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A Comparison of the Rate of Convergence, Efficiency and Condition Number of Chebyshev and Legendre Polynomial Series with Prolate Spheroidal, Kosloff/Tal-Ezer and Theta-Mapped Fourier Basis Sets

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Pseudospectral Methods: Good & Bad

Good:

- Geometric Converge: $E(N) \propto \exp(-qN)$
[“Infinite Order”, “Exponential”, “Spectral”]
- With domain decomposition, (“spectral elements”), parallelizes well

Bad:

- “Stiffness”: CFL limit is $O(1/N^2)$ vs. $O(1/N)$ for equispaced finite difference.
- Highly Non-Uniform Resolution:
Linear-Density-in-Interior,
Quadratic-Density-in-Boundary Layers
- Highly Non-Uniform Grid in each element.

THEMES:

- All five spectral basis sets here are
cosines-with-change-of-coordinate

$$u(f[t]) = \sum_{j=1}^{\infty} a_j \cos(jt[x])$$

only the mapping $t(x)$ is different.

- Non-Chebyshev mappings can
improve grid uniformity which implies
Much longer stable timestep

$$O\left(\sqrt{N}\right)$$

Better accuracy

[asymptotically $(\pi/2)$ per dimension]

- Multiple non-Chebyshev choices:
Kosloff/Tal-Ezer basis, prolate spheroidal,
theta-mapped cosines [NEW]

Mapped-Cosine Basis Functions

Ancient identity:

$$T_n(x) \equiv \cos(nt[x]), \quad t(x) \equiv \arccos(x) \quad (1)$$

Legendre polynomials are mapped cosines too:

$$P_n(x) \sim \{\text{sign}(x)\}^n \frac{\sqrt{\arccos(|x|)}}{(1-x^2)^{1/4}} \times J_0 \left(\left[n + \frac{1}{2} \right] \arccos(|x|) \right) + O \left(\frac{0.062}{n^{3/2}} \right) \quad (2)$$

Except near $x = \pm 1$, this simplifies to

$$P_n(x) \sim \frac{\{\text{sign}(x)\}^n}{(1-x^2)^{-1/4}} \sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{n+1/2}} \times \cos \left(\left[n + \frac{1}{2} \right] \arccos(|x|) - \frac{\pi}{4} \right)$$

Query: Is the Mapping $t = \arccos(x)$
Optimum?

Three Mapped-Cosine Species That are Better
Than Chebyshev/Legendre:

1. **Kosloff/Tal-Ezer Basis**
2. **Prolate Spheroidal Functions**
3. **Theta-Mapped Cosines**

Advantages

- Nearly-Uniform Grid Improves Resolution by $\pi/2 \approx 1.57$ *per dimension*
- Timestep Lengthened by Order-of-Magnitude

Limits on Mapped Cosine Functions

Theorem 1 (Mapping Constraints) *Let*

$$u(f[t]) = \sum_{j=1}^{\infty} a_j \cos(jt[x]) \quad (3)$$

Infinite order convergence requires

- 1. All odd derivatives of $f(\tau)$ are zero at both $\tau = 0$ and $\tau = \pi$*
- 2. $f(\tau)$ is symmetric with respect to both $\tau = 0$ and $\tau = \pi$*
- 3. $f(\tau)$ is periodic with period 2π .*
- 4. The inverse function, $\tau = f^{-1}(x)$, has branch points at $x = \pm 1$; if $d^2f/d\tau^2 \neq 0$ at $\tau = 0, \pi$, then the branch points are square roots.*

Implications:

1. Grid cannot be completely uniform.
2. dt/dx must rise to vertical at $x = \pm\pi$.

Quasi-uniform, better-than-Chebyshev grid

IS POSSIBLE

Resolution of Mapped Cosines

- Evenly-spaced t -grid \Rightarrow non-uniform x -grid.
- Larger $dt/dx \Rightarrow$ smaller δx
- Higher minimum resolution (by $\pi/2$) than Chebyshev.

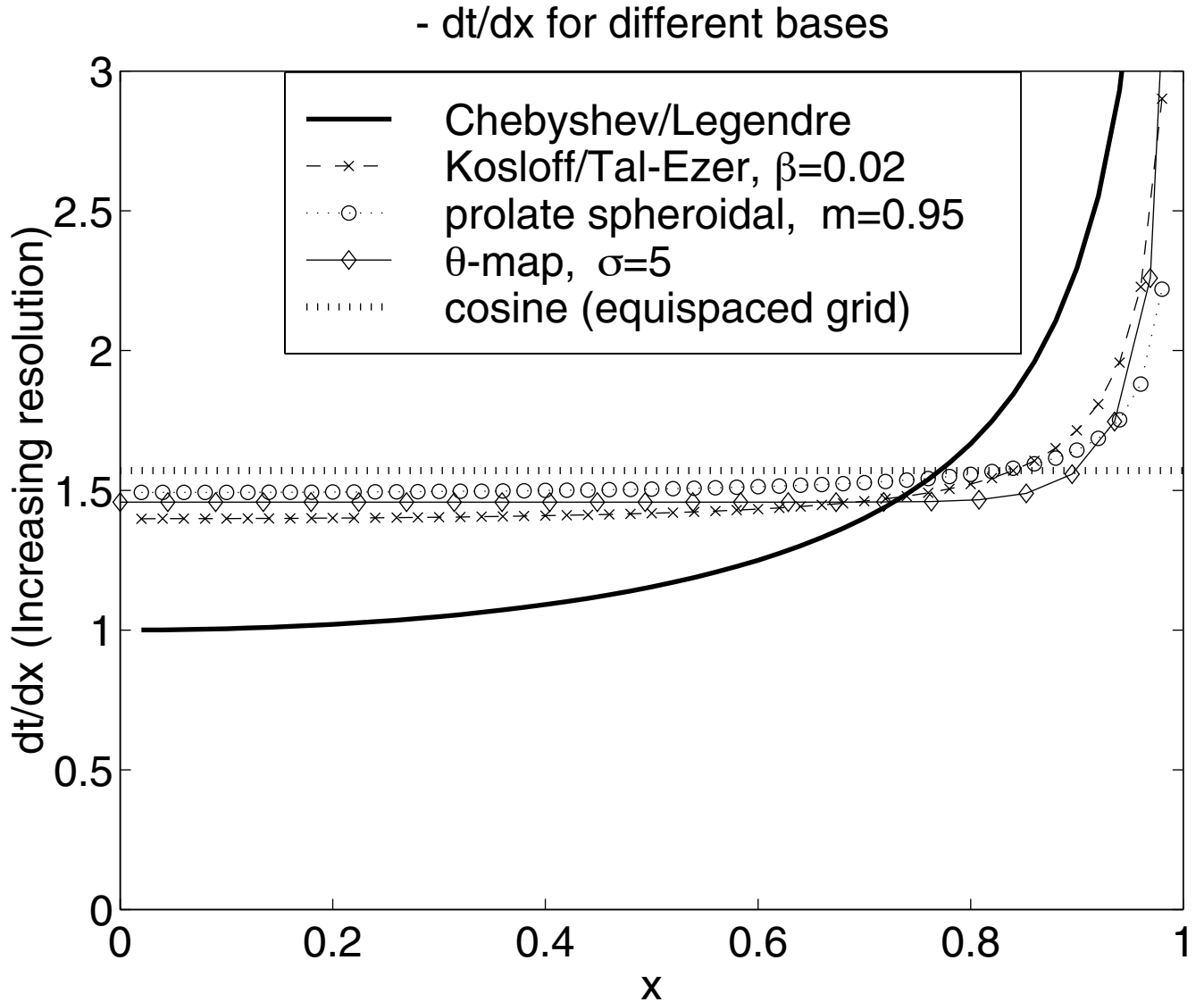


Figure 1: The slopes, $d\tau/dx$, for six different basis sets.

