

6. L -functions in the constant terms.

Recall the global intertwining operator $M(s, \pi) = \otimes_v A(s, \pi_v, w_0)$. For $f \in V(s, \pi)$,

$$f \in \otimes_{v \in S} V(s, \pi_v) \otimes \otimes_{v \notin S} f_v^0,$$

where S is a finite set of places, including archimedean places, and f_v^0 is the spherical vector such that $f_v^0(k_v) = 1$ for $k_v \in \mathbf{G}(\mathcal{O}_v)$. Suppose $f = \otimes f_v$, where $f_v = f_v^0$ for $v \notin S$. Then

$$M(s, \pi)f = \otimes_{v \in S} A(s, \pi_v, w_0)f_v \otimes \otimes_{v \notin S} A(s, \pi_v, w_0)f_v^0.$$

Langlands observed that L -functions show up in the calculation of $A(s, \pi_v, w_0)f_v^0$. First we need inducing in stages of induced representations. Suppose $\mathbf{P}_1 = \mathbf{M}_1\mathbf{N}_1 \subset \mathbf{P}_2 = \mathbf{M}_2\mathbf{N}_2$ are two parabolic subgroups. Then $\mathbf{P}_1 \mapsto \mathbf{M}_2 \cap \mathbf{P}_1$ is a bijection from $\{\mathbf{P}_1 : \mathbf{P}_1 \subset \mathbf{P}_2\}$ onto the set of standard parabolic subgroups of \mathbf{M}_2 . Then $\mathbf{M}_2 \cap \mathbf{N}_1$ is the unipotent radical of $\mathbf{M}_2 \cap \mathbf{P}_1$; $\mathbf{M}_2 \cap \mathbf{P}_1 = \mathbf{M}_1 \cdot (\mathbf{M}_2 \cap \mathbf{N}_1)$ is the Levi decomposition.

Lemma 6.1 (inducing in stages). *Suppose $\mathbf{P}_1 \subset \mathbf{P}_2$. Let $\mathbf{Q} = \mathbf{M}_2 \cap \mathbf{P}_1$. Let π be a representation of $\mathbf{M}_2(F)$ such that $\pi = \text{Ind}_Q^{M_2} \sigma \otimes \exp(\langle \Lambda_0, H_Q(\cdot) \rangle)$. Then*

$$\text{Ind}_{P_2}^G \pi \otimes \exp(\langle \Lambda, H_{P_2}(\cdot) \rangle) = \text{Ind}_{P_1}^G \sigma \otimes \exp(\langle \Lambda + \Lambda_0, H_{P_1}(\cdot) \rangle),$$

where $\Lambda \in X^*(M_2) \otimes \mathbb{C}$ extends to $X^*(M_1) \otimes \mathbb{C}$.

Suppose π_v is spherical. Then $\pi_v \hookrightarrow I(\chi_v)$. By inducing in stages,

$$I(s, \pi_v) \subset \text{Ind}_B^G \chi_v \otimes \exp \langle s\tilde{\alpha}, H_B(\cdot) \rangle = I(s\tilde{\alpha}, \chi_v).$$

We denote by $A(s\tilde{\alpha}, \chi_v, w_0)$ the intertwining operator for $I(s\tilde{\alpha}, \chi_v)$. Then $A(s, \pi_v, w_0) = A(s\tilde{\alpha}, \chi_v, w_0)|_{I(s, \pi_v)}$. Here $f_v^0 \in I(s\tilde{\alpha}, \chi_v)$, and

$$f_v^0(tuk) = \chi(t) \exp \langle s\tilde{\alpha} + \rho_B, H_B(t) \rangle,$$

for $t \in \mathbf{T}(F_v)$, $u \in \mathbf{U}(F_v)$, $k \in \mathbf{G}(\mathcal{O}_v)$. We need to calculate

$$A(s\tilde{\alpha}, \chi_v, w_0)f_v^0(e) = \int_{\mathbf{N}(F_v)} f_v^0(w_0^{-1}n) dn.$$

(for simplicity, we assume that \mathbf{P} is self-conjugate.)

