

An Organizing Center for Tracefiring in the Oregonator

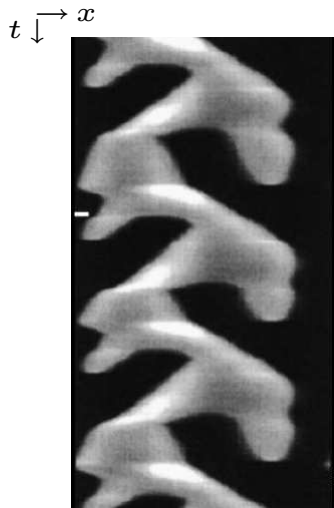
Building blocks and instability

Jens Rademacher

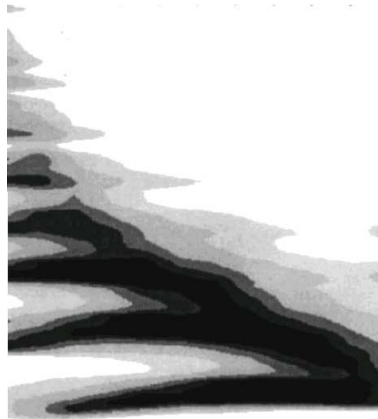
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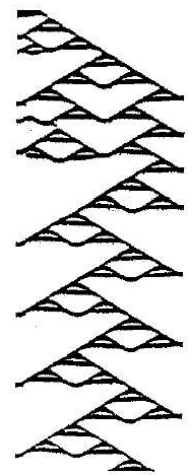
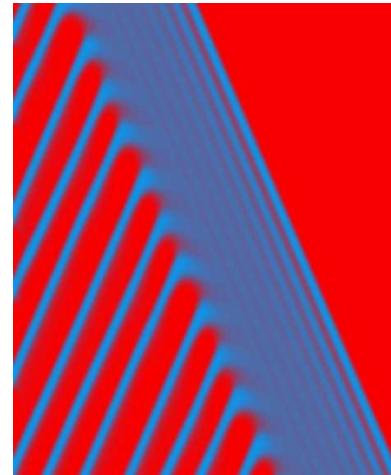
Backfiring and Tracefiring



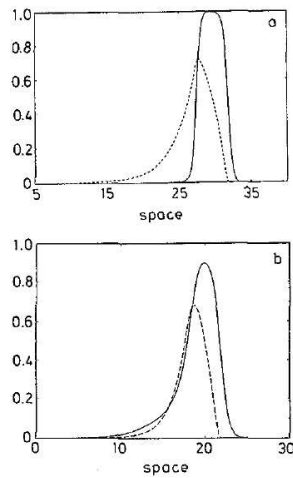
[Annamalai et al '99]
countercurr. flow reactor



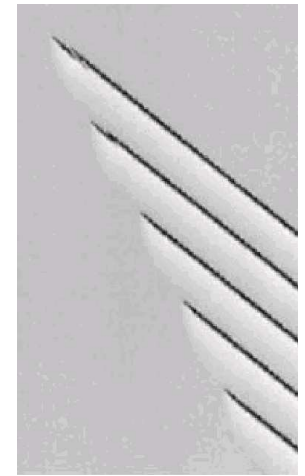
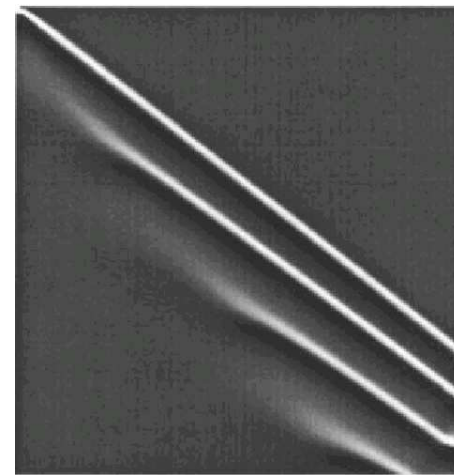
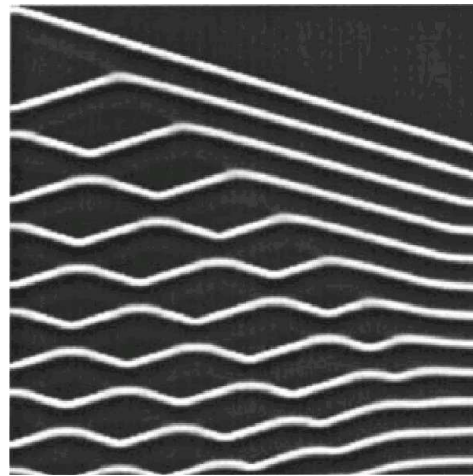
[Otterstedt et al '98]
active/passive cobalt