Kosloff/Tal-Ezer Basis

$$\phi_n^{KT}(x; \beta) \equiv \cos(n t^{KT}(x))$$

$$t^{KT} = \arccos \left\{ \frac{\sin \arcsin(1 - \beta)x}{1 - \beta} \right\} \quad (4)$$

$$\frac{dt^{KT}}{dx} \approx -\frac{\pi}{2} \quad \forall |x| < 1 - O(\sqrt{\beta}), \quad \beta \ll 1$$

i. e., maximum gridpoint separation ($\pi/2$) smaller than Chebyshev

Table 1: Theory and Applications of Kosloff/Tal-Ezer Mapping

| References | Comments |
|--|---|
| Kosloff&Tal-Ezer (1993) | Introduction and numerical experiments |
| Tal-Ezer(1994) | Theory; optimization of map parameter |
| Carcione(1994a) | Compares standard Chebyshev grid with Kosloff/Tal-Ezer grid |
| Renaut&Frohlich(1996) | 2D wave equations, one-way wave equation at boundary |
| Carcione(1996) | wave problems |
| Godon(1997b) | Chebyshev-Fourier polar coordinate model, stellar accretion disk |
| Renaut(1997) | Wave equations with absorbing boundaries |
| Don&Solomonoff(1997) | Accuracy enhancement and timestep improvement, especially for higher derivatives |
| Renaut&Su(1997) | 3rd order PDE; mapping was not as efficient as standard grid for $N < 16$ |
| Don&Gottlieb(1998) | Shock waves, reactive flow |
| Mead&Renaut(1999) | Analysis of Runge-Kutta time-integration |
| Hesthaven, Dinesen & Lynov(1999) | Diffraction optical elements; chose $\beta = 1 - \cos(1/2)$ to double timestep versus standard grid |
| Abril-Raymundo & García-Archilla(2000) | Theory and experiment for convergence of the mapping |
| Zhan&Ng | cardial modelling in 2D |

Kosloff/Tal-Ezer Basis: Virtues & Vices

Virtues

If $\beta \sim \text{constant}/N^2$:

1. Nearly-uniform grid; $\pi/2$ better than Chebyshev/Legendre
2. CFL limit $O(1/N)$, same as finite difference.

Vices

1. Mapping is *singular*;
branch point moves to $x \in [-1, 1]$ as $\beta \rightarrow 0$
2. $\beta \sim O(1/N^2)$ destroys spectral accuracy

Kosloff/Tal-Ezer: Error in Series of $u(x) \equiv x$

- Geometric convergence is saved only if β is *independent* of N (bottom curve)
- CFL limit is still $O(1/N^2)$, but may gain a factor of two.

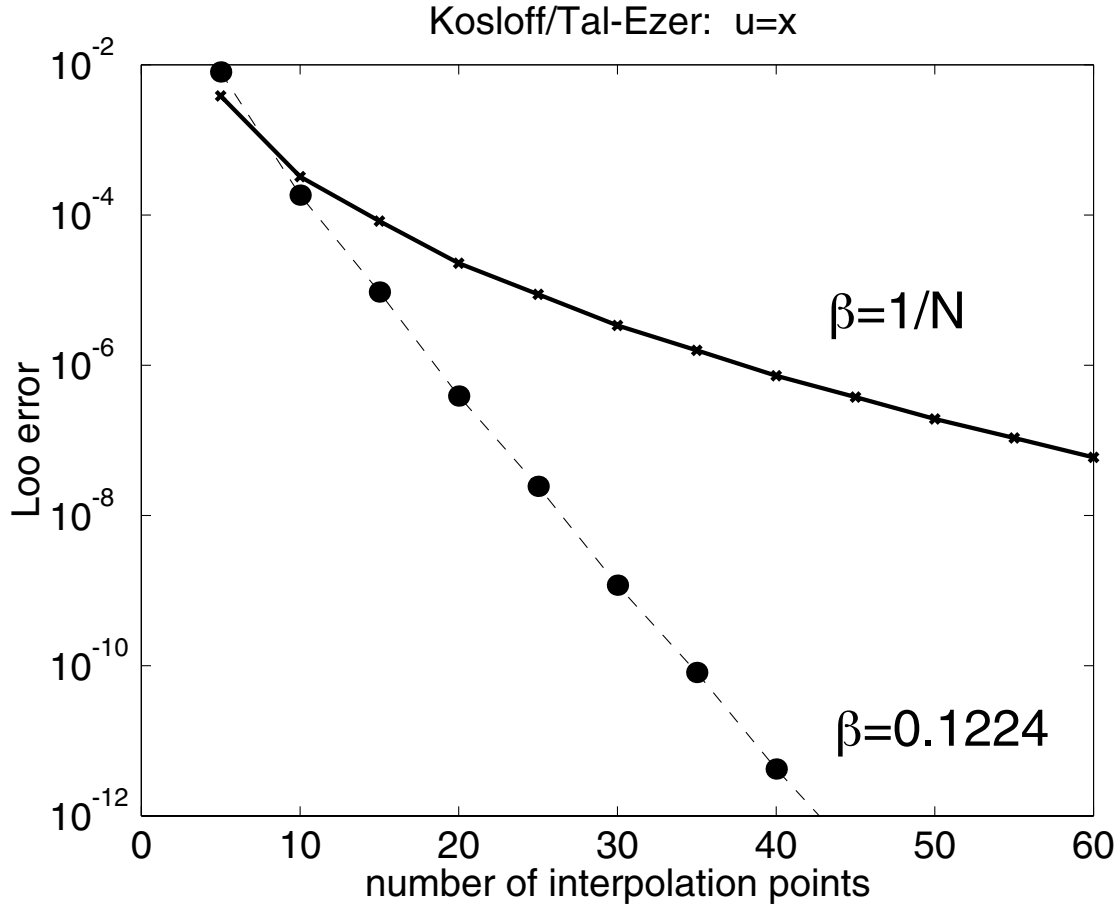


Figure 2: The error in the approximation of the linear function, $u(x) \equiv x$, by a Fourier cosine series using the Kosloff/Tal-Ezer mapping. When $\beta = 0.1224$, the choice of Hesthaven, Dinesen and Lynov(1999), the timestep limit is increased only by a factor of two, but the approximation still converges geometrically. When $\beta = 1/N$, the timestep limit for a first order hyperbolic problem shrinks only as $O(1/N^{3/2})$ versus the more severe Chebyshev/Legendre/Hesthaven *et al.* limit of $O(1/N^2)$. However, the usual rate of geometric convergence with N has been slowed to a subgeometric rate (upper curve) with an error falling as $\exp(-\text{constant}N^{1/2})$.