We reduce the calculation to SL_2 case by using cocycle relation of intertwining operators: Let $\Lambda \in \mathfrak{a}_{\mathbb{C}}^*$, χ a character of $\mathbf{T}(F)$. Let $A(\Lambda, \chi, w)$ be the intertwining operator from $I(\Lambda, \chi)$ to $I(w\Lambda, w\chi)$. Let $w = w_2w_1$ be a minimal length decomposition, i.e., $l(w) = l(w_1) + l(w_2)$. Then

Theorem 6.2. $A(\Lambda, \chi, w) = A(w_1 \Lambda, w_1 \chi, w_2) A(\Lambda, \chi, w)$.

Hence if $w = w_{\alpha_n} \cdots w_{\alpha_1}$, where $\alpha_1, \dots, \alpha_n \in \Delta$ and $l(w) = n$, then

$$A(\Lambda, \chi, w) = A(w_{\alpha_{n-1}} \cdots w_{\alpha_1} \Lambda, w_{\alpha_{n-1}} \cdots w_{\alpha_1} \chi, w_{\alpha_n}) \cdots A(\Lambda, \chi, w_{\alpha_1}).$$

Each $A(w_{\alpha_i} \cdots w_{\alpha_1} \Lambda, w_{\alpha_i} \cdots w_{\alpha_1} \chi, w_{\alpha_{i+1}})$ is an intertwining operator for SL_2 : Suppose α is a simple root. Recall the homomorphism $\phi_\alpha : SL_2 \rightarrow \mathbf{G}$ such that

$$\phi_\alpha \begin{pmatrix} 1 & u \\ 0 & 1 \end{pmatrix} = e_\alpha(u), \quad \phi_\alpha \begin{pmatrix} u & 0 \\ 0 & u^{-1} \end{pmatrix} = h_\alpha(u), \quad \phi_\alpha \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = w_\alpha.$$

Note that $\exp(\langle \Lambda, H_B(h_\alpha(u)) \rangle) = |u|^{\langle \Lambda, \alpha^\vee \rangle}$. Also $\langle w\Lambda, \alpha^\vee \rangle = \langle \Lambda, (w^{-1}\alpha)^\vee \rangle$ and $\exp(\langle w\Lambda, H_B(h_\alpha(u)) \rangle) = \exp(\langle \Lambda, H_B(h_{w^{-1}\alpha}(u)) \rangle)$.

If $w = w_2 w_1$, $l(w) = l(w_2) + l(w_1)$, then

$$\{\alpha > 0, w\alpha < 0\} = w_1^{-1}\{\beta > 0, w_2\beta < 0\} \cup \{\gamma > 0, w_1\gamma < 0\},$$

where it is a disjoint union. Hence by induction, if $w = w_{\alpha_n} \cdots w_{\alpha_1}$, where $\alpha_1, \dots, \alpha_n \in \Delta$, $l(w) = n$, then

$$\{\alpha > 0, w\alpha < 0\} = \{\alpha_1, w_{\alpha_1}(\alpha_2), w_{\alpha_1}w_{\alpha_2}(\alpha_3), \dots, w_{\alpha_1} \cdots w_{\alpha_{n-1}}(\alpha_n)\}.$$

Therefore, $\langle w_{\alpha_{i-1}} \cdots w_{\alpha_1} \Lambda, \alpha_i^\vee \rangle = \langle \Lambda, (w_{\alpha_1} \cdots w_{\alpha_{i-1}}(\alpha_i))^\vee \rangle$.

Now we compute the intertwining operator in the case of SL_2 . Let $G = SL_2(\mathbb{Q}_p)$ and χ_p is a character of \mathbb{Q}_p^\times . Let

$$A(s, \chi_p) f_p(g) = \int_{N_p} f_p(w_0^{-1} n g) dn,$$

be the intertwining operator for $I(s, \chi_p)$, where $N_p = \left\{ \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \right\}$, and $w_0 =$

$$\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}.$$

Theorem 6.3 (Computation in the case of $SL_2(\mathbb{Q}_p)$; Gindikin-Karpelevich formula). *Let χ_p be an unramified character. Then*

$$A(s, \chi_p) f_p^0(e) = \frac{1 - \chi_p(p) p^{-s-1}}{1 - \chi_p(p) p^{-s}} = \frac{L(s, \chi_p)}{L(s+1, \chi_p)},$$

where f_p^0 is the spherical function such that $f_p^0(e) = 1$. It is the unique $SL_2(\mathbb{Z}_p)$ -fixed function satisfying $f_p^0 \left(\begin{pmatrix} a & x \\ 0 & a^{-1} \end{pmatrix} g \right) = \chi_p(a) |a|^{s+1} f_p^0(g)$.

