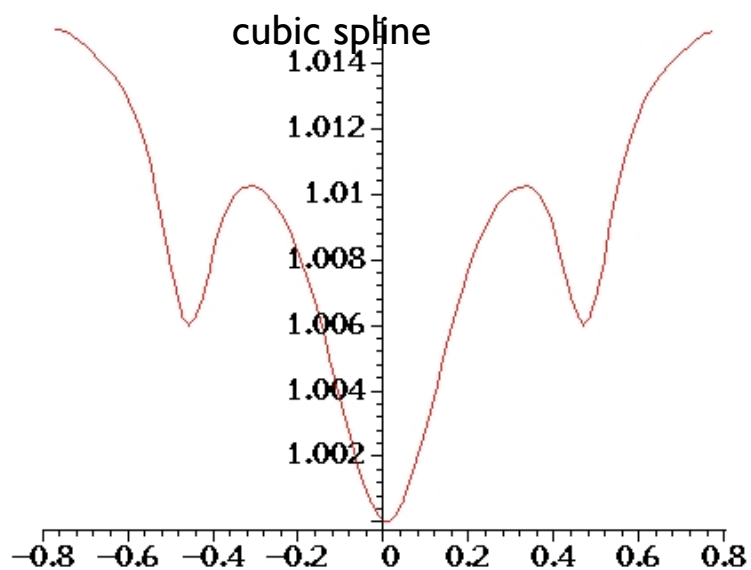
$$\gamma(\theta) = 1 + \frac{\epsilon}{2}(1 - \cos^3(4\theta)), \epsilon = 0.04$$



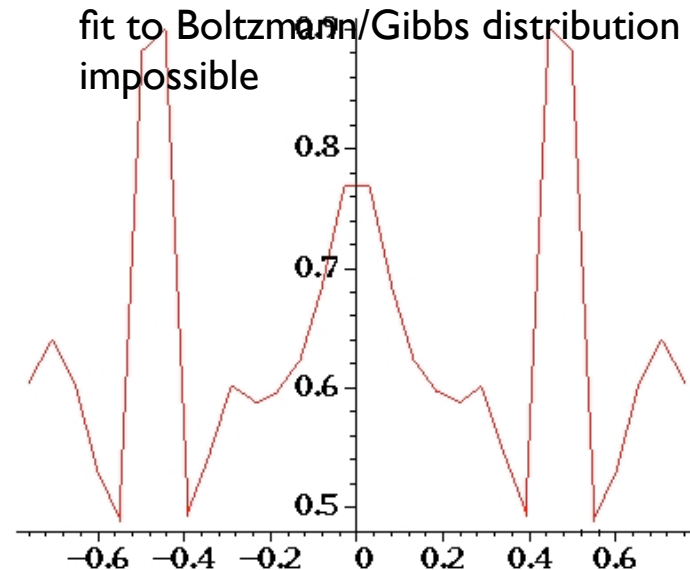
fit to Boltzmann/Gibbs distribution
cannot fit

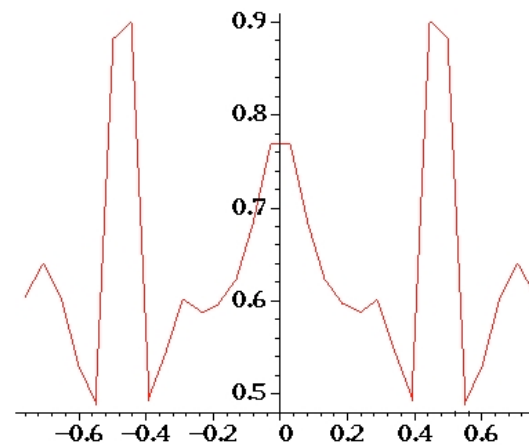
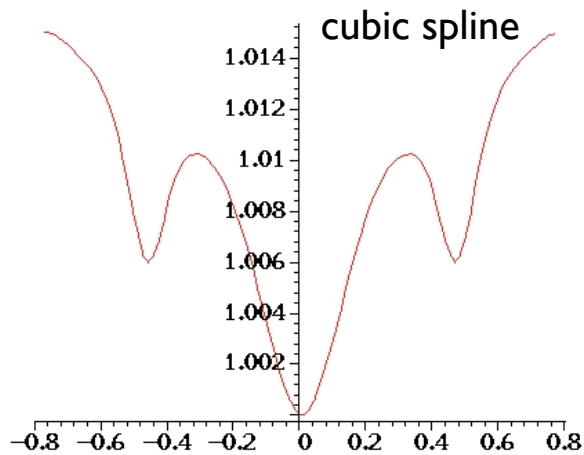


cubic spline



fit to Boltzmann/Gibbs distribution
impossible





- grain boundary character distribution depends on anisotropy
- simple cases may be interpreted as independent trials with respect to the Boltzmann/Gibbs distribution of the energy
- but all of our data for all of our materials show that distribution is not the Boltzmann/Gibbs
- complex energy landscape enables solution to remain in a region of phase space from which equilibrium is inaccessible
- understanding this issue will be a major goal of the project!