Prolate Spheroidal Functions of Zeroth Order

“Prolate Spheroidal Wave Functions are likely to be a better tool for the design of spectral and pseudo-spectral techniques than the orthogonal polynomials and related functions”

— Xiao, Rokhlin & Yarvin(2001), pg. 837.

- Defined as solutions $\psi_n(x, c)$ of

$$(1 - x^2)\psi_{xx} - 2x\psi_x + \left\{ \chi - c^2x^2 \right\} \psi = 0$$

n is the mode number, χ_n is eigenvalue and c is a constant, the “bandwidth” parameter.

- $\psi_n(x; c = 0) = P_n(x)$
- Complete asymptotic expansions are messy
- Relevant asymptotic expansion is simple.

Prolate Spheroidal Functions

“Transition bandwidth parameter”

$$c_*(n) \equiv \frac{\pi}{2} (n + 1/2)$$

Prolate functions span $x \in [-1, 1]$ only if

$$c \leq c_*(n)$$

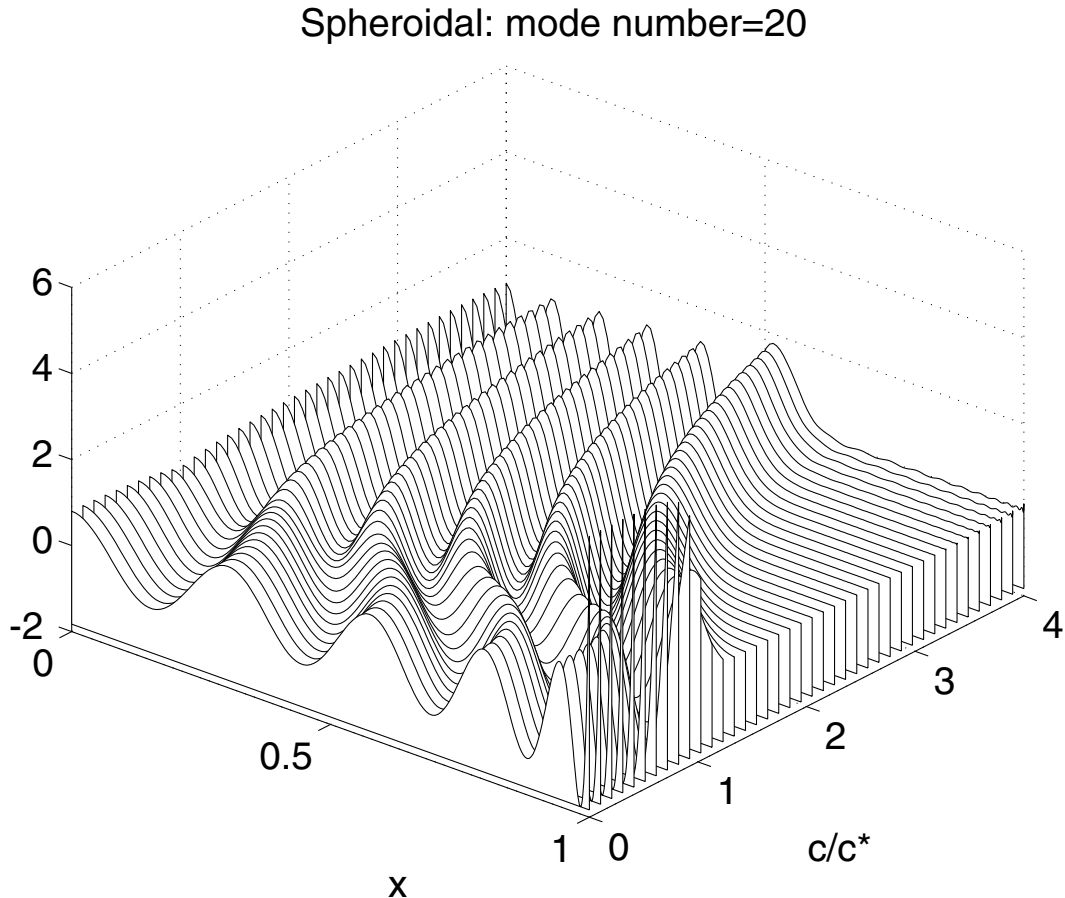


Figure 3: $\psi_{20}(x; c)$ in the $x - c$ plane.

Prolate Spheroidal Asymptotics

$$\psi_n(x; c) \sim \frac{\sqrt{\mathcal{E}(x; m)}}{(1-x^2)^{1/4} (1-mx^2)^{1/4}} \times J_0 \left(\frac{c}{\sqrt{m}} \mathcal{E}(x; m) \right)$$

$$\mathcal{E}(x; m) \equiv \int_x^1 dt \frac{\sqrt{1-mt^2}}{\sqrt{1-t^2}}; \quad m \equiv c^2/\chi_n$$

When $|x| < 1 - O(1/\sqrt{c})$,

$$\psi_n(x; c) \sim \sqrt{\frac{2}{\pi}} \frac{m^{1/4}}{\sqrt{\chi_n}} \frac{1}{(1-x^2)^{1/4} (1-mx^2)^{1/4}} \times \cos \left(\frac{c}{\sqrt{m}} \mathcal{E} - \pi/4 \right)$$

If $c = c_*(n)$, then

$$\psi_n(x; c = c_*(n)) \sim \sqrt{\frac{2}{\pi}} \frac{1}{c} \frac{1}{(1-x^2)^{1/2}} \times \cos \left(\frac{\pi}{2} n (1-x) \right) \quad (5)$$

Prolate Spheroidal Basis: Virtues & Vices

Virtues

1. Nearly-uniform grid; $\pi/2$ better than Chebyshev/Legendre
2. CFL limit $O(1/N^{3/2})$
3. Orthogonal with unit weight, like Legendre

Vices

1. Complicated to precompute function values & grid points
 - Symmetric tridiagonal Legendre-Galerkin
 - Newton-Raphson iteration for grid
2. Poorly-developed theory

Theta-Mapped Cosines

$$\begin{aligned}\phi_n^\theta(x; \sigma) &\equiv \cos(n t^\theta(x)) \\ t^\theta &= \Xi^{-1}(x; \sigma)\end{aligned}\tag{6}$$

$$\begin{aligned}\Xi(t; \sigma) &\equiv \sum_{m=-\infty}^{\infty} (-1)^m V(t - \pi m) \\ &\quad / \sum_{m=-\infty}^{\infty} (-1)^m V(\pi m)\end{aligned}$$

where

$$V = \frac{\pi}{2} t \operatorname{erf}(\sigma t) + \frac{\sqrt{\pi}}{2\sigma} \exp(-\sigma^2 t^2) \tag{7}$$