Proof. We compute $A(s, \chi_p) f_p^0(e) = \int_{\mathbb{Q}_p} f_p^0(w_0^{-1} n) dx$. The Haar measure dx satisfies; $\int_{\mathbb{Q}_p} f(xy) dy = |x|^{-1} \int_{\mathbb{Q}_p} f(y) dy$, and $\int_{\mathbb{Z}_p} dx = 1$. Hence $\int_{\mathbb{Z}_p^\times} dx = \int_{\mathbb{Z}_p} dx - \int_{p\mathbb{Z}_p} dx = 1 - \frac{1}{p} = \frac{p-1}{p}$. Also $\int_{p^{-m}\mathbb{Z}_p^\times} dx = p^m \int_{\mathbb{Z}_p^\times} dx$.

Any element of $x \in \mathbb{Q}_p$ can be written as $x = p^{-m}u$, $u \in \mathbb{Z}_p^\times$ (units in \mathbb{Z}_p), and $m \in \mathbb{Z}$. Note that for $n = \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}$,

$$w_0^{-1}n = \begin{pmatrix} 0 & -1 \\ 1 & x \end{pmatrix} = \begin{pmatrix} x^{-1} & 0 \\ 0 & x \end{pmatrix} \begin{pmatrix} 1 & -x \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ x^{-1} & 1 \end{pmatrix}.$$

Hence if $x \in \mathbb{Z}_p$, $f_p^0(w_0^{-1}n) = 1$. If $x = p^{-m}u$ with $m \geq 1$, then $f_p^0(w_0^{-1}n) = \chi_p(p^m)(p^{-m})^{s+1}$. Therefore,

$$\begin{aligned} \int_{\mathbb{Q}_p} f_p^0(w_0^{-1}n) dx &= \int_{\mathbb{Z}_p} f_p^0(w_0^{-1}n) dx + \sum_{m=1}^{\infty} \int_{p^{-m}\mathbb{Z}_p^\times} f_p^0(w_0^{-1}n) dx \\ &= 1 + \sum_{m=1}^{\infty} \chi_p(p^m)(p^{-m})^{s+1} p^m \left(1 - \frac{1}{p}\right) \\ &= 1 + \left(1 - \frac{1}{p}\right) \sum_{m=1}^{\infty} (\chi_p(p)p^{-s})^m = 1 + \left(1 - \frac{1}{p}\right) \frac{\chi_p(p)p^{-s}}{1 - \chi_p(p)p^{-s}} \\ &= \frac{1 - \chi_p(p)p^{-s-1}}{1 - \chi_p(p)p^{-s}}. \end{aligned}$$

More generally over an arbitrary algebraic number field, we can show

$$\int_{F_v} f_v^0(w_0^{-1}n) dn = \frac{L(s, \chi_v)}{L(s+1, \chi_v)}.$$

where $L(s, \chi_v) = (1 - \chi_v(\varpi_v)q_v^{-1})^{-1}$, ϖ_v is a uniformizer in F_v , and q_v is the number of elements in $\mathcal{O}_v/\mathfrak{p}_v$.

Corollary 6.4.

$$A(\Lambda, \chi_v, w) f_v^0(e) = \prod_{\beta > 0, w\beta < 0} \frac{L(\langle \Lambda, \beta^\vee \rangle, \chi_v \circ \beta^\vee)}{L(1 + \langle \Lambda, \beta^\vee \rangle, \chi_v \circ \beta^\vee)}.$$

If $F = \mathbb{R}$, we leave it as an exercise to show that

$$A(s, 1) f_\infty^0(e) = \sqrt{\pi} \frac{\Gamma(\frac{s}{2})}{\Gamma(\frac{s+1}{2})},$$

where f_∞^0 is the spherical function satisfying $f_\infty^0\left(\begin{pmatrix} a & x \\ 0 & a^{-1} \end{pmatrix} g\right) = |a|^{s+1} f_\infty^0(g)$.