MB model for Co-ox. on Pt110: [Bär et al '94]



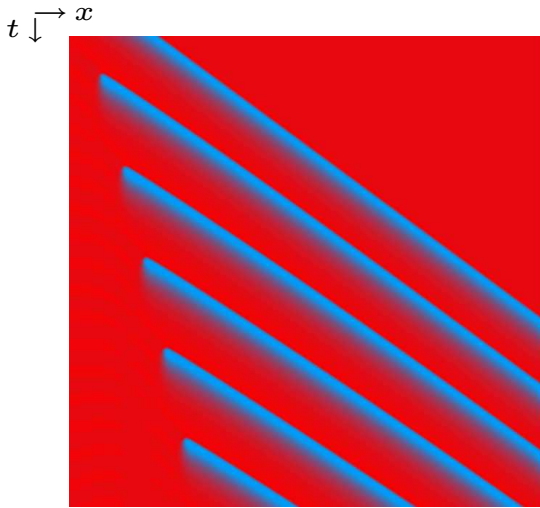
MB model



2-state / 3-state Calcium and IP_3 -receptor [Sneyd et al '00]

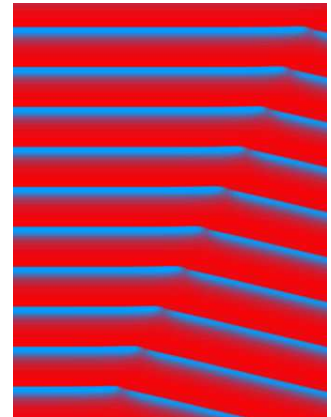
CICR [Snita et al '00]

Tracefiring in the Oregonator model



Two characteristic speeds:

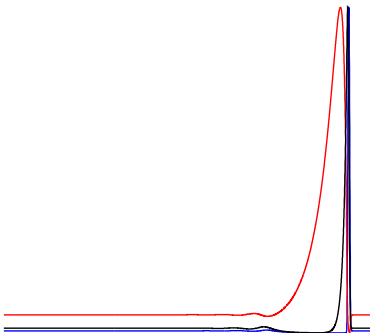
- **pulse-chain** ~ 3.4 (**sink**)
 ω -limit: invading pulse-chain
- **defect** ~ 0.6 (**source**)
 ω -limit: rest state connected to space-time periodic



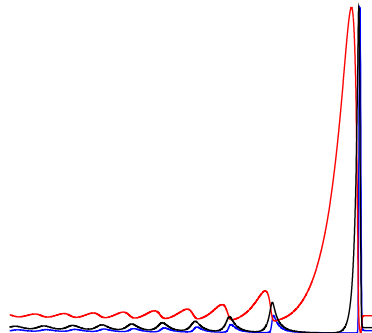
other defect for
 $\phi = 0.00016$

Kinetics parameter $\phi \sim$ light intensity \sim excitability
(b.c. in comoving frame: left Neumann, right Dirichlet)