- Ξ is 2d integral of Jacobian θ -function.
- $\Xi(t, 0) \equiv \cos(t)$; basis $\Rightarrow T_n(x)$
- Unlike Kosloff/Tal-Ezer basis,
 θ -map is free of singularities

$$\frac{dt^\theta}{dx} \approx -\frac{\pi}{2} \quad \forall |x| < 1 - O(1/\sigma), \quad \sigma \gg 1$$

i. e., maximum grid point separation ($\pi/2$)
smaller than Chebyshev

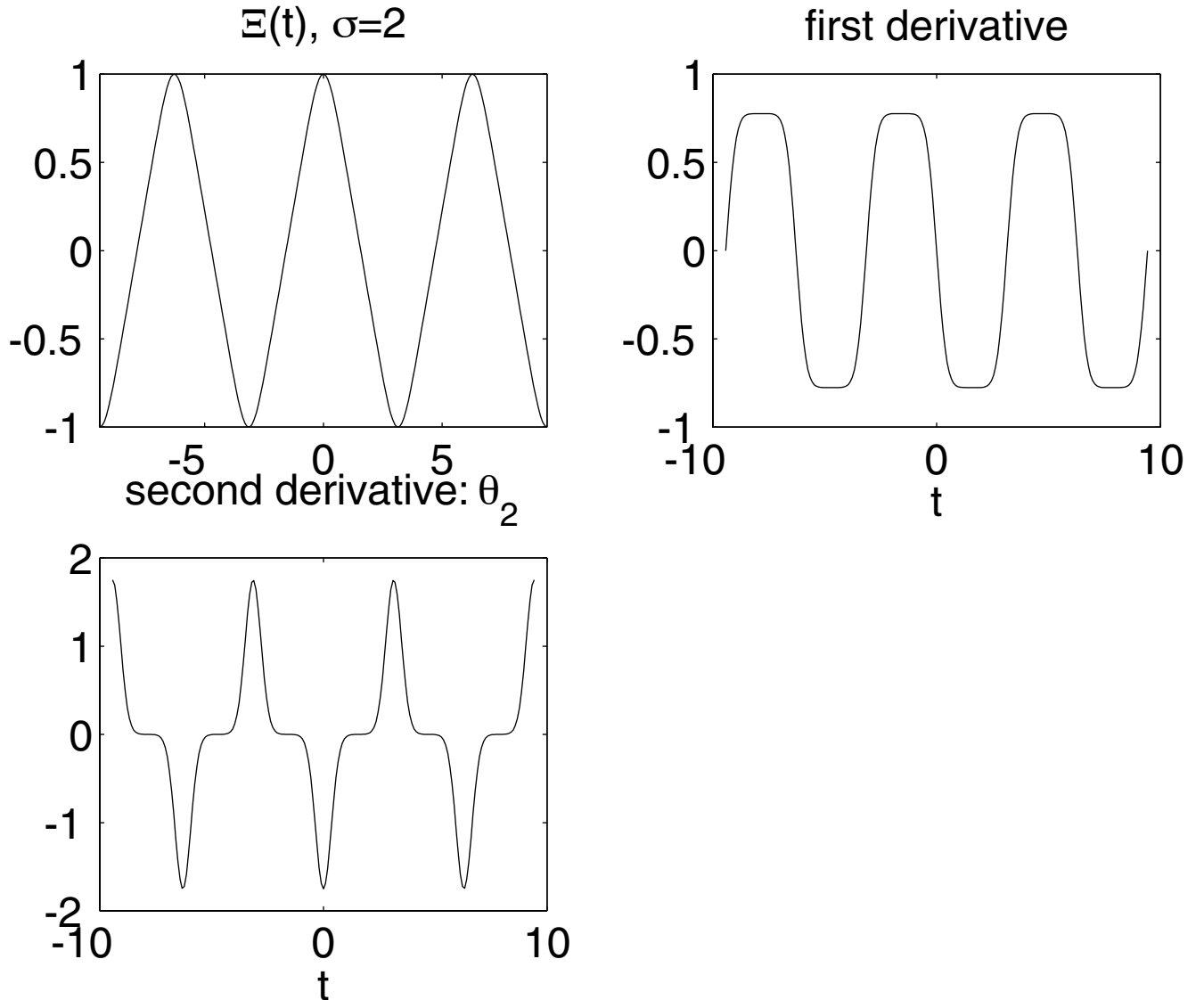


Figure 4: Illustrated over three periods to explicitly show the periodicity of the map.

Theta-Mapped Cosines

Minor disadvantage:

lack of explicit inverse $t(x)$

Easy to compute numerically by bisection

Table 2: Inverse of Theta-Map: $t(x)$

```
function t=finverse_thetamap(x,sigma);
itermax=50;      epsilon = 1.E-12;   t1=0;   t2=pi;
ff=Theta Map(t1,sigma) - x;
fmiddle=ThetaMap(t2,sigma) - x;
if ff < 0,   t=t1;   deltax=t2-t1;
else       t=t2;   deltax=t1-t2;           end \% if
for j=1:itermax
deltax=deltax*0.5; tmiddle=t + deltax;
fmiddle= ThetaMap(tmiddle,sigma) - x;
if(fmiddle <= 0), t=tmiddle; end \% if
if ((abs(deltax) < epsilon) | (fmiddle == 0)),
break; end \% if !
end \% j loop
```