Use the fact that

$$w_0^{-1}n = \begin{pmatrix} 0 & -1 \\ 1 & x \end{pmatrix} = \begin{pmatrix} \Delta_x^{-1} & -x\Delta_x^{-1} \\ 0 & \Delta_x \end{pmatrix} \kappa_{\theta(x)},$$

where $\Delta_x = \sqrt{1+x^2}$, $\theta(x) = \arctan(-\frac{1}{x})$, and $\kappa_\theta = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$. Hence

$$A(s, 1)f_\infty^0(e) = \int_{-\infty}^{\infty} (1+x^2)^{-\frac{s+1}{2}} dx.$$

Proof of Theorem 6.2. Let $N(w) = U \cap w\overline{U}w^{-1}$, where \overline{U} is the opposite unipotent radical. By the commutator relation, each element of $N(w)$ can be expressed in the form $\prod_{\alpha>0, w^{-1}\alpha<0} e_\alpha(u_\alpha)$, where $u_\alpha \in F$ (in any order).

Some properties of $h_\alpha, e_\alpha, w_\alpha$;

$$\begin{aligned} e_\alpha(t_1)e_\alpha(t_2) &= e_\alpha(t_1 + t_2) \\ [e_\alpha(t), e_\beta(u)] &= \begin{cases} 1, & \text{if } \alpha + \beta \notin \Phi \\ e_{\alpha+\beta}(c_{\alpha\beta}tu), & \text{if } \alpha + \beta \in \Phi \end{cases} \\ h_\alpha(t)h_\alpha(u) &= h_\alpha(tu) \\ h_\alpha(t)e_\beta(u)h_\alpha(t)^{-1} &= e_\beta(t^{\langle\beta, \alpha^\vee\rangle}u) \\ w_\alpha e_\beta(u)w_\alpha^{-1} &= e_{w_\alpha\beta}(u). \end{aligned}$$

Consider

$$A(\Lambda, \chi, w)f(g) = \int_{N(w)} f(w^{-1}ng) dn = \int_{N(w)} f(w^{-1}nw w^{-1}g) dn = \int_{\overline{N}(w)} f(nw^{-1}g) dn,$$

where $\overline{N}(w) = w^{-1}Uw \cap \overline{U}$. Here each element of $\overline{N}(w)$ can be expressed in the form $\prod_{\alpha>0, w\alpha<0} e_{-\alpha}(u_\alpha)$, where $u_\alpha \in F$ (in any order). Let $w = w_{\alpha_n} w'$. Then

$$\{\alpha > 0, w\alpha < 0\} = \{w'^{-1}(\alpha_n)\} \cup \{\beta > 0, w'\beta < 0\},$$

and $\overline{N}(w) = \overline{N}(w') \cdot U_{-w'^{-1}(\alpha_n)}$. Hence

$$A(\Lambda, \chi, w)f(g) = \int_{U_{-w'^{-1}(\alpha_n)}} \int_{\overline{N}(w')} f(n'e_{-w'^{-1}(\alpha_n)}(u)w'^{-1}w_{\alpha_n}g) dn' du.$$

Here $n'e_{-w'^{-1}(\alpha_n)}(u)w'^{-1} = n'w'^{-1}w'e_{-w'^{-1}(\alpha_n)}(u)w'^{-1} = n'w'^{-1}e_{-\alpha_n}(u)$. Hence

$$\begin{aligned} A(\Lambda, \chi, w)f(g) &= \int_{U_{-\alpha_n}} \int_{\overline{N}(w')} f(n'w'^{-1}e_{-\alpha_n}(u)w_{\alpha_n}g) dn' du \\ &= \int_{U_{-\alpha_n}} A(\Lambda, \chi, w')f(e_{-\alpha_n}(u)w_{\alpha_n}g) du = A(w'\Lambda, w'\chi, w_{\alpha_n})A(\Lambda, \chi, w')f(g). \end{aligned}$$

The above cocycle relation of intertwining operators is quite general. Suppose $\mathbf{P} = P_\theta$, $\theta \in \Delta$. Let $A(\Lambda, \pi, \theta, w)$ be the intertwining operator for $\text{Ind}_P^G \pi \otimes \exp(\langle \Lambda, H_P(\cdot) \rangle)$, where $w \in W$ is such that $w\theta \in \Delta$. We can write $A(\Lambda, \pi, \theta, w)$ as a product of rank-one intertwining operators (i.e., intertwining operators for maximal parabolic subgroups).