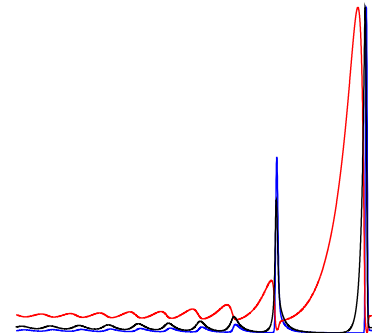
$\phi = 0.0006$



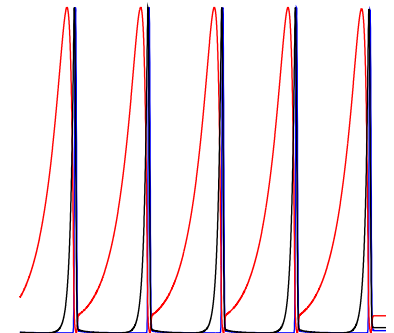
$\phi = 0.000155$



$t = 0$



$t = 50$



$t = 1000$

Spatial Dynamics

Constituents of pulse splitting, (back-) tracefiring

$$U \in \mathbb{R}^N, \quad U_t = DU_{xx} + F(U; \phi)$$

$$\xi = x - ct \quad \rightarrow \quad U_t = DU_{\xi\xi} + cU_\xi + F(U; \phi)$$


equilibria in these coordinates satisfy travelling wave ODE


$$0 = DU_{\xi\xi} + cU_\xi + F(U; \mu) \quad \Leftrightarrow$$

$$u_\xi = f(u; c, \phi) = f(u; \mu), \quad u \in \mathbb{R}^{2N}$$

bounded ODE solution \leftrightarrow **travelling wave, speed c**

homoclinic  \leftrightarrow **pulse** 

n -homoclinic \leftrightarrow **n -pulse** 

heteroclinic: pulse-train to equilibrium \leftrightarrow **invading pulse-chain** 

T-point in other models

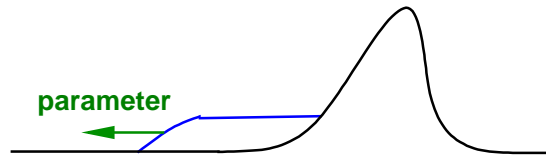
**Organizing Center: Constituents exist nearby,
pulse loses stability nearby**

MB [Zimmermann et al '97], Ca-flow [Sneyd et al '00] models:

Codim-2 heteroclinic cycle between two equilibria

(T-point [Glendinning, Sparrow '86])

**Interaction of pulse and equilibrium: plateau in pulse's wake
from other, (absolutely) unstable rest state**



Instability: MB [Sandstede, Scheel '00], Ca [Romero, Jones '03]

What happens in the Oregonator ?

The light-sensitive Oregonator model

$$u_t = D_u u_{xx} + (u(1 - u) - v(u - q))/\epsilon$$

$$v_t = D_v v_{xx} + (fw + \phi - v(u + q))/\delta$$

$$w_t = u - w$$

$$D_v = 1.12, \epsilon = 0.09, \delta = \epsilon/8, q = 0.001, f = 1.5$$

Chemical species in \mathbb{R}^3 : $u \sim$ bromous acid, $v \sim$ bromide, $w \sim$ catalyst

Travelling wave ODE in \mathbb{R}^5 ($c \approx 3.4$):