Theta-Mapped Cosines: Comparisons with Chebyshev Polynomials

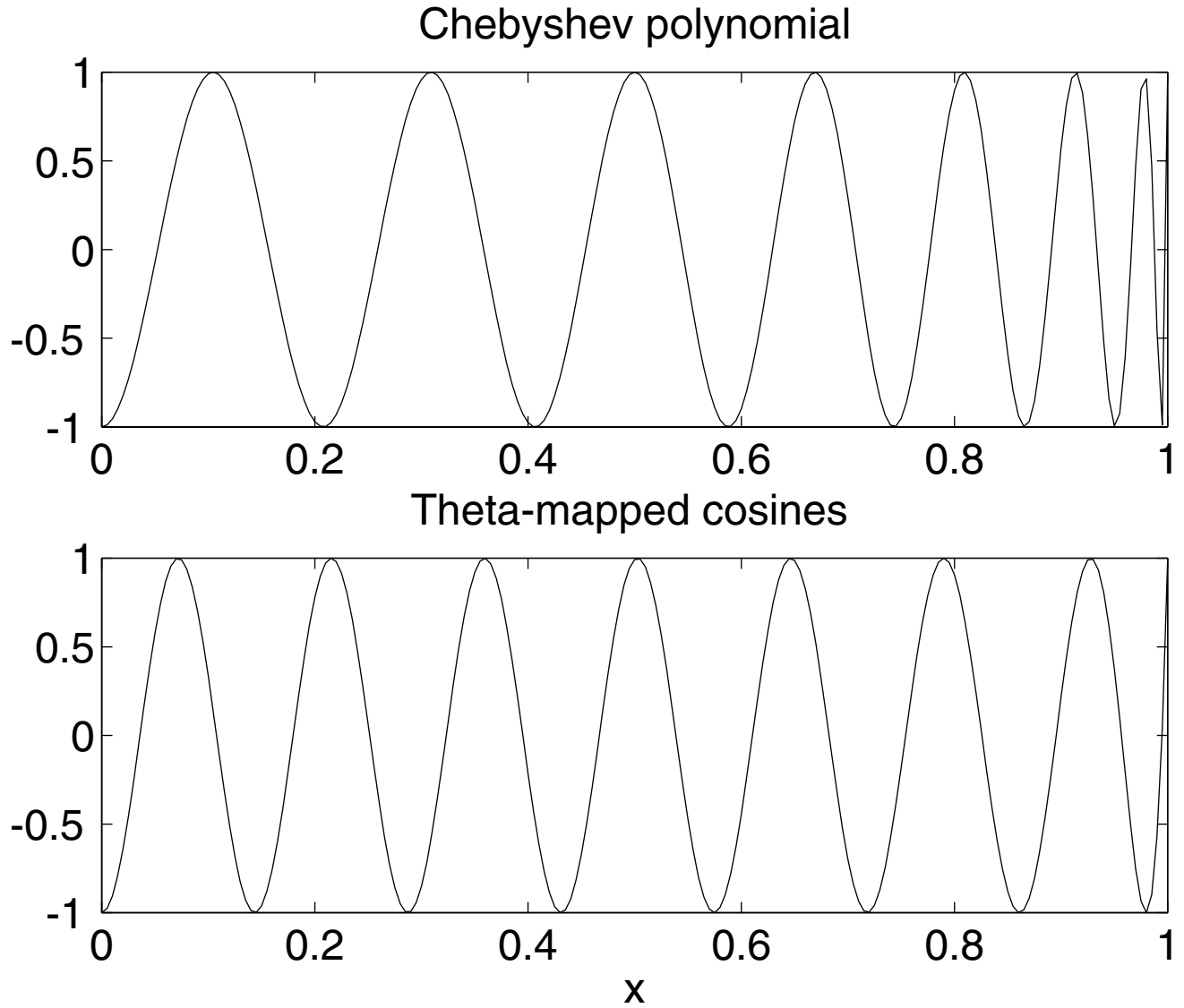


Figure 5: $N = 30$. For the θ -mapping, $\sigma = 5$

Theta-Mapped Cosines: Comparisons with Chebyshev Grid

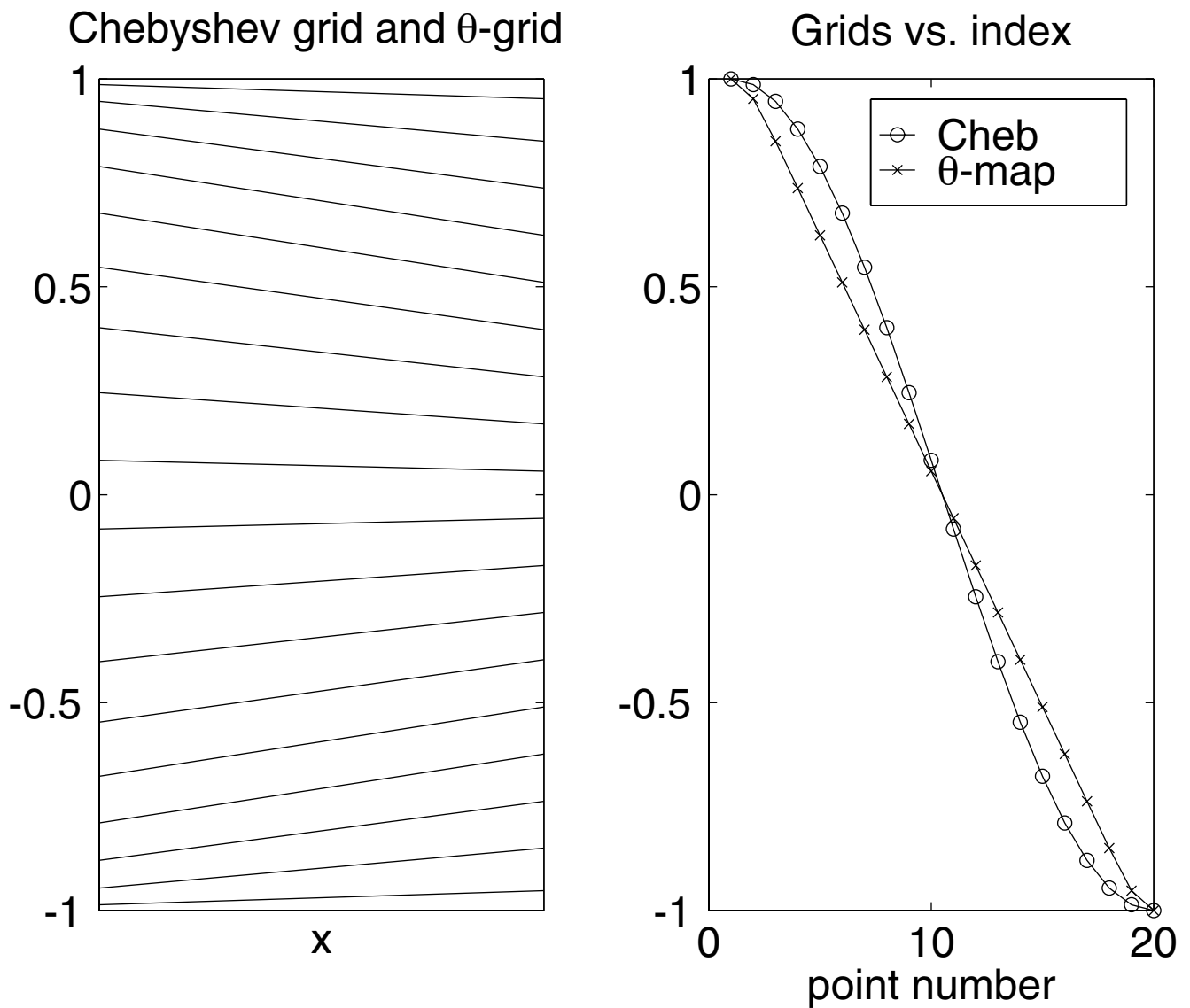


Figure 6: Left panel: Chebyshev grid on left, θ -grid on right.