Theorem 6.5. *There exists a finite sequence $\alpha_1, \dots, \alpha_n \in \Delta$ such that $\Omega_i = \theta_i \cup \{\alpha_i\}$, $\theta_{i+1} = w_i \theta_i$, where $w_i(\theta_i) \subset \Omega_i$, $w_i(\alpha_i) < 0$. Then $w = w_n \cdots w_1$, and*

$$A(\Lambda, \pi, \theta, w) = A(w_{n-1} \cdots w_1 \Lambda, w_{n-1} \cdots w_1 \pi, \theta_n, w_n) \cdots A(\Lambda, \pi, \theta, w_1).$$

Here let $\Phi(A_\theta, G)$ be the set of roots of G with respect to A_θ , and $\Phi_r(A_\theta, G)$ be the set of reduced roots. Let $\Phi_r^+(\theta, w) = \{\alpha \in \Phi_r^+(A_\theta, G) : w\alpha < 0\}$. Then

$$\Phi_r^+(\theta, w) = \{\alpha_1, w_1^{-1}\alpha_2, (w_2 w_1)^{-1}\alpha_3, \dots, (w_{n-1} \cdots w_1)^{-1}\alpha_n\}.$$

We explain the theorem using one example. Suppose $\mathbf{P} = \mathbf{MN} \subset Sp(2n)$, where $\mathbf{M} \simeq GL_{n_1} \times \cdots \times GL_{n_k} \times Sp(2l)$, and $\pi_1 \otimes \cdots \otimes \pi_k \otimes \sigma$ be a representation of $\mathbf{M}(F)$. Let

$$Ind_{\mathbf{P}}^G \pi_1 |det|^{s_1} \otimes \cdots \otimes \pi_k |det|^{s_k} \otimes \sigma,$$

be the induced representation. Then $A(\Lambda, \pi_1 \otimes \cdots \otimes \pi_k \otimes \sigma, w)$ is a product of $A(s_i, \pi_i \otimes \sigma)$, $i = 1, \dots, k$, $A(s_i - s_j, \pi_i \otimes \pi_j)$, $A(s_i + s_j, \pi_i \otimes \tilde{\pi}_j)$, $1 \leq i < j \leq k$, where

$A(s_i, \pi_i \otimes \sigma)$ is the rank-one operator for $Ind \pi_i |det|^{s_i} \otimes \sigma$, $GL_{n_i} \times Sp(2l) \subset Sp(2(n_i + l))$
 $A(s_i - s_j, \pi_i \otimes \pi_j)$ is the rank-one operator for $Ind \pi_i |det|^{s_i} \otimes \pi_j |det|^{s_j}$, $GL_{n_i} \times GL_{n_j} \subset GL_{n_i + n_j}$
 $A(s_i + s_j, \pi_i \otimes \tilde{\pi}_j)$ is the rank-one operator for $Ind \pi_i |det|^{s_i} \otimes \tilde{\pi}_j |det|^{-s_j}$, $GL_{n_i} \times GL_{n_j} \subset GL_{n_i + n_j}$

We return to spherical representations and maximal parabolic case. Let π_v be a spherical representation of $\mathbf{M}(F_v)$, $\mathbf{P} = \mathbf{MN}$ is a maximal parabolic subgroup. Let χ_v be the inducing unramified character. Then

$$A(s\tilde{\alpha}, \pi_v, w_0) f_v^0(e) = \prod_{\beta > 0, w_0 \beta < 0} \frac{L(s\langle \tilde{\alpha}, \beta^\vee \rangle, \chi_v \circ \beta^\vee)}{L(1 + s\langle \tilde{\alpha}, \beta^\vee \rangle, \chi_v \circ \beta^\vee)}.$$

Langlands observed that $\langle \tilde{\alpha}, \beta^\vee \rangle = i$, $i = 1, \dots, m$: Let V_i be the subspace of ${}^L \mathbf{n}$, generated by $E_{\beta^\vee} \in {}^L \mathbf{n}$, $\langle \tilde{\alpha}, \beta^\vee \rangle = i$, where ${}^L \mathbf{n}$ is the Lie algebra of ${}^L \mathbf{N}$. For each i , the adjoint action of ${}^L M$ leaves V_i stable. Let r be the adjoint representation of ${}^L M$ on ${}^L \mathbf{n}$, and $r_i = r|_{V_i}$.