$$u_\xi = \tilde{u}$$

$$\tilde{u}_\xi = -cu_\xi - (u(1 - u) - v(u - q))/\epsilon$$

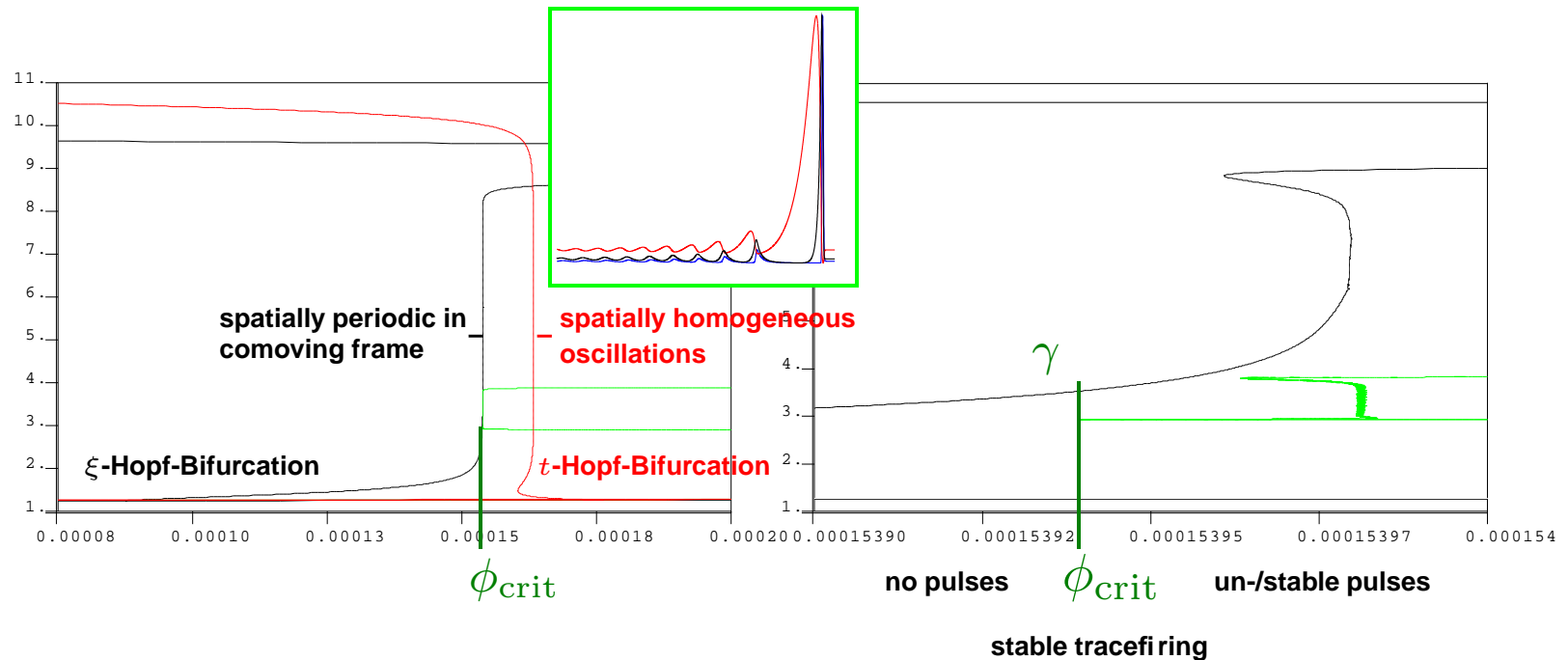
$$v_\xi = \tilde{v}$$

$$\tilde{v}_\xi = -(cv_\xi + (fw + \phi - v(u + q))/\delta)/D_v$$

$$w_\xi = (w - u)/c$$

Bifurcations and Tracefiring Onset

Hopf-bifurcations of kinetics ODE and travelling wave ODE, and path of homoclinic - instability when *decreasing* ϕ :