Theta-Mapped Cosines: Error for $u(x) \equiv x$

- Geometric convergence requires $\sigma \sim O(\sqrt{N})$
- Minimum grid spacing is then $O(N^{3/2})$
- $\sigma \sim O(N) \Rightarrow$ uniform grid,
but exponential convergence is destroyed

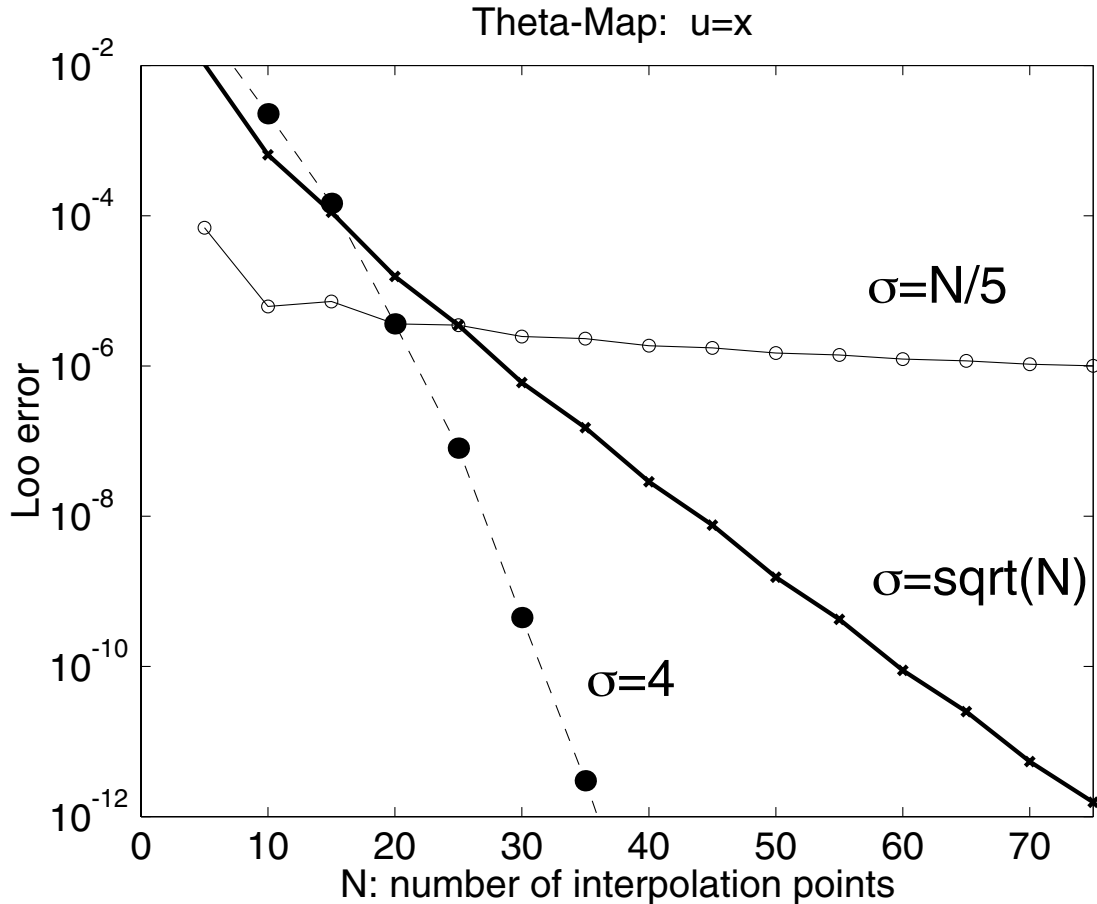


Figure 7:

θ -Mapped Cosines: Derivative Condition Numbers

- Chebyshev:

$$\max \left(\left| \text{eig} \left(\vec{\delta}_2 \right) \right| \right) \sim 0.045 N^4 \quad (8)$$

- θ -map with $\sigma = \sqrt{N}$:

$$\max \left(\left| \text{eig} \left(\vec{\delta}_2 \right) \right| \right) \sim 0.100 N^3 \quad (9)$$

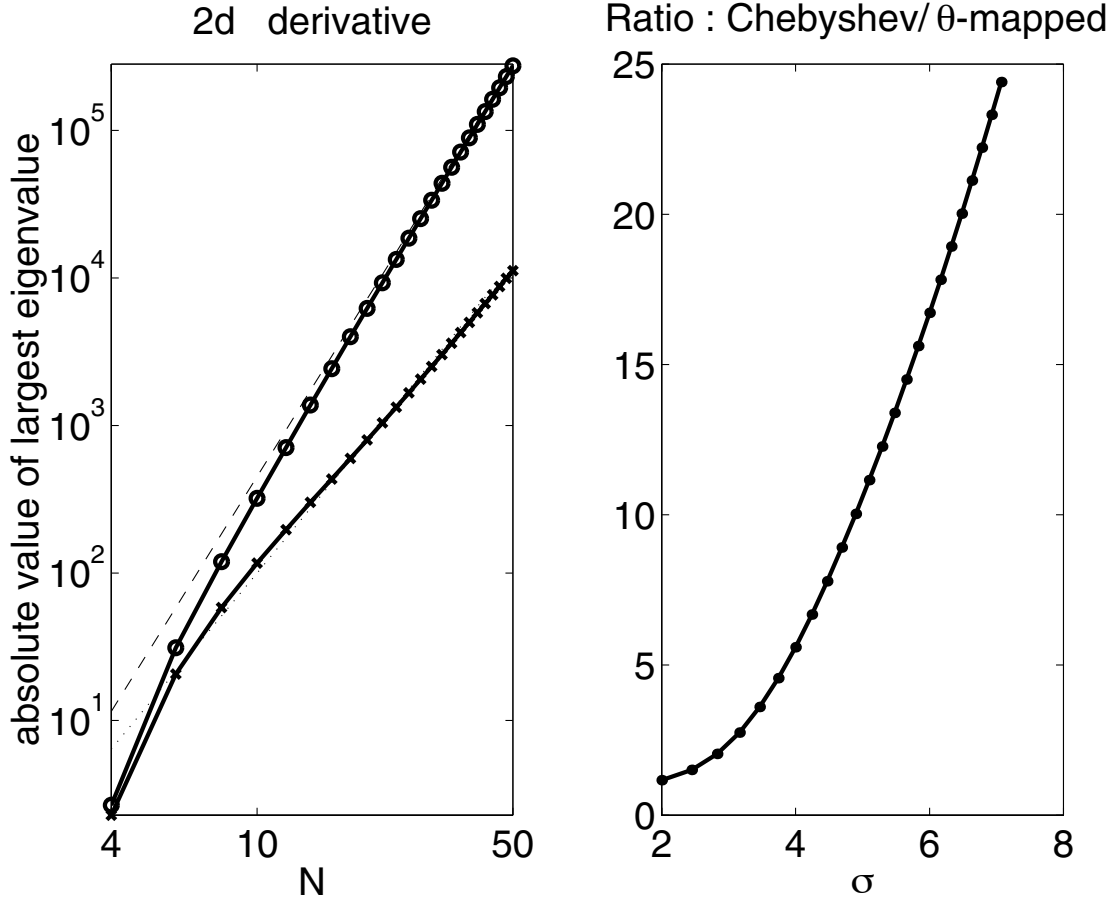


Figure 8: The guidelines (dashed) are the asymptotic forms.