Theorem 6.6 (Langlands-Shahidi). *r_i is irreducible for each i , and the weights of r_i are the roots β^\vee in ${}^L \mathbf{n}$ which restricts to $i\alpha^\vee$ in ${}^L A$.*

Let $r = \bigoplus_{i=1}^m r_i$. Let t_v be the semi-simple conjugacy class in ${}^L T$. Recall that $\chi_v \circ \beta^\vee(\varpi_v) = \beta^\vee(t_v)$.

Theorem 6.7 (Langlands-Shahidi).

$$A(s\tilde{\alpha}, \pi_v, w_0) f_v^0(e) = \prod_{i=1}^m \frac{L(is, \pi_v, r_i)}{L(1 + is, \pi_v, r_i)},$$

where

$$L(s, \pi_v, r_i) = \prod_{\beta > 0, \langle \tilde{\alpha}, \beta^\vee \rangle = i} L(s, \chi_v \circ \beta^\vee), \quad L(s, \chi_v \circ \beta^\vee) = (1 - \chi_v \circ \beta^\vee(\varpi_v) q_v^{-s})^{-1}.$$

Examples 6.8. (1) $\mathbf{G} = GL_{k+l}$, and $\mathbf{M} \simeq GL_k \times GL_l$, $\mathbf{N} = \begin{pmatrix} I_k & X \\ 0 & I_l \end{pmatrix}$. Then ${}^L M \simeq GL_k(\mathbb{C}) \times GL_l(\mathbb{C})$, and $r(\text{diag}(g_1, g_2))X = g_1 X g_2^{-1}$. Here r is irreducible. So in this case $m = 1$. More explicitly, let $\pi_1 = \pi(\mu_1, \dots, \mu_k)$, $\pi_2 = \pi(\nu_1, \dots, \nu_l)$ be spherical representations of $GL_k(F)$, $GL_l(F)$, resp. Then

$$L(s, \pi_1 \otimes \pi_2, r) = L(s, \pi_1 \times \tilde{\pi}_2) = \prod_{i=1}^k \prod_{j=1}^l L(s, \mu_i \nu_j^{-1}).$$

(2) $\mathbf{G} = Sp(2n)$, $\mathbf{M} = GL_n$ (Siegel parabolic subgroup). Then \mathbf{N} is generated by $\{e_i + e_j, i < j, 2e_i, i = 1, \dots, n\}$, and $\tilde{\alpha} = e_1 + \dots + e_n$. We can see that $m = 2$; $r = r_1 \oplus r_2$; ${}^L \mathbf{n} = V_1 \oplus V_2$, where V_1 is spanned by $e_i, i = 1, \dots, n$ (note that e_i is the coroot of $2e_i$); V_2 is spanned by $e_i + e_j, i < j$. Hence, for a spherical representation π of $GL_n(F)$,

$$L(s, \pi, r_1) = L(s, \pi), \text{ standard } L\text{-function of } GL_n$$

$$L(s, \pi, r_2) = L(s, \pi, \wedge^2), \text{ exterior square } L\text{-function of } GL_n$$

(3) Consider $E_7 - 1$ case in Example 1.36. We saw that there is a rational injection $f : \mathbf{M} \longrightarrow GL_2 \times GL_3 \times GL_4$. Then f induces an injection $f : \mathbf{M}(\mathbb{A}) \longrightarrow GL_2(\mathbb{A}) \times GL_3(\mathbb{A}) \times GL_4(\mathbb{A})$ such that $\mathbf{M}(\mathbb{A})(\mathbb{A}^*)^2$ is co-compact in $GL_2(\mathbb{A}) \times GL_3(\mathbb{A}) \times GL_4(\mathbb{A})$, where $(\mathbb{A}^*)^2$ is embedded as a center of the first two factors.