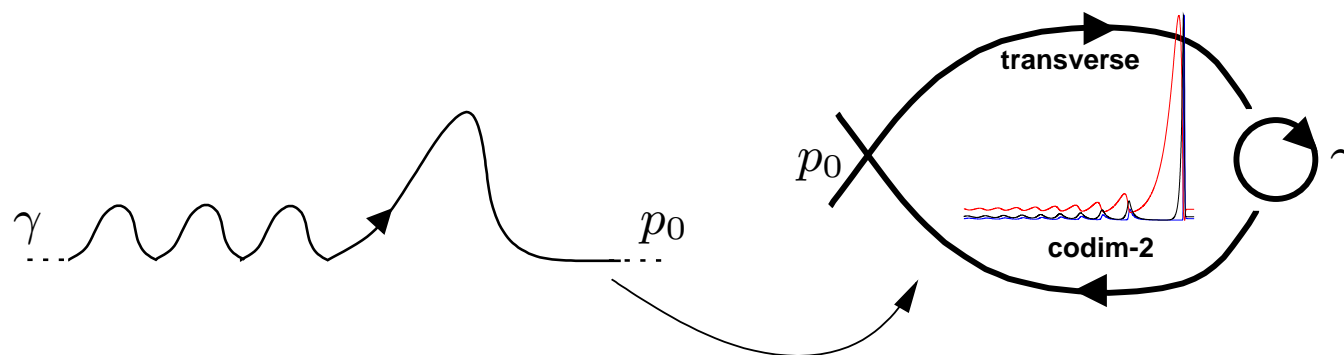


Periodic or Neumann boundary conditions: transient tracefiring

Dirichlet b.c. ahead of primary pulse: permanent tracefiring

Organizing Center for Oregonator's Tracefiring

Idea: codim-2 heteroclinic cycle with *periodic* orbit



Oregonator Morse indices:

$$i(p_0) = \dim(W^u(p_0)) = 3, \quad i(\gamma) := \dim(W^{cu}(\gamma)) = 2$$

Generally for ODE in \mathbb{R}^n :

periodic orbit: $i(\gamma) \geq 2, \dim(W^{sc}(\gamma)) = n - i(\gamma) + 1 \geq 2$

codim-2: $i(p_0) = i(\gamma) + 1 \rightarrow i(p_0) \geq 3$

get 2D transversality: $i(p_0) + \dim(W^{sc}(\gamma)) = n + 2$

for heteroclinic cycle: $\dim(W^s(p_0)) \geq 1 \rightarrow n \geq 4$

Such cycles have recently been found in other models [Sieber '02, Sneyd '03].

Existence of Constituents near Organizing Center

Theorem (R.)

Assume such a heteroclinic cycle with maximal transversality. Then a family of curves of locally unique 1-pulses exist, emanating from heteroclinic cycle in parameter space.

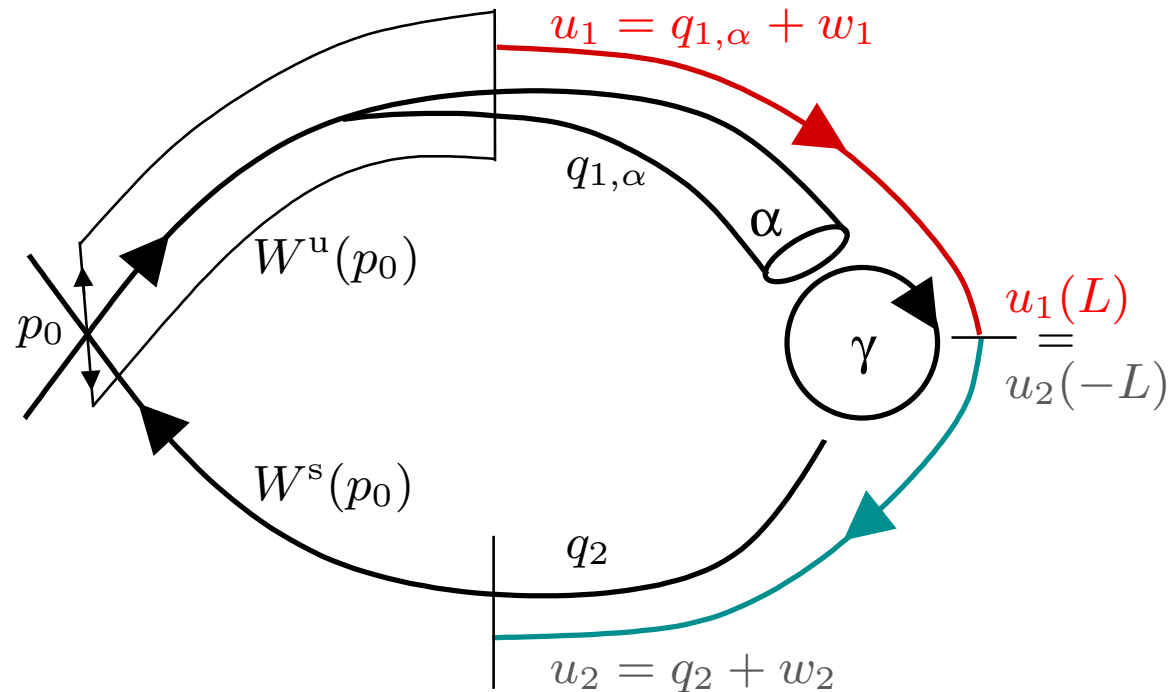
(Also for other codim and more general heteroclinic nets.)