Theta-Mapped Cosines: Error for $u(x) \equiv \cos(kx)$

- Chebyshev needs *minimum* of $N = k$ pts.

Map reduces this to $N = (2/\pi)k$
[asymptotically as $N \rightarrow \infty$]

- For well-resolved oscillations

Map dramatically reduces error

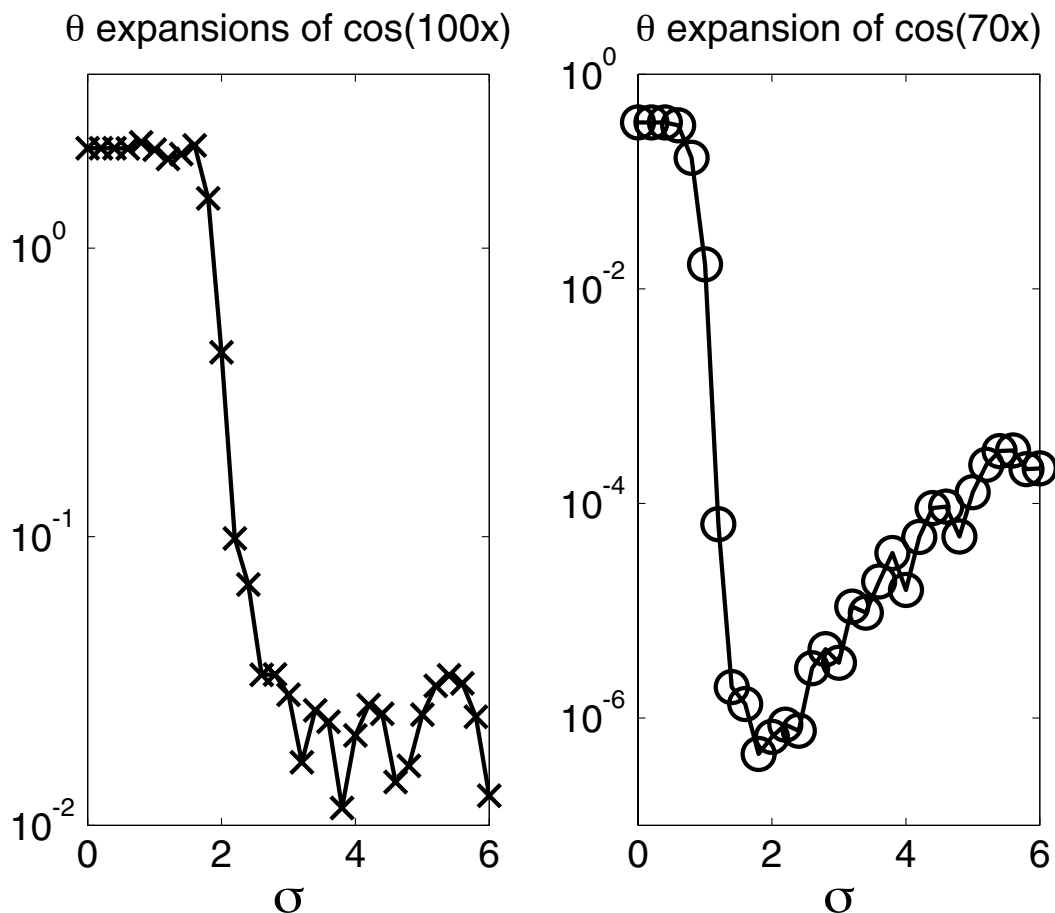


Figure 9: $N = 100$ point expansions of $\cos(100x)$ and $\cos(70x)$.

Theta-Mapped Cosines: Error for $u(x) \equiv \cos(kx)$

- $\sigma = 0$ [left axis] is Chebyshev polynomials
- Error grows monotonically with k for fixed σ .
- As $N \rightarrow \infty$, error contours bunch up
- As $N \rightarrow \infty$, contours asymptote to $\pi/2$ from below.

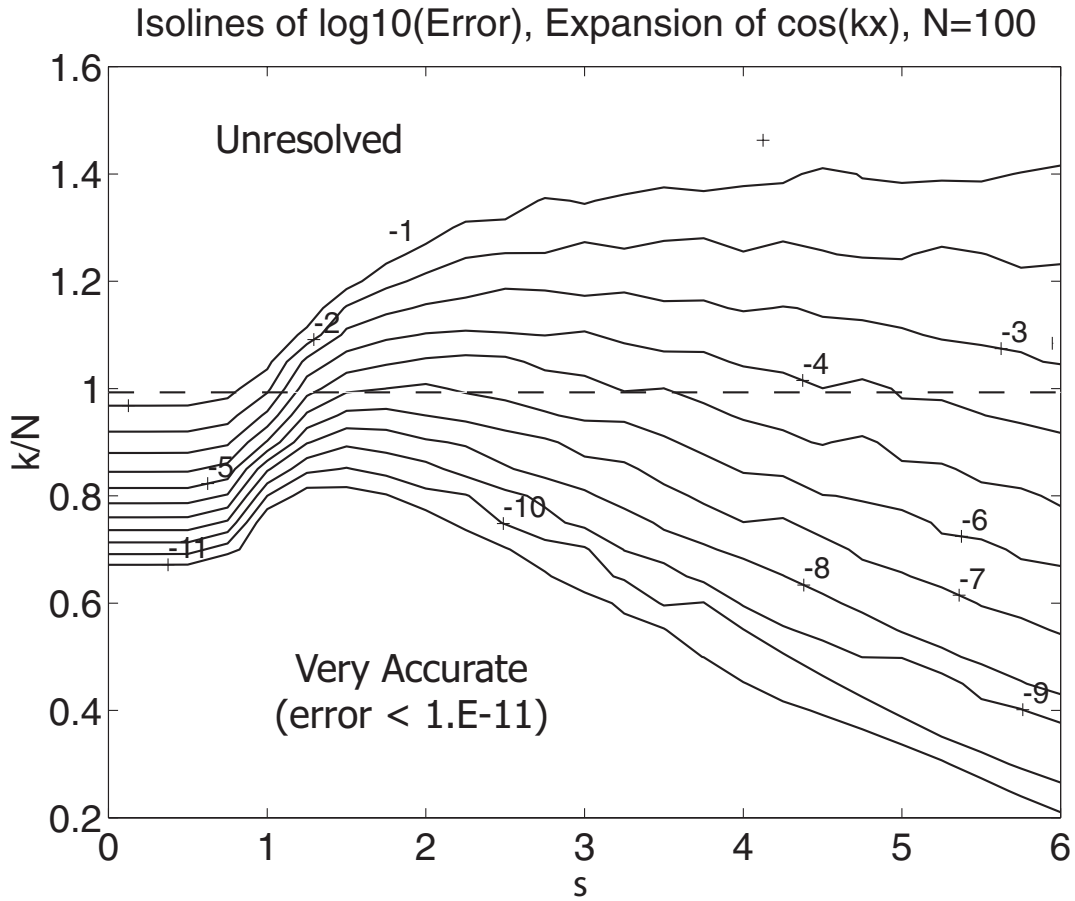


Figure 10: $N = 100$ expansions of $\cos(kx)$ for various k with the θ -map with various σ . Dashed line is the asymptotic Chebyshev limit for the highest resolved wavenumber.