Let $\pi_i, i = 1, 2, 3$, be cuspidal representations of $GL_{1+i}(\mathbb{A})$ with central characters ω_i , resp. Then $\pi_1 \otimes \pi_2 \otimes \pi_3|_{f(\mathbf{M}(\mathbb{A}))}$ decomposes as a direct sum of irreducible representations of $\mathbf{M}(\mathbb{A})$. Let π be any cuspidal irreducible constituent of this direct sum. Then $\omega_\pi = \omega_1^6 \omega_2^4 \omega_3^3$. Let $\pi = \otimes_v \pi_v$. If π_v is a spherical representation, we can see that $m = 4$, and

$$L(s, \pi_v, r_1) = L(s, \pi_{1v} \times \pi_{2v} \times \pi_{3v}),$$

$$L(s, \pi_v, r_2) = L(s, \tilde{\pi}_{2v} \otimes \pi_{3v}, (\rho_3 \otimes \omega_1 \omega_2) \otimes \wedge^2 \rho_4),$$

$$L(s, \pi_v, r_3) = L(s, (\pi_{1v} \otimes \omega_1 \omega_2 \omega_3) \times \tilde{\pi}_{3v}),$$

$$L(s, \pi_v, r_4) = L(s, \pi_{2v} \otimes \omega_1^2 \omega_2 \omega_3).$$

List of L -functions via Langlands-Shahidi method (split reductive groups).

A_n case.

$GL_m \times GL_n \subset GL_{m+n}$ gives the Rankin-Selberg L -function $L(s, \pi_1 \times \pi_2)$, where π'_i 's are cuspidal representations of GL_m, GL_n , resp.

B_n case.

$B_n - 1$: $m = 2$; r_1 gives the Rankin-Selberg L -function of $GL_k \times GSpin(2l+1)$; r_2 gives the twisted symmetric square L -function of GL_k . If $G = SO(2n+1)$, r_1 gives the Rankin-Selberg L -function of $GL_k \times SO(2l+1)$; r_2 gives the symmetric square L -function of GL_k .

 C_n case.

$C_n - 1$: $m = 2$; r_1 gives the Rankin-Selberg L -function of $GL_k \times Sp(2l)$; r_2 gives the exterior square L -function of GL_k , if $k \neq 1$ (If $k = 1$, then $m = 1$).

 D_n case.

$D_n - 1$: $m = 2$; r_1 gives the Rankin-Selberg L -function of $GL_k \times GSpin(2l)$; r_2 gives the twisted exterior square L -function of GL_k . If $G = SO(2n)$, r_1 gives the Rankin-Selberg L -function of $GL_k \times SO(2l)$; r_2 gives the exterior square L -function of GL_k (If $k = 1$, then $m = 1$).

$D_n - 2$: $m = 2$; r_1 gives the triple L -function of $GL_{n-2} \times GL_2 \times GL_2$; r_2 gives the twisted exterior square L -function of GL_{n-2} .

$D_n - 3$: $m = 2$; r_1 gives $L(s, \sigma \otimes \tau, \rho_{n-3} \otimes \wedge^2 \rho_4)$; r_2 gives the twisted exterior square L -function of GL_{n-3} .

 F_4 case.

$F_4 - 1$: $m = 4$; r_1 gives $L(s, \sigma \times \tau)$ which is entire; r_2 gives $L(s, \sigma \otimes \tau, Sym^2 \rho_2 \otimes \rho_3)$ which has a pole at $s = 1$ when $\tilde{\tau} \simeq Sym^2 \sigma$.

$F_4 - 2$: $m = 3$; r_1 gives $L(s, \sigma \otimes \tau, Sym^2 \rho_3 \otimes \rho_2)$; r_2 gives $L(s, \sigma, Sym^2 \rho_3 \otimes \omega_\tau)$ which has a pole at $s = 1$ always.

(xviii): $m = 2$; $M = GSpin(7) \subset F_4$; $dim r_2 = 1$; r_1 is the 14-dim'l irreducible representation of $Sp_6(\mathbb{C})$, called spherical harmonic.

(xix): $m = 2$; $M = GSp_6 \subset F_4$; $r_1 = 8$ -dim'l spin representation of $Spin(7, \mathbb{C})$; r_2 gives the standard L -function of $SO_7(\mathbb{C})$ (7-dimensional).

 E_6 case.