Assume in addition $W_1 := W^u(p_0) \cap W^{cs}(\gamma)$ contains a curve homotopic to γ in $W_1 \cup \gamma$ (holds e.g. near spatial Hopf Bifurcation). Then nearby a smooth curve of unique 1-pulses exist, which emanates from the heteroclinic cycle.

Under Shil'nikov condition a (family of) curve(s) of n -pulses, and invading pulse-chain exists. (In Oregonator pulses are Shil'nikov-type and γ close to spatial Hopf bifurcation).

Idea of Proof: Adapt Lin's Method

1. Seek solutions near the heteroclinics and 'glued' close to γ
2. Match these away from γ with un/stable manifolds from p_0



Problems here:

- by phase shift: lack of hyperbolicity (essential spectrum)
- periodic distance of forward vs. backward approach to γ

Spectra of Pulses on Large Bounded Domains

and approaching heteroclinic bifurcation

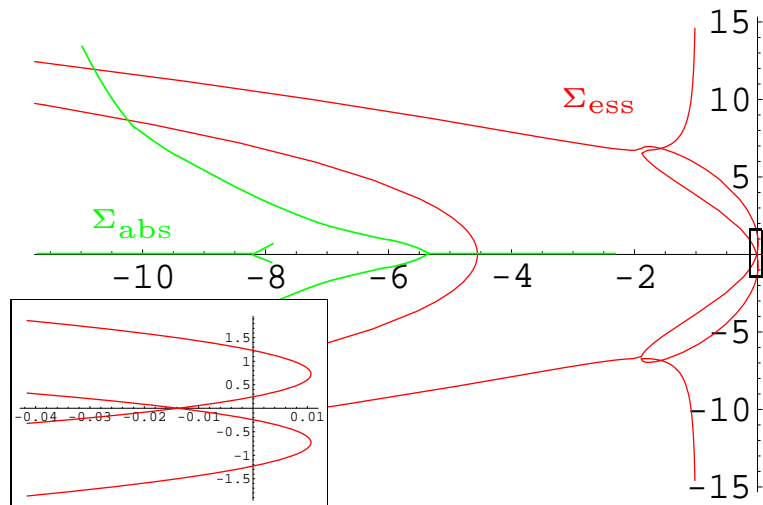
Theorems [Sandstede, Scheel '00] Typically:

- point spectrum from unbounded domain persists
- point spectrum for increasing domain length clusters:
near essential spectrum Σ_{ess} for *periodic* b.c.
near absolute spectrum Σ_{abs} for *separated* b.c.
- If stable pulse approaches heteroclinic orbit with absolutely unstable equilibrium (or periodic orbit) p_2 , then its critical point spectrum clusters near absolute spectrum of p_2 .