Theta-Mapped Cosines: Error for $u(x) \equiv \cos(kx)$

For a given error, σ was chosen to push the error contour as high in k as possible; the maximum k for that error is then plotted.

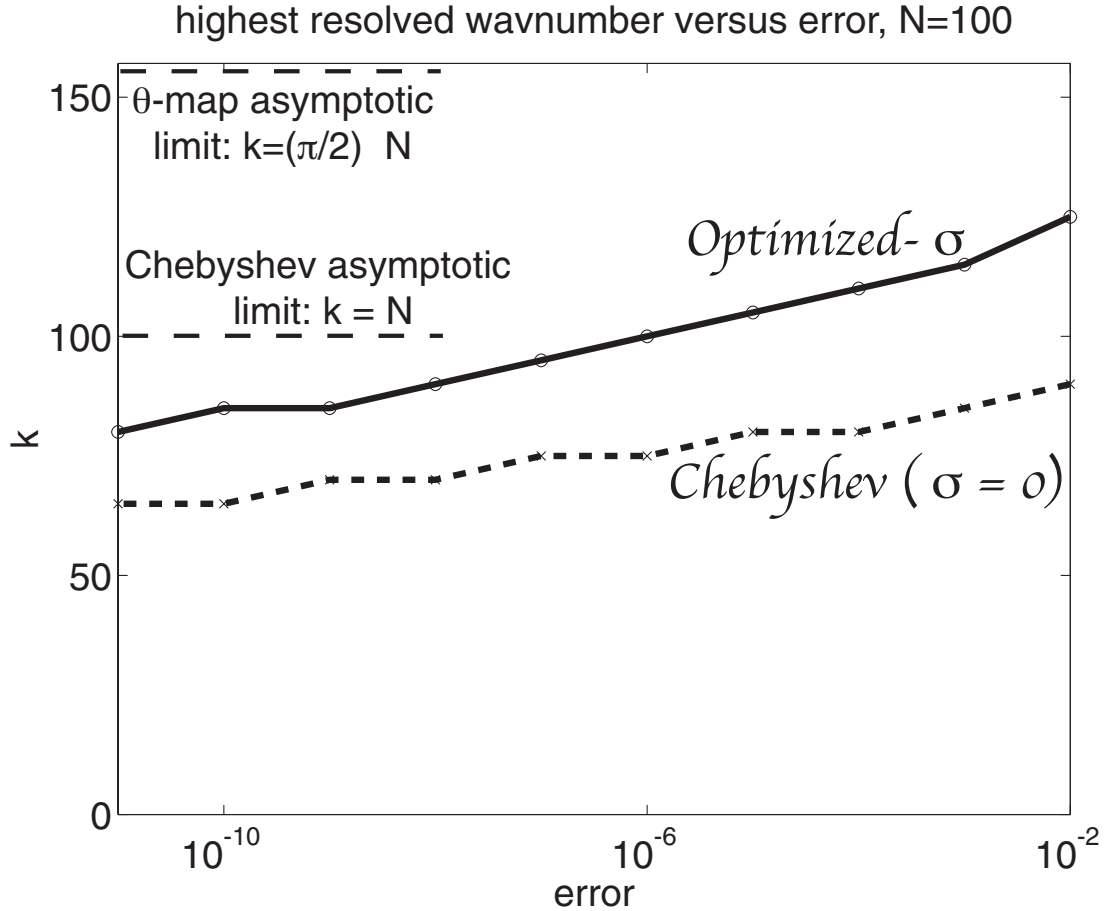


Figure 11: Number of resolved wavenumbers when σ has its optimum value, the value that pushes that particular error contour highest in k .

Optimizing Parameters

- No comprehensive theory as yet
- Theta-mapped theory harder than unbounded domain via steepest descent analysis
- Prolate theory *really* hard because
prolate \Rightarrow Legendre for *fixed* c ;
- Empirical, problem-dependent experimentation
is best strategy for now
- Experimentation cost-effective for
community models

Spectral Elements

- Prolate orthogonal with UNIT WEIGHT
 \Rightarrow trivial to replace Legendre
- Needs only grid, weights, derivatives-at-grid
- grid & weights: Newton-Raphson iteration
functions/derivatives: Symmetric
tridiagonal Legendre-Galerkin
- Theta-mapped cosines are *almost* orthogonal
easy to Gram-Schmidt orthogonalize
- Best basis for multi-domains?

inner product of basis ["Gram" or "mass" matrix], $\sigma=4$

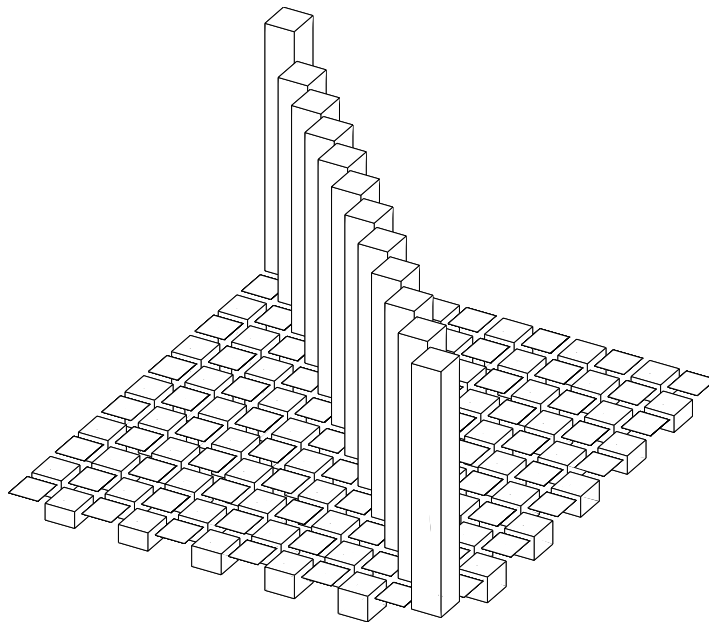


Figure 12:

Virtues of Cosine-Mapped Functions

1. Better resolution by $\pi/2$
2. More uniform grid
3. Much larger CFL timestep limit

Vices

1. All contain a free parameter
2. No theory for choosing parameter
3. More complicated to program than Chebyshev

Additional Conclusions

- Kosloff/Tal-Ezer inferior because of map singularity
- Prolate are orthogonal with unit weight
easy to drop into spectral elements
- θ -map is best for single-domain (simplest!)

Future Problems

- *theory for optimizing σ or c*
- Empirical guidelines for σ or c
- Practical experience with “prolate elements”, “theta elements”, etc., on hard problems