$E_6 - 1$: $m = 3$; r_1 gives the triple L -function of $GL_3 \times GL_2 \times GL_3$; r_2 gives the standard L -function of $GL_3 \times GL_3$.

$E_6 - 2$: $m = 2$; $r_1 = \wedge^2 \rho_5 \otimes \rho_2$; r_2 gives the Rankin-Selberg L -function of $GL_5 \times GL_2$ which is entire.

(x): $m = 2$; $dim r_2 = 1$; r_1 gives the exterior cube L -function of $GL_6(\mathbb{C})$ (20 dim'l irreducible representation of $GL_6(\mathbb{C})$)

(xi): $m = 1$; $r_1 = 16$ -dimensional half-spin representation of $Spin(10)$.

 E_7 case.

$E_7 - 1$: $m = 4$; r_1 gives the triple L -function of $GL_3 \times GL_2 \times GL_4$; r_2 comes from $D_6 - 3$ case.

$E_7 - 2$: $m = 3$; $r_1 = \wedge^2 \rho_5 \otimes \rho_3$; r_2 gives the Rankin-Selberg L -function of $GL_5 \times GL_3$ which is entire.

$E_7 - 3$: $m = 2$; r_1 gives the L -function $L(s, \sigma \otimes \tau, Spin^{16} \otimes \rho_2)$, where σ, τ are cuspidal representations of $GSpin_{10}, GL_2$, resp. and $Spin^{16}$ is the 16-dimensional half-spin representation of $Spin_{16}(\mathbb{C})$.

$E_7 - 4$: $m = 3$; r_1 gives the L -function $L(s, \sigma \otimes \tau, \wedge^2 \rho_6 \otimes \rho_2)$, where σ, τ are cuspidal representations of GL_6, GL_2 , resp.

(xi): $m = 2$; r_1 gives the exterior cube L -function of $GL_7(\mathbb{C})$ (35-dimensional representation of $GL_7(\mathbb{C})$); r_2 gives the standard L -function of $GL_7(\mathbb{C})$ which is entire.

($xxvi$): $m = 2$; $\dim r_2 = 1$ and r_1 gives the degree $32 = 2^5$ spin L -function of $Spin_{12}$

(xxx): $m = 1$; r_1 gives the standard L -function of E_6 .

E_8 case.

$E_8 - 1$: $m = 6$; r_1 gives the triple L -function of $GL_3 \times GL_2 \times GL_5$; r_2 comes from $E_7 - 2$ case.

$E_8 - 2$: $m = 5$; r_1 gives the L -function $L(s, \sigma \otimes \tau, \wedge^2 \rho_5 \otimes \rho_4)$, where σ, τ are cuspidal representations of GL_5, GL_4 , resp.

$E_8 - 3$: $m = 4$; r_1 gives the L -function $L(s, \sigma \otimes \tau, Spin^{16} \otimes \rho_3)$, where σ, τ are cuspidal representations of $GSpin_{10}, GL_3$, resp. and $Spin^{16}$ is the 16-dimensional half-spin representation of $Spin_{16}(\mathbb{C})$.

$E_8 - 4$: $m = 3$; r_1 gives the standard L -function of $E_6 \times GL_2$; r_2 gives the standard L -function of E_6 ((xxx) case).

$E_8 - 5$: $m = 4$; r_1 gives the L -function $L(s, \sigma \otimes \tau, \wedge^2 \rho_7 \otimes \rho_2)$, where σ, τ are cuspidal representations of GL_7, GL_2 , resp.

($xiii$): $m = 3$; r_1 gives the degree 56 exterior cube L -function of GL_8

($xxviii$): $m = 2$; r_1 gives the degree $64 = 2^6$ spin L -function of $Spin_{14}$

($xxxi$): $m = 2$; $\dim r_2 = 1$; r_1 gives the standard L -function of E_7 .

G_2 case.

(xv) (attached to the maximal parabolic subgroup generated whose unipotent radical contains the long simple root subgroup) : $m = 2$; $\dim r_2 = 1$, r_1 gives the third symmetric power L -function of GL_2

(xvi) (attached to the maximal parabolic subgroup generated whose unipotent radical contains the short simple root subgroup): $m = 3$; $\dim r_2 = 1$, r_1 gives the standard L -function of GL_2