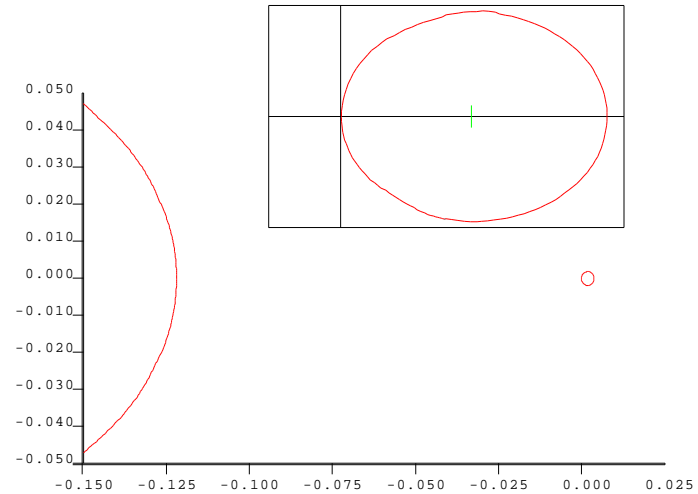
Σ_{abs} : **Solutions** $e^{\lambda t + \nu_1 \xi}$, $e^{\lambda t + \nu_2 \xi}$ **with** $\text{Re}(\nu_1) = \text{Re}(\nu_2)$ **and index condition. Distinguishes absolute and convective instabilities.**

Spectra in the Oregonator

Spectra of p_0 (base state)



Spectra of γ for ϕ near onset



Observations

- $\Sigma_{\text{ess}}(p_0)$ is already unstable: convective instability of pulse on unbounded domain or Neumann b.c. – *not* onset of tracefiring
- $\Sigma_{\text{abs}}(p_0)$ remains stable: no absolute instability of background
- $\Sigma_{\text{ess}}(\gamma)$, and in particular $\Sigma_{\text{abs}}(\gamma)$ unstable

Framework for Oregonator's Tracefiring Instability

Conclusions in light of organizing center idea

Assuming stable isolated point spectrum, the pulse becomes unstable when approaching the organizing center with absolutely unstable γ . In fact, winding around γ increases Morse-index. Similar for n -pulses.

Geometric indication of real-type absolute instability:

Theorem (R.)

All assumptions of theorem 1. Path of homoclinic spirals (logarithmically) into point of heteroclinic cycle, if and only if $0 \in \Sigma_{\text{abs}}$ and Σ_{abs} unstable.

Idea of proof

Use Floquet form of variation about γ and to compute leading order expansion of homoclinic path. This oscillates, if and only if $0 \in \Sigma_{\text{abs}}$ and Σ_{abs} unstable.

Discussion

- **several models exhibit tracefiring (or backfiring)**
- **Oregonator's transient tracefiring permanent with mixed b.c.**
- **idea for organizing center of Oregonator's tracefiring instability: theory and observations are coherent, and analogous to T-point:**
 - **constituents exist in unfolding**
 - **instability caused by abs. unstable 'wiggles' in pulse's wake**
 - **spiraling path (saddle-nodes) iff real-type absolute instability**

Perspectives

- **apply analytic method to bifurcations in other heteroclinic networks with periodic orbits**
- **periodic nature of tracefiring (also for T-points)**

Spectra and Stability

For a solution U_* to $U_t = DU_{\xi\xi} + cU_\xi + F(U; \phi)$, we say $\lambda \in \mathbb{C}$ is in the spectrum $\Sigma(U_*)$ of U_* , if

$$\mathcal{T}(\lambda) = (\lambda - D\partial_{\xi\xi} - c\partial_\xi - \partial_U F(U_*))$$

is not boundedly invertible.

We say $\lambda \in \Sigma(U_*)$ is in the point spectrum, if $\mathcal{T}(\lambda)$ is Fredholm with index 0; otherwise λ lies in the essential spectrum $\Sigma_{\text{ess}}(U_*)$.

Travelling waves spectra can be characterized in terms of exponential dichotomies of

$$v_\xi = \left(\partial_u f(u_*(\xi); \mu) + \lambda \begin{pmatrix} 0 & 0 \\ -\text{Id} & 0 \end{pmatrix} \right) v$$

Here spectral stability ($\text{Re}(\Sigma(U_*)) < 0$) implies nonlinear